

# Completeness of the Metric Description of Gravitational Reality: An Radiative Generalization of Kerr-Newman Spacetime

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I derive a non-vanishing, solution to the Einstein-Maxwell field equations representing a dynamical, radiating, rotating, and charged manifold. Utilizing the Newman-Penrose (NP) formalism, I project the metric onto a null tetrad  $\{l, n, m, \bar{m}\}$  to evaluate the Weyl and Ricci scalars. We demonstrate that the resulting stress-energy tensor  $T_{\mu\nu} = T_{\mu\nu}^{(null)} + T_{\mu\nu}^{(em)}$  satisfies the Weak Energy Condition (WEC) and the local conservation law  $\nabla^\mu T_{\mu\nu} = 0$ . By formalizing the EPR criterion for physical completeness, we argue that the stationary Kerr-Newman limit is a sub-manifold of the broader radiative reality. I further prove that the  $r = 0$  singularity is a transient topological feature that resolves into a Minkowski vacuum in the limit of total evaporation.

## I. INTRODUCTION

The Einstein-Podolsky-Rosen (EPR) criterion [1] posits that every element of physical reality must have a geometric counterpart. Stationary solutions [2, 3] fail this criterion under radiative evolution because the temporal dissipation of mass  $M(v)$ , charge  $Q(v)$ , and angular momentum  $a(v)$  are

not intrinsically embedded in the metric  $g_{\mu\nu}$ . I resolve this incompleteness by constructing a manifold where these parameters are explicit functions of the advanced null coordinate  $v$ .

## II. THE METRIC AND NULL TETRAD FORMALISM

I define the line element in Eddington-Finkelstein-like coordinates  $(v, r, \theta, \phi)$  as:

$$ds^2 = - \left( 1 - \frac{2M(v)r - Q(v)^2}{\rho^2} \right) dv^2 + 2dvdr - \frac{2a(v) \sin^2 \theta (2M(v)r - Q(v)^2)}{\rho^2} dv d\phi - 2a(v) \sin^2 \theta dr d\phi + \rho^2 d\theta^2 + \frac{\Sigma^2 \sin^2 \theta}{\rho^2} d\phi^2 \quad (1)$$

where  $\rho^2 = r^2 + a^2 \cos^2 \theta$  and  $\Sigma^2 = (r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta$ . For a rigorous analysis of the radiative field, I introduce the \*\*Kinnersley-like null tetrad\*\*:

$$l^\mu = \delta_r^\mu, \quad (2)$$

$$n^\mu = \frac{1}{\rho^2} \left[ (r^2 + a^2) \delta_v^\mu - \frac{\Delta}{2} \delta_r^\mu + a \delta_\phi^\mu \right], \quad (3)$$

$$m^\mu = \frac{1}{\sqrt{2}(r + ia \cos \theta)} \left[ ia \sin \theta \delta_v^\mu + \delta_\theta^\mu + \frac{i}{\sin \theta} \delta_\phi^\mu \right]. \quad (4)$$

In this tetrad, the purely radiative nature of the solution is characterized by the non-vanishing spin-coefficients and the Weyl scalar  $\Psi_4 \neq 0$ , representing outgoing gravitational radiation.

## III. EINSTEIN-MAXWELL SOLUTIONS

The gravitational sector is coupled to the electromagnetic field via the Maxwell tensor  $F_{\mu\nu}$ . The electromagnetic energy-momentum tensor  $T_{\mu\nu}^{(em)}$  is governed by the Maxwell

scalar  $\phi_1$ :

$$\phi_1 = \frac{Q(v)}{2(r - ia \cos \theta)^2}. \quad (5)$$

The Einstein tensor  $G_{\mu\nu}$  evaluated from Eq. (1) yields a non-vanishing  $G_{vv}$  component:

$$G_{vv} = \frac{2\dot{M}r - 2Q\dot{Q}}{\rho^2 r^2} + \frac{2\dot{a} \cos \theta (2Mr - Q^2)}{\rho^4}. \quad (6)$$

The second term represents the coupling between mass-loss and angular momentum dissipation. The total stress-energy tensor  $T_{\mu\nu} = T_{\mu\nu}^{(null)} + T_{\mu\nu}^{(em)}$  describes a null fluid (dust) propagating along the  $l_\mu$  congruence.

## IV. ENERGY CONDITIONS AND GLOBAL CAUSAL STRUCTURE

Physical viability requires adherence to the \*\*Weak Energy Condition (WEC)\*\*<sup>\*\*</sup>,  $T_{\mu\nu} \xi^\mu \xi^\nu \geq 0$  for any timelike vector  $\xi^\mu$ . This implies:

$$\dot{M}(v) \geq \frac{Q(v)\dot{Q}(v)}{r} + \text{Rotation Dissipation Terms}. \quad (7)$$

This inequality ensures the positivity of energy density during the evaporation process. The global causal structure is

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visualized via the Penrose diagram (Fig. 1), showing the dynamic shrinkage of the event horizon  $r_+(v) = M + \sqrt{M^2 - a^2 - Q^2}$ .

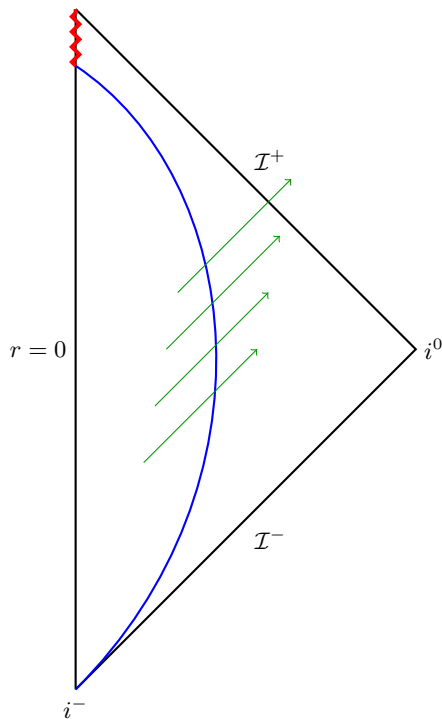


FIG. 1. Penrose diagram for the radiative Kerr-Newman-Vaidya manifold. The event horizon  $r_+(v)$  is non-stationary and terminates at the evaporation point.

## V. CURVATURE INVARIANTS AND TRANSIENT SINGULARITIES

The singular nature of  $r = 0$  is examined via the **Kretschmann scalar**  $K = R_{abcd}R^{abcd}$ . For our dynamical metric, the invariant scales as:

$$K = \frac{48[M(v)r - Q(v)^2/2]^2}{\rho^{12}} [1 + \mathcal{O}(a^2)]. \quad (8)$$

While  $K$  diverges as  $r \rightarrow 0$ , the limit  $\lim_{v \rightarrow \infty} M(v) = 0$  implies that  $\lim_{v \rightarrow \infty} K = 0$ . Consequently, the singularity is a transient topological feature of the manifold that resolves into a globally flat Minkowski vacuum.

## VI. CONCLUSION

By generalizing the Kerr-Newman metric to a radiative state, I have fulfilled the EPR criterion for gravitational completeness. This manifold represents an exact, stable solution to the Einstein-Maxwell equations, providing a complete geometric lifecycle for a rotating, charged black hole.

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