

Dark Mass is Potential Energy

Nicolas Poupart, Independent Researcher (2025)
99 Chemin Scraire, Mille-Isles, Québec, Canada (J0R 1A0)
(450) 939-2167
nicolas.poupart@yahoo.fr

Abstract — This paper demonstrates that gravitational potential energy ($E = -GmM/d$) accounts for the phenomenon commonly attributed to dark matter, by applying the mass-energy equivalence relation ($E = mc^2$). No additional assumptions are made beyond the principle of volume conservation, ensuring that the gravitational field remains conservative and uninfluenced by forces other than gravity. We develop a straightforward equation and algorithm to accurately compute the potential energy of a stellar system. The theoretical implications of this model are explored with respect to energy production by various types of stars and black holes in galaxies. The model is empirically tested against the SPARC database comprising 175 galaxies to assess its validity. We conclude by deriving the logical and mathematical consequences of the hypothesis that dark mass is stored as potential energy in the gravitational field, a claim further validated using the SPARC data. The final consequence being the production of dark energy and the reproduction of Λ CDM cosmology.

Introduction

Since Vera Rubin postulated the existence of galactic dark mass to explain the flatness of galactic rotation curves (1; 2; 3), no convincing explanation for the nature of this mass has been provided. Attempts to attribute the missing mass to undetectable baryonic matter were largely refuted by the AGAPE (4), MACHO (5), and EROS (6) programs. Similar explanations involving non-baryonic or exotic particles have also failed to account for the discrepancy. Numerous detection efforts, including those by the LUX (7), PICASSO (8), PICO (9), and SuperCDMS (10) collaborations, have thus far been unsuccessful. Likewise, results from CERN's latest accelerator suggest that physics remains consistent with the Standard Model, making the existence of exotic particles increasingly unlikely. It also remains extremely difficult to explain this phenomenon using current gravitational theory, whether Newtonian gravitation or general relativity.

In cosmology, the prevailing framework is the Λ CDM model, which postulates the existence of cold dark matter. For most physicists, the concept of mass remains inseparable from that of matter. Moreover, the term “dark matter” is often used categorically, although it more accurately refers to “dark mass.” One alternative to the existence of real dark mass is to modify gravitation at galactic scales. However, such efforts are challenged by the remarkable empirical adequacy of general relativity in describing observed phenomena and the inferred presence of dark mass (11; 12; 13).

The explanation proposed in this article takes a fundamentally different approach. Dark mass is neither a form of real matter nor the result of modified gravitation. Rather, it is a secondary effect inherent to the current formulation of gravitational theory—specifically, the gravitational potential energy stored in the gravitational field. The only axioms employed are $E = -GmM/d$ and $E = mc^2$. Thus, the explanation relies solely on the judicious application of classical physics.

Useful Gravitational Potential Energy

If we consider the Newtonian gravitational force equation, $F = GmM/d^2$, we obtain—by integrating from d to infinity—the potential energy formula $E_p = -GmM/d$. While this formulation yields negative potential energy which can be used for calculating the motion of celestial bodies, it is unsuitable for evaluating the total physical potential energy of a system.

To determine the physically meaningful potential energy, we must compute the energy difference between two states of the system—analogueous to lifting a mass m from position d by a height h . In such a case: $\Delta E_p = E_p(d+h) - E_p(d) = GmMh/(d^2 + dh)$. This expression represents the *usable* form of gravitational potential energy. The concept of absolute negative energy E_p is not directly interpretable in physical terms.

However, this formulation becomes impractical for systems involving two celestial bodies (e.g., planets, stars, or galaxies), because the initial reference distance d is undefined or ambiguous. To resolve this, we compute the energy difference between a compact state—a single solid spherical mass $M_t = m + M$ and a configuration of two distinct spherical bodies, m and M , separated by a distance d .

Let the two bodies have radii r and R , and volumes $v = 4\pi r^3/3$ and $V = 4\pi R^3/3$, with corresponding densities m/v and M/V . The total volume of the single solid sphere is $V_t = V + v = 4\pi/3 (R^3 + r^3)$, and its radius $R_t = (R^3 + r^3)^{1/3}$. Assuming equal volumetric densities, $m/v = M/V$, that of the single compact ball will be identical $m/v = M/V = M_t/V_t$; if not, it assumes an average density weighted by the contributions of m and M .

The gravitational potential energy of a homogeneous solid sphere is given by $E_p = -3GM^2 / 5R$. The gravitational potential energy of the system can be evaluated by comparing the initial state (a single compact mass), the final state (two separated masses), and the difference between them:

$$E_i = -\frac{3GM_t^2}{5R_t}$$

$$E_f = -\frac{3GM^2}{5R} - \frac{3Gm^2}{5r} - \frac{GMm}{d}$$

$$\Delta E_p = E_f - E_i = \left(-\frac{3GM^2}{5R} - \frac{3Gm^2}{5r} - \frac{GMm}{d} \right) + \frac{3GM_t^2}{5R_t}$$

This expression quantifies the energy change from a compact configuration to a system of two separated masses. It represents a redistribution of matter that preserves the original component densities, thereby ensuring that the transformation involves only gravitational forces.

Gravitational Potential Energy is Massive

Gravitational potential energy can be expressed as $E_p = m_p c^2$, indicating that it must be stored as mass within the system. This statement is readily illustrated through a simple thought experiment. Consider a nuclear reactor that converts a mass m_n into electrical energy, which is then used to raise a mass m to a height h .

There should be no debate regarding whether the Earth's mass remains unchanged before and after the transformation. Given the conservation of mass-energy m_n within the Earth system, general relativity ensures that its gravitational field remains unchanged. The question of whether this constitutes “real” mass is meaningless, since mass is defined by its measurable gravitational or inertial effects.

Moreover, since the equivalence of inertial and gravitational mass has never been empirically violated (14), we may, for now, treat gravitational potential energy as fully equivalent to its mass representation.

Calculation of the Potential Energy of Celestial Bodies

	m (kg)	M (kg)	r (m)	R (m)	d (m)	R_t (m)
Moon + Earth	7.348E+22	5.972E+24	1.738E+06	6.378E+06	3.844E+08	6.421E+06
Earth + Sun	5.972E+24	1.989E+30	6.378E+06	6.963E+08	1.496E+11	6.963E+08
Jupiter + Sun	1.898E+27	1.989E+30	6.991E+07	6.963E+08	7.780E+11	6.966E+08
Sun + M80	1.989E+30	9.985E+35	6.963E+08	5.534217E+10	4.541E+17	5.534221E+10
Sun + Galaxy	1.989E+30	5.000E+40	6.963E+08	2.040E+12	5.000E+20	2.040E+12

This table contains the standard values for the celestial bodies under consideration, except for the radii for the globular cluster M80 and a representative galaxy (indicated by gray boxes). We modeled the initial state as a solid sphere of mass $M + m$, assuming solar density. The final state is the sun at distance d of a mass M of solar density. For the Sun–galaxy pair, the values used are, $R = 2.03985712710655 \times 10^{12}$ and $R_t = 2.03985712713359 \times 10^{12}$. The rationale behind this construction will be addressed in a subsequent section. In either case, the radius is calculated using the relation: $R_t = (R^3 + r^3)^{1/3}$.

	E_i (j)	E_f (j)	ΔE_p (j)	ΔE_p (kg)	$m / \Delta E_p$
Moon + Earth	-2.2795E+32	-2.2413E+32	3.8176E+30	4.2477E+13	0.000000001
Earth + Sun	-2.2751E+41	-2.2751E+41	1.3024E+36	1.4491E+19	0.000002427
Jupiter + Sun	-2.2787E+41	-2.2751E+41	3.5519E+38	3.9520E+21	0.000002082
Sun + M80	-7.2141E+50	-7.2141E+50	2.3949E+45	2.6647E+28	0.013396966
Sun + Galaxy	-4.9079E+58	-4.9079E+58	3.2540E+48	3.6206E+31	18.20

The table presents the gravitational potential energy differences between current physical configurations and their hypothetical fusion into a single mass. For context, the complete annihilation of one kilogram of matter yields an energy output approximately equivalent to that of a hydrogen bomb. For example, 2 kg of deuterium fused with 3 kg of tritium results in a mass loss of roughly 1 kg—comparable to the energy released by the Tsar Bomba.

Based on this equivalence, the merger of the Moon with the Earth would release energy equivalent to approximately 40 trillion hydrogen bombs. This potential energy corresponds to a mass of about 40 billion metric tons—a value that is far from negligible.

If we consider the energy released by merging the Earth into the Sun, the resulting energy would correspond to over 14 million billion metric tons. However, this value is still less than the mass the Sun loses each hour due to nuclear fusion. Depending on its spatial distribution, this energy-equivalent mass could, in principle, be measurable.

If we examine the ratio $m/\Delta E_p$, representing the mass of the smaller body relative to the potential energy of the system, we find that it is negligible in most cases—except at the galactic scale, where it approaches the order of magnitude of the observed dark mass ratio. For clusters and superclusters, this ratio is on the order of 1,000 and 10,000, respectively. However, as seen in the galactic case, the calculation reaches the precision limit of double-precision floating-point arithmetic (approximately 15 significant digits).

The Potential Energy of Celestial Systems

The first observation is that the term $-GmM/d$ of ΔE_p is negligible at all relevant scales, where $d \gg R \gg r$. This result is intuitive from a physical standpoint: altering the distance between two celestial bodies, such as the Earth and the Moon, leads to minimal changes in gravitational energy compared to the energy released by their hypothetical fusion. This approximation holds across all celestial systems.

The energy expression thus simplifies to:

$$\Delta E_p = \frac{3G}{5} \left(\frac{M_t^2}{R_t} - \frac{M^2}{R} - \frac{m^2}{r} \right)$$

This implies that the spatial arrangement of bodies relative to one another is irrelevant for calculating total gravitational potential energy. If we consider a system of n masses m_i with radii r_i , merged sequentially, we obtain the following expressions: $M_i = \sum_{j \leq i} m_j$, $R_i = (R_{i-1}^3 + r_i^3)^{1/3}$ and

$$\Delta E_i = \frac{3G}{5} \left(\frac{M_i^2}{R_i} - \frac{M_{i-1}^2}{R_{i-1}} - \frac{m_i^2}{r_i} \right), \Delta E_p = \sum_{i \leq n} \Delta E_i$$

It is important to note that the order of mergers is irrelevant, as the gravitational field is conservative.

Simplified Calculation with Identical Mass and Radius

By assuming identical mass m and radius r for each of the n bodies, the calculation simplifies considerably. In this case, the cumulative mass at the i^{th} merger step is $M_i = im$, and the corresponding radius is:

$$R_i = \left(\frac{M_i r^3}{m} \right)^{1/3}$$

since M_i retains the same density as the original body of mass m .

The change in gravitational potential energy at each step is given by:

$$\Delta E_i = \frac{3G}{5} \left(\frac{M_i^2}{R_i} - \frac{M_{i-1}^2}{R_{i-1}} - \frac{m^2}{r} \right), \Delta E_p = \sum_{i \leq n} \Delta E_i$$

Due to cancellation of intermediate terms, the telescoping sum simplifies to:

$$\Delta E_p = \frac{3G}{5} \left(\frac{M_n^2}{R_n} - n \frac{m^2}{r} \right)$$

Since $M_n = nm$, we find:

$$M_n^2 = n^2 m^2 \quad \text{and} \quad R_n = \left(\frac{M_n r^3}{m} \right)^{1/3} = n^{1/3} r$$

Substituting into the equation:

$$\Delta E_p = \frac{3G}{5} \left(\frac{n^2 m^2}{n^{1/3} r} - \frac{nm^2}{r} \right) = \frac{3G}{5} (n^{5/3} - n) \frac{m^2}{r}$$

This final expression shows that the total gravitational potential energy of the system depends solely on the mass, radius, and number of individual bodies. Consequently, a small radius and large mass—i.e., high density—of the component solid spheres significantly increases the total potential energy of the system.

Practical Estimation from a Stellar Population Histogram

In practice, the stellar population of a galaxy is typically known through a mass distribution histogram. The total stellar mass M can be partitioned into k discrete mass intervals, each denoted by a triplet (M_i, m_i, r_i) , $M_i = n_i m_i$ representing the cumulative mass, individual stellar mass and stellar radius respectively.

The first step involves calculating the contribution from each interval as:

$$\Delta E_{pa} = \sum_{i \leq k} \frac{3G}{5} \left(\frac{M_i^2}{R_i} - n_i \frac{m_i^2}{r_i} \right), R_i = \left(\frac{M_i r_i^3}{m_i} \right)^{1/3}$$

where R_i is the radius associated with mass M_i , and n_i is the number of stars of mass m_i in interval i .

Next, we compute the residual potential energy ΔE_{pb} by merging all k compact balls (M_i, R_i) . The total gravitational potential energy of the system is then:

$$\Delta E_p = \Delta E_{pa} + \Delta E_{pb}$$

Alternatively, this process can be formalized via a function $f(m, r, M, M_i, R_i) \mapsto (\Delta E_p, M_t, R_t)$ which takes as input the mass m , radius r , and total mass M of the stellar units, along with an optional compact configuration (M_i, R_i) that may be initialized to zero. The function returns the total potential energy ΔE , along with the updated total mass M_t and radius R_t . The algorithm is simply:

$$\begin{aligned} f(m, r, M, M_i, R_i) &\mapsto (\Delta E_p, M_t, R_t) : \\ n &= M / m \\ M_t &= M + M_i \\ R &= n^{1/3} r \\ R_t &= (R^3 + R_i^3)^{1/3} \\ \Delta E_p &= (3G/5) [M_t^2 / R_t - n m^2 / r] \end{aligned}$$

Hierarchical Computation of Gravitational Potential Energy

To calculate the gravitational potential energy ΔE_p of larger-scale structures—such as galaxy clusters or superclusters—each constituent galaxy must first be reduced to a compact configuration using the previously described method. This process yields three key quantities for each galaxy: the potential energy ΔE_p , the total mass M_t , and the radius R_t . The total mass of the resulting compact object is then given by $M = M_t + \Delta E_p / c^2$ where $\Delta E_p / c^2$ accounts for the mass-equivalent of the stored gravitational potential energy. The radius R of this composite object is now that of the galaxy with the dark mass distributed in the gravitational field.

This fusion procedure can then be reapplied at the next hierarchical scale, treating each galaxy as a single unit in a larger gravitational system—whether a cluster or supercluster. By repeating the process iteratively from smaller to larger scales, one constructs a consistent model of gravitational energy accumulation across cosmic structures.

This hierarchical approach is conceptually analogous to a system of nested dolls, with each level of structure encapsulating the gravitational characteristics of its constituent components.

Treatment of Galactic Gas in Gravitational Potential Energy Calculations

Galactic gas presents a modeling challenge, as it does not form distinct compact objects suitable for direct merger calculations. However, it is known to possess at least as much gravitational potential energy as the stellar component, since it is responsible for star formation through gravitational collapse.

This process can be conceptually divided into two phases: first, the collapse of gas into stellar cores; and second, the merger of these stellar objects into an initial compact configuration. The young stellar population of a galaxy reflects the outcome of this evolutionary pathway and therefore serves as a practical proxy for the gas component. As such, the simplest and most consistent approach is to treat the gas as contributing directly to this population.

Alternatively, if one models the gas as collapsing entirely into the final stellar configuration, its relative sparsity justifies approximating each molecule as nearly isolated. In this scenario, the conservative displacement of individual gas molecules within the compact configuration increases the system's radius while preserving its overall density.

Consequently, the gas should be treated as contributing indistinguishably to the stellar mass distribution. This modeling assumption is supported by our numerical tests and is necessary to minimize errors in the interpretation of observational data.

The Initial Compact Ball Problem

We define the initial compact configuration as a solid sphere that conserves the total volume of the bodies it replaces. This strict conservation of volume ensures physical consistency in the state of matter before and after compaction.

In reality, gravitational forces would compress this configuration even further. However, doing so would introduce additional repulsive forces—such as degeneracy pressure or radiation pressure—that are not included in the current gravitational model. Enforcing volume conservation is therefore the only way to preserve the conservative nature of the gravitational transformation without invoking new interactions.

Any deviation from the volume-conserving radius in an attempt to improve the estimate of potential energy would require a physically justified rationale. Otherwise, it would introduce arbitrariness into the model and compromise its internal consistency.

Empirical consistency is observed when applying the volume-conserving model to systems composed of dense astrophysical bodies. Less dense configurations would underestimate the potential energy, which must be at least equal to the gravitational energy associated with the current structure. Conversely, using a denser configuration leads to exaggerated potential energy values, as confirmed by theoretical models. This is problematic, as the observed quantity of dark mass is not currently underrepresented in our models.

Why the Volume-Conserving Compact Ball Is Preferred over a Black Hole

The exclusion of non-gravitational forces from the model is a foundational assumption. Under this constraint, one might ask whether a black hole could serve as the initial compact configuration for evaluating gravitational potential energy.

To assess this, consider a comparison between the volume-conserving compact ball—denoted ΔE_b —which aggregates the potential energy of all stellar components, and a non-rotating black hole of equivalent mass. The black hole is assumed to have no angular momentum, since the random distribution of stellar rotation vectors leads to approximate cancellation, resulting in a non-rotating configuration.

The gravitational collapse of the compact ball into a black hole—effectively a time-reversed scenario—releases an energy difference given by: $\Delta E_{bh} = \Delta E_b - E_{bh}$

Because the black hole's rest mass accounts for its total energy, and we define this state as the zero-potential-energy reference, we set $E_{bh} = 0$. Hence: $\Delta E_{bh} = \Delta E_b$.

Accordingly, the gravitational potential energy of the system is given by ΔE_b , corresponding to the volume-conserving compact configuration. This approach ensures that no additional forces are introduced, because in the static configurations before and after compaction, any internal forces—such as pressure or other repulsive interactions—are exactly balanced by gravity, in accordance with the principle of action and reaction. Their net contribution is therefore zero. By enforcing volume conservation, we remain within a purely gravitational framework, avoiding the need to model forces that, while present internally, cancel out in the static equilibrium of the system.

Theoretical Model

We apply our gravitational potential energy theory to generate theoretical models in the form of curves, representing the dark mass ratio produced by galaxies with varying numbers of stars. Each curve corresponds to a different galaxy size (Figures 1 and 2).

The x-axis represents the average stellar mass M_s , expressed in solar masses. Stellar radii are determined using the standard mass–radius relationship for main-sequence stars. The y-axis shows the ratio of total gravitationally induced dark mass M_G to the total baryonic stellar mass M_B , that is, M_G / M_B .

For example, a galaxy with 100 billion stars (thin dotted line) yields a dark mass ratio of approximately 28 when the average star has $1 M_\odot$ and $1 R_\odot$. A population of 200 billion stars yields a ratio of 45, and 300 billion stars results in a ratio of 60. These values exceed the commonly accepted ratio of around 20 for the Milky Way.

Nevertheless, the order of magnitude is consistent, despite the exclusion of gas and black hole components in this simplified analysis. This result reinforces the central hypothesis that dark mass is not a distinct form of matter, but rather emerges from the gravitational potential energy of the system.

Impact of Stellar Density Variations on Dark Mass Production

The inclusion of less dense stars, such as red giants, significantly reduces the modeled production of gravitationally induced dark mass (Figures 3 and 4). This scenario is based on the same stellar population used in previous cases, composed primarily of stars with $1 M_\odot$ and $1 R_\odot$, to which a variable fraction of red giants is added, each with $4 M_\odot$ and $100 R_\odot$. The resulting decrease in average stellar density lowers the total gravitational potential energy of the system and, consequently, the dark mass ratio.

Conversely, the addition of dense stellar remnants, such as white dwarfs ($0.6 M_\odot$, $0.0085 R_\odot$) and neutron stars ($1.35 M_\odot$, 10.8 km), leads to a dramatic increase in dark mass production (Figures 5 and 6). The extremely high densities of these objects substantially elevate the system's total potential energy. In realistic stellar populations, however, this increase is moderated by the presence of low-density stars (e.g., red giants) and black holes.

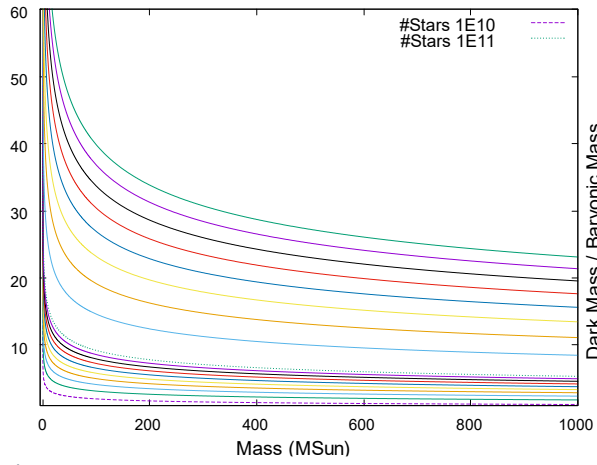


Figure 1

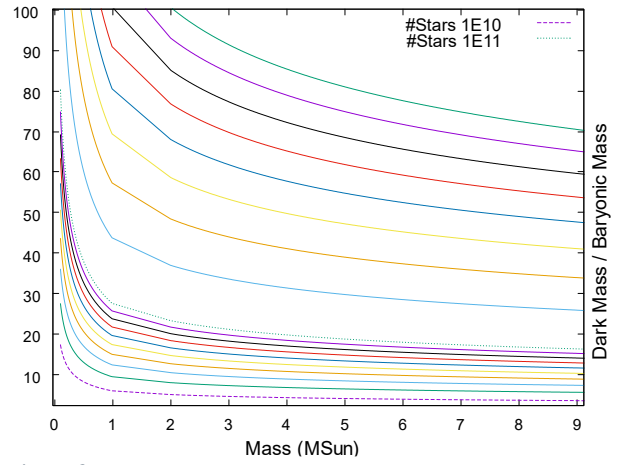


Figure 2

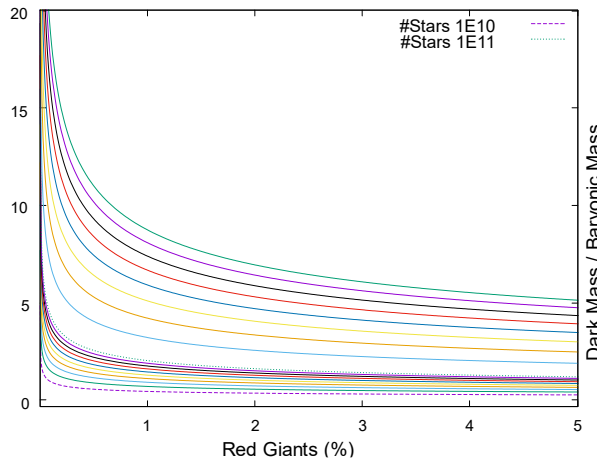


Figure 3

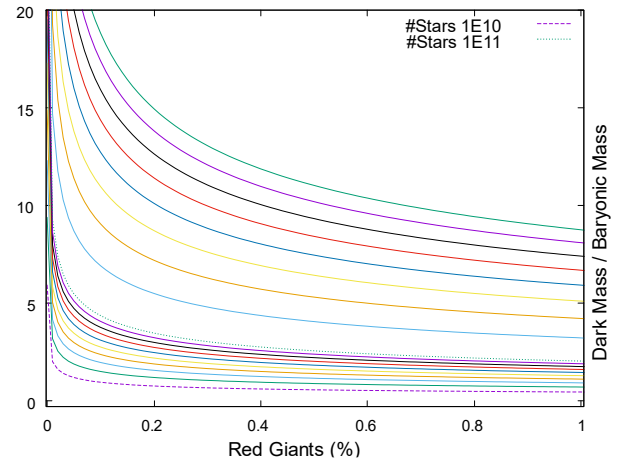


Figure 4

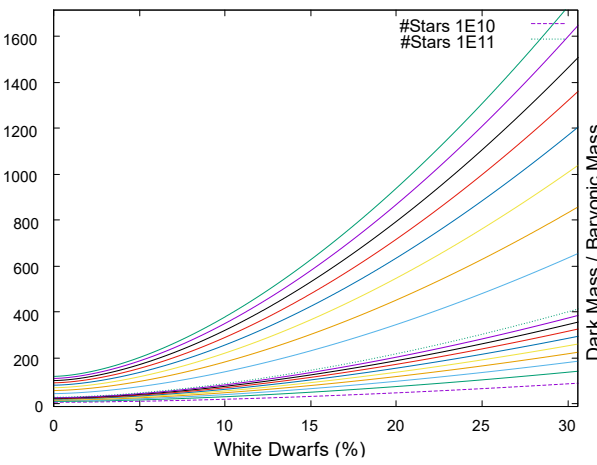


Figure 5

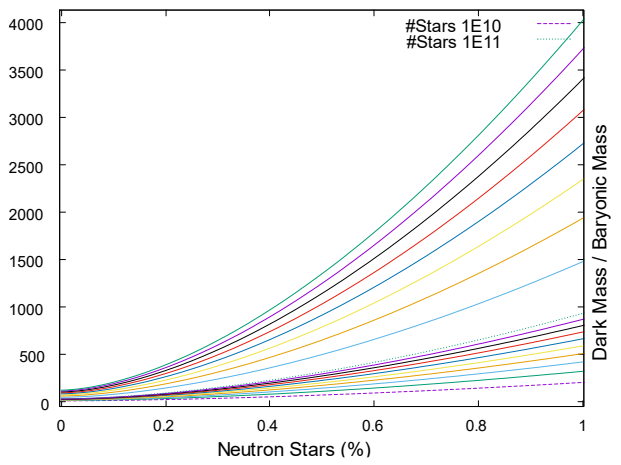


Figure 6

Black Holes

While the potential energy of stellar systems can now be calculated using the methods previously described, extending the analysis to black holes introduces additional complexity. However, the underlying principle remains the same: the gravitational potential energy is defined as the difference between two configurations—the initial, merged black hole and the final state consisting of two separate black holes.

As in the stellar case, the relevant quantity is the energy released during merger. According to general relativity (see references (15; 16)), the final mass M_f of the remnant black hole is given by:

$$M_f = (M + m)(1 - \eta \epsilon_{\text{rad}}(q, \chi))$$

where:

$$\eta = \frac{Mm}{(M + m)^2}, \quad q = \frac{m}{M} \leq 1, \quad \chi = \text{spin parameter}$$

and $\epsilon_{\text{rad}}(q, \chi)$ is an empirical function representing the efficiency of gravitational wave emission. For non-spinning black holes, a widely used approximation is:

$$\epsilon_{\text{rad}}(q) = 0.048 \cdot \frac{(1 - q)^2}{(1 + q)^4}$$

The effective spin of the remnant is given by:

$$\chi_{\text{eff}} = \frac{M\chi_M + m\chi_m}{M + m}$$

If the component black holes have randomly oriented spins, the net spin of the merged object will, on average, tend toward zero, i.e. $\chi_{\text{eff}} \approx 0$.

Let $f(m, M) \mapsto M_f$ be a function representing the final mass of a black hole formed by merging two black holes of masses m and M , according to the relativistic formula described earlier.

To extend this to the merger of n black holes, we reduce the problem to the case of n identical black holes, each of mass m , so that the total initial mass is $M = nm$. Define a recursive function $g(n, m) \mapsto M_f$ where $g(n, m)$ returns the final mass after all n black holes of mass m are merged. The function proceeds as follows:

```

g(n, m)  $\mapsto$   $M_f$ :

    IF n = 2 THEN  $M_f = f(m, m)$ 
    IF n = 3 THEN  $M_f = f(m, m), M_f = f(m, M_f),$ 
    ELSE
        isOdd = false
        IF Odd(n) THEN n = n - 1, isOdd = true
         $M_f = g(n/2, m)$ 
         $M_f = f(M_f, M_f)$ 
        IF isOdd THEN  $M_f = f(m, M_f)$ 

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This recursive structure is well-suited for dynamic programming. By storing intermediate results of $g(a, b)$ in a memoization table, the overall number of distinct computations is reduced significantly, achieving logarithmic efficiency in the number of mergers, i.e. $O(\log^2 n)$.

Negligible Potential Energy Contribution from Black Hole Mergers

The gravitational potential energy produced by black hole mergers is extremely limited, as it is drawn directly from the initial mass of the black holes. By definition, the radiated energy—primarily in the form of gravitational waves—can never exceed the total mass-energy of the system.

Even under ideal conditions, such as an aligned prograde spin configuration (which maximizes radiative efficiency), the fraction of mass converted into gravitational wave energy remains low. Using a semi-analytical approximation derived from numerical simulations (17; 18), the radiated energy fraction is given by:

$$\frac{E_{\text{gw}}}{M} \approx \eta(1 - 4\eta)[1 - 0.0686(1 - \chi_{\text{eff}})^2]$$

where $\eta = \frac{Mm}{(M+m)^2}$ is the symmetric mass ratio, and χ_{eff} is the effective dimensionless spin.

In the most favorable scenario—mergers involving black holes exceeding 500 solar masses and spins near unity—the maximum energy fraction remains below 16%.

Given both this upper bound and the relatively low abundance of black holes in most galactic environments, their contribution by fusion to the total gravitational potential energy of celestial systems is negligible.

Black Holes and the Problem of Volume Conservation

Despite their relatively small contribution to gravitational wave energy, it is impossible to exclude black holes from the calculation of gravitational potential energy. According to Newtonian gravity, any massive object—including black holes—contributes to the total potential energy via the standard expression $E_p = -GmM/d$.

Consider the construction of the initial compact configuration, which conserves the total volume of gas and stars and results in a body of mass M and volume V . If the system also contains n black holes of mass m and individual volume v , then—by analogy—they should be included in the compact configuration as well, leading to a revised mass and volume: $M' = M + nm$, $V' = V + nv$.

However, this treatment presents a paradox. If black holes are introduced into the compact configuration, they would, under gravitational attraction, eventually coalesce. But unlike normal matter, *the merger of black holes reduces their total volume*—a property not shared by stars or gas. Thus, treating them analogously within a volume-conserving framework breaks the model's fundamental assumption of conservation.

This highlights a unique property of black holes: they violate volume conservation through merger. Their inclusion in gravitational potential energy calculations is necessary, but their geometric behavior introduces a conceptual inconsistency that challenges the volume-conserving assumption used to construct the compact configuration.

Black Holes and the Limiting Case of Gravitational Potential Energy

This analysis suggests that black holes must be treated analogously to ordinary matter within the volume-conserving compact configuration—but with the important caveat that their unique volume behavior under merger must be respected. In this framework, we merge all black holes present in the system into a single equivalent black hole.

The resulting object may be highly dense (e.g., density ≈ 1000 for a mass of $4.3 \times 10^6 M_\odot$) or extremely diffuse on a geometric scale (e.g., density ≈ 0.001 for a mass of $4.3 \times 10^9 M_\odot$) due to the non-linear scaling of Schwarzschild radius with mass. Paradoxically, the presence of a larger number of black holes—when merged into a single object—can decrease the total gravitational potential energy of the system.

This outcome is unavoidable. As the number and total mass of black holes increase, the system approaches the physical limit of forming a single, final black hole. According to the no-hair theorem, such a black hole has no internal structure—no mass distribution, no gravitational self-interaction, and thus no gravitational potential energy in the classical sense. It is entirely defined by its mass, spin, and charge. Consequently, black holes represent the endpoint of gravitational compaction, beyond which no further gravitational binding energy can be extracted. As their dominance in a system grows, the accessible gravitational potential energy inevitably declines.

Model Results with Varying Black Hole Fractions

Figures 7 through 10 present the results of our theoretical model for galaxies containing the same number of stars as previously considered, now fixed to solar mass and radius: $1 M_{\odot}$, $1 R_{\odot}$. The x-axis shows the fraction of the total stellar mass that is assigned to black holes, expressed as a percentage.

At a black hole mass fraction of approximately 1.3%, the dark mass ratios for galaxies with 100, 200, and 300 billion stars converge to values of 19.7, 22, and 23, respectively. These results align well with observational constraints—particularly the commonly cited 95% dark matter proportion for the Milky Way, which corresponds to a dark mass ratio of ~ 20 , and estimates placing the stellar mass fraction in black holes at around 1% (19).

The trend of decreasing dark mass production with increasing black hole proportion persists up to galaxy sizes of approximately one billion stars, beyond which the effect begins to reverse slightly. Remarkably, with only a 5% black hole mass fraction, most galaxies become indistinguishable from one another in terms of their dark mass ratio. At this point, only the smallest galaxies remain distinct, while larger systems converge toward a universal behavior in dark mass production.

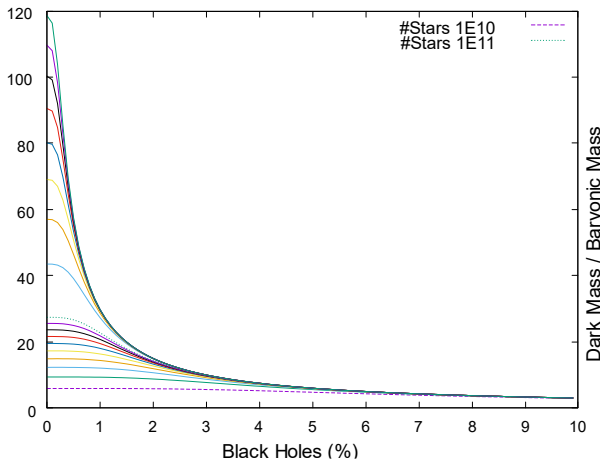


Figure 7

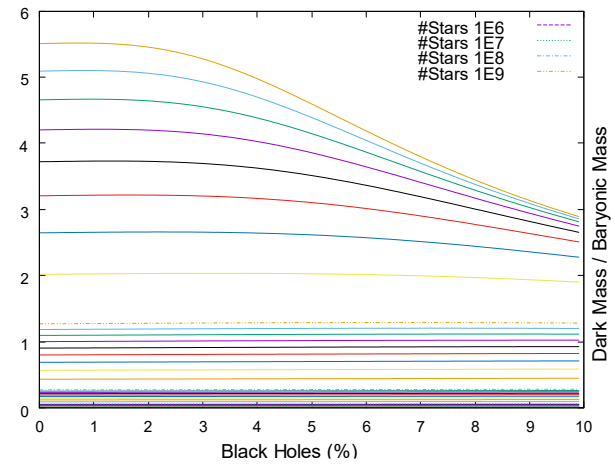


Figure 8

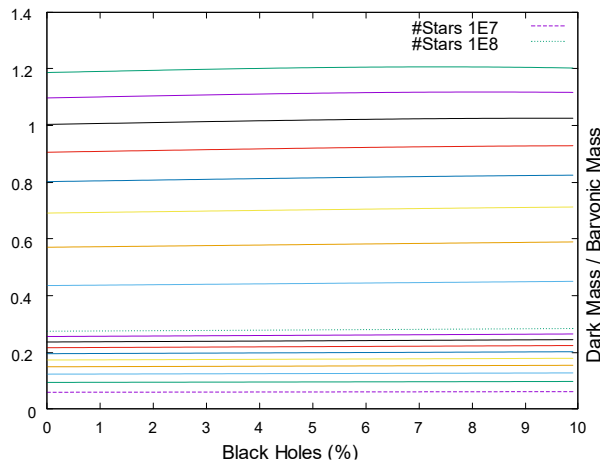


Figure 9

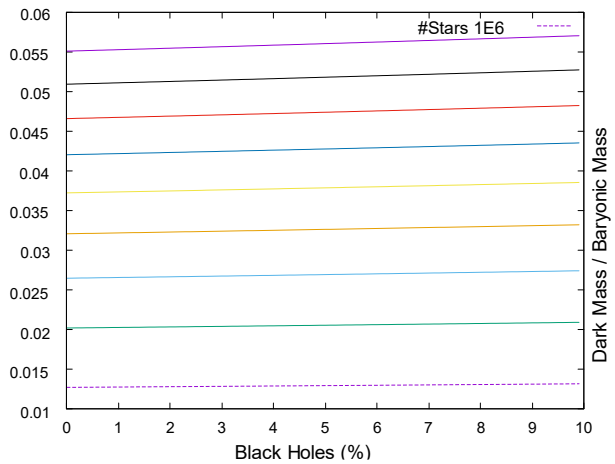


Figure 10

Galactic Models

To evaluate whether gravitational potential energy alone can account for the observed dark mass in galaxies, we tested our theoretical model against the SPARC sample (20), which includes rotation curve data for 175 galaxies along with their mass models.

For each data point, the dataset provides the observed velocity V_{obs} and the contribution of the baryonic mass contained in the orbital radius V_{bar} as well as its division in bulge stars V_{bul} , disk stars V_{disk} and gas V_{gas} . These values are used to estimate the mass contained in the orbital radius M_{bul} , M_{disk} , M_{gas} using the mass models. Subsequently, the required dark mass M_{dark} is calculated by comparing the observed acceleration to that predicted by visible matter. Each point also includes an estimate of the uncertainty eM_{dark} for the inferred dark mass.

Two galaxies (UGC 4305 and UGC 6628) were excluded due to negative global dark mass, leaving a final sample of 173 galaxies. This dataset contains 3,362 kinematic data points, of which 3,039 (approximately 90%) were retained for analysis. Data points with negative dark mass values or negative errors were removed to ensure consistency and reliability.

Stellar Population Fitting Method

For each data point in the SPARC sample, we apply our model to determine the optimal composition of stellar remnants and evolved stars—specifically black holes, white dwarfs ($0.6 M_{\odot}$, $0.0085 R_{\odot}$), red dwarfs ($0.4 M_{\odot}$, $0.5 R_{\odot}$), and red giants ($4 M_{\odot}$, $100 R_{\odot}$)—that can reproduce the required dark mass M_{dark} as closely as possible.

For each configuration, we compute the gravitational potential energy E_{dark} predicted by the model and evaluate the absolute deviation: $\Delta_{dark} = |E_{dark} - M_{dark}|$

The fitting process aims first to bring Δ_{dark} within the observational uncertainty: $\Delta_{dark} < eM_{dark}$ or $\Delta_{dark} < 2eM_{dark}$

While a more precise approach would account for both the uncertainty in the observed dark mass M_{dark} and the uncertainty in the model prediction E_{dark} , we treat eM_{dark} as the effective error margin for E_{dark} as a practical simplification.

If no composition satisfies this criterion, we minimize Δ_{dark} as a second step. Each data point is treated independently.

Model Fit Results

The model demonstrates excellent agreement with observational data. Among the 173 galaxies analyzed, 86% (149 galaxies) exhibit no residual errors across all included data points when evaluated using the relaxed threshold $\Delta_{dark} < 2eM_{dark}$. This corresponds to 81% of the total data points (2,467 out of 3,039) meeting the same condition.

Under the stricter criterion $\Delta_{dark} < eM_{dark}$, 74% of galaxies (128 out of 173) are fully resolved with no individual data point exceeding the error threshold, accounting for 71% of the retained data points.

Overall, 95% of all individual data points meet the $2eM_{dark}$ condition, and 92% fall within the stricter $1eM_{dark}$ threshold. These results confirm the robustness of the model in reproducing observed dark mass distributions across a wide range of galaxy types.

Moreover, when examining the positioning of the predicted value E_{dark} within the error bar of M_{dark} —where 0% corresponds to an exact match and 100% lies at the edge of the error bar—the results are highly concentrated near the center. For the 86% of galaxies that satisfy $\Delta_{dark} < 2eM_{dark}$, the average offset is $3\% \pm 9\%$. For the 74% of galaxies with $\Delta_{dark} < eM_{dark}$, the offset improves to $1.4\% \pm 1.6\%$.

All error margins ($\pm e$) in this analysis correspond to one standard deviation.

Population Composition in Successful Fits

One of the most compelling outcomes of the model is the physical plausibility of the stellar compositions required to reproduce the observed dark mass. For the 86% of galaxies where the model fit satisfies $\Delta_{dark} < 2eM_{dark}$, the algorithm converges to the following average stellar population composition (expressed as a fraction of the total baryonic mass):

Black holes:	$5.4\% \pm 4.6\%$
White dwarfs:	$30\% \pm 7.7\%$
Red giants:	$20\% \pm 6.7\%$
Red dwarfs:	$45\% \pm 7.7\%$

These proportions remain consistent across galaxies meeting the stricter criterion $\Delta_{\text{dark}} < eM_{\text{dark}}$, as well as those with slightly larger residuals. This consistency supports the conclusion that gravitational potential energy can account for the observed dark mass using realistic distributions of stellar types.

The relatively high fraction of white dwarfs inferred by the algorithm might be moderated by including a small proportion of neutron stars, which offer high density at lower mass. Similarly, the contribution from red giants may be partially substituted by rarer but very low density and more massive hypergiants. These adjustments would preserve the energetic balance while remaining astrophysically plausible.

Intra-Galactic Consistency and Stellar Distribution Profiles

The previous model demonstrates that it is possible to reproduce the observed dark mass entirely through gravitational potential energy generated by common stellar populations—namely black holes, white dwarfs, red dwarfs, and red giants. However, this approach optimizes each data point independently and does not enforce consistency across points within the same galaxy.

To introduce intragalactic coherence, we compute m_{bul} , m_{disk} , m_{gas} and m_{dark} , which represent the masses enclosed within the orbital slice corresponding to each data point. The proportions of stellar types and black holes assigned to a given point can only depend on the total local baryonic mass, defined as $m_{\text{bul}} + m_{\text{disk}} + m_{\text{gas}}$.

This constraint reduces the available degrees of freedom, making the fitting process more complex, but it ensures global physical consistency across the galaxy.

Under this constraint, the model achieves successful fits—with no residual errors—within $\Delta_{\text{dark}} < 2eM_{\text{dark}}$ for 62% of the galaxies (108 out of 173), covering 49% of the total data points. While this represents a reduction compared to the unconstrained model, it provides a key physical advantage: it enables the generation of internally consistent stellar composition profiles for each galaxy.

Figures 11 and 12 present two example galaxies illustrating the recovered stellar distribution across orbital radii, as predicted by this constrained optimization.

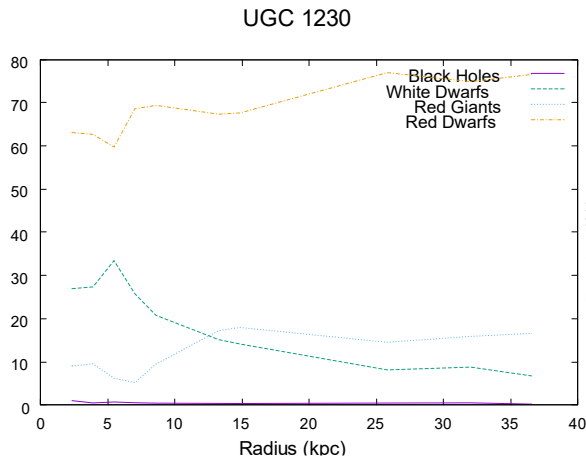


Figure 11

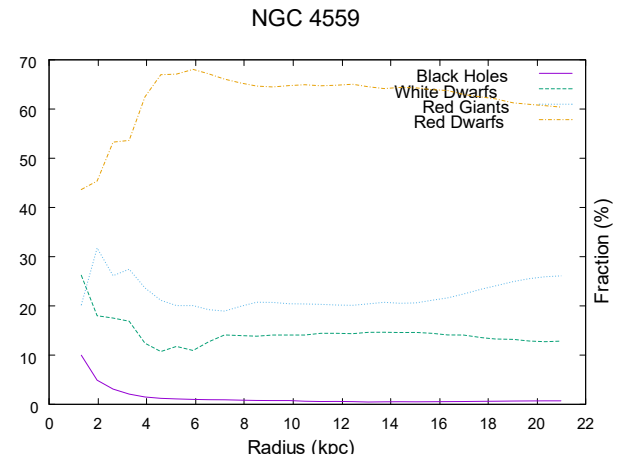


Figure 12

The Mass Energy Field

To understand how gravitational potential energy—interpreted here as dark mass—affects galactic rotation curves, we must first consider the spatial distribution of this mass-energy. Unlike baryonic matter, this invisible component is not localized in discrete particles, but rather is distributed throughout the gravitational field itself.

The persistent invisibility and intangibility of this form of mass-energy presents a conceptual challenge. However, this difficulty is not entirely new. As noted by Leon Brillouin (21; 22):

“All energy has mass, but it seems that the case of potential energy has been omitted. The founders of Relativity hardly mention it. In fact, the corresponding energy is spread throughout space, and its mass cannot be exactly localized. The symmetry of the distribution suggests dividing the mass between the various interacting particles. It is therefore necessary, from classical Relativity onwards, to revise the values of the masses. Well before quanta, *renormalization* was essential (and was omitted) in Einstein's Relativity.”

Brillouin's insight underscores a key theoretical omission in classical relativity: the mass-equivalence of potential energy is spatially distributed and cannot be attributed to any single point. This aligns with our interpretation: the gravitational potential energy of a system manifests as an extended, nonlocal mass-energy field. As such, the field itself contributes to the curvature of spacetime, producing the dynamical effects attributed to dark mass.

Field Proportionality and the Localization of Dark Mass

Although the problem of the mass associated with gravitational potential energy was clearly articulated by Brillouin in 1965, it has not received a satisfactory treatment in either classical or relativistic physics. In this work, we postulate that the mass stored in the gravitational field is directly proportional to the intensity of the field at a given point.

Let $\Phi(E_a)$ denote the gravitational field produced by a dark mass E_a at position a , and $\Phi(M_a)$ the field produced by a baryonic mass M_a at the same point. Likewise, at position b , let $\Phi(M_b)$ and $\Phi(E_b)$ denote the baryonic and dark field components, respectively. We assume the following proportionality relationship:

$$\frac{\Phi(E_b)}{\Phi(E_a)} = \frac{\Phi(M_b)}{\Phi(M_a)}$$

Using the standard expression for the gravitational field produced by a point mass,

$$\Phi(M) \propto \frac{M}{R}$$

where R is the characteristic radius or distance from the source, we obtain:

$$\frac{\Phi(M_b)}{\Phi(M_a)} = \frac{M_b R_a}{M_a R_b} \quad \text{and} \quad \frac{\Phi(E_b)}{\Phi(E_a)} = \frac{E_b R_a}{E_a R_b}$$

Equating the two ratios leads to:

$$\frac{E_b R_a}{E_a R_b} = \frac{M_b R_a}{M_a R_b}$$

which simplifies to:

$$\frac{E_b}{E_a} = \frac{M_b}{M_a} \quad \Rightarrow \quad E_b = E_a \cdot \frac{M_b}{M_a}$$

This relation provides a simple but powerful rule: the local dark mass is proportional to the local baryonic mass, under the assumption that gravitational field strength governs the distribution of field-stored mass-energy.

Component Contributions to Dark Mass Production

In the galactic context, the production of gravitationally induced dark mass must be analyzed component by component, as each baryonic contributor exhibits a different efficiency in generating potential energy. In practice, the available observational data typically separate the baryonic mass into at least two components: the gas mass and the stellar disc mass.

If the gas component contributes a dark mass E_{gas} from a baryonic mass M_{gas} , and the stellar disc contributes E_{disk} from a baryonic mass M_{disk} it is necessary to compute their relative contributions to the total gravitationally induced dark mass. These can be expressed as:

$$e_{\text{gas}} = \frac{E_{\text{gas}} m_{\text{gas}}}{M_{\text{gas}}}, \quad e_{\text{disk}} = \frac{E_{\text{disk}} m_{\text{disk}}}{M_{\text{disk}}}$$

The total contribution is then: $e_{\text{tot}} = e_{\text{gas}} + e_{\text{disk}}$

However, this division is still simplistic. As shown earlier, the gravitational potential energy varies significantly depending on stellar type (e.g., red dwarfs, red giants, white dwarfs, black holes). A more accurate treatment would require dividing the baryonic mass into multiple categories, each with distinct mass and radius parameters. These variations determine the gravitational binding energy and thus the amount of dark mass effectively “generated” by each category. A coarse two-component model can serve as an approximation, but more detailed modeling is essential for precision.

Two-Component Model with a Single Fitting Parameter

In the absence of detailed information about stellar sub-populations, we restrict the analysis to the two primary baryonic components typically available in observational data: gas and the stellar disc. Their respective contributions to the gravitationally induced dark mass are modeled using a single free parameter f , which represents the fraction of E_{dark} attributed to the gas component.

The total modeled dark mass is expressed as:

$$E_{\text{dark}} = E_{\text{gas}} + E_{\text{disk}} = f E_{\text{dark}} + (1 - f) E_{\text{dark}}$$

From this, we define the relative contributions for each data point as:

$$e_{\text{gas}} = \frac{f E_{\text{dark}} m_{\text{gas}}}{M_{\text{gas}}}, \quad e_{\text{disk}} = \frac{(1 - f) E_{\text{dark}} m_{\text{disk}}}{M_{\text{disk}}}$$

where:

- E_{dark} , M_{gas} , and M_{disk} are values measured at the midpoint radius of the galaxy.
- m_{gas} and m_{disk} are the local baryonic mass values at the data point.

The parameter f is optimized once per galaxy and applied uniformly across all data points. Despite its simplicity, this model reproduces the observed dark mass within the margin of error for 76 of the 175 galaxies in the SPARC sample—corresponding to 43% of the dataset.

To remain conservative in our error analysis, we assign to each calculated value e_{tot} the same uncertainty as that of the corresponding observed dark mass p_{dark} .

Figures 13 and 14 illustrate representative examples of galaxies well fitted by this single-parameter, two-component model.

Incremental Refinement of the Two-Component Model

It is possible to increase the accuracy of the model by introducing additional f parameters, each associated with a distinct radial region of the galaxy. These parameters allow the relative weighting between gas and stellar disc contributions to vary spatially, thereby improving the fit to the observed dark mass profile.

For each segment, the scaling is performed using the midpoint of the segment as the local reference, with its own values of E_{dark} , M_{gas} , and M_{disk} . This ensures that the proportional calculation within each region remains physically consistent and properly normalized.

Using two f parameters, we achieve error-free fits for 29 additional galaxies. Allowing three f parameters yields an additional 6 galaxies with perfect fits. In total, this approach provides error-free solutions for 111 out of 175 galaxies, or approximately 63% of the SPARC sample.

The ability to reproduce the observed dark mass distributions across such a large fraction of galaxies—using only local baryonic mass and gravitational potential energy scaling—strongly supports the existence of a direct, physically meaningful relationship between baryonic mass and the inferred dark mass. This reinforces the core hypothesis that gravitational potential energy, when correctly modeled, accounts for the so-called dark matter without requiring matter.

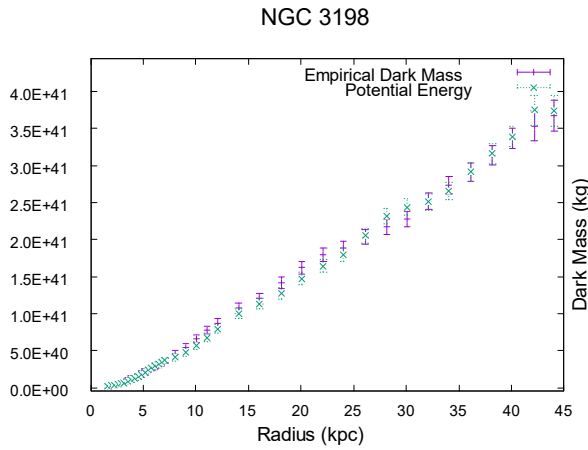


Figure 13

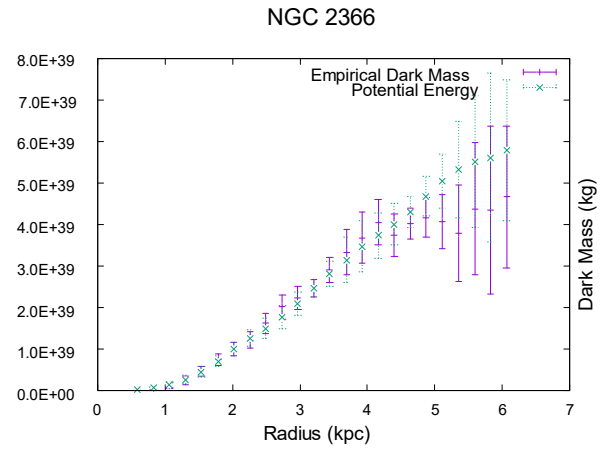


Figure 14

Dark Energy

In previous models for calculating the gravitational potential energy of self-gravitating systems (galaxies, clusters, etc.), the large-scale direct interaction term between distant bodies was considered negligible. This term, of the form $-GmM/d$, where d can be interpreted as the mean separation between bodies, is indeed extremely small on stellar scales. However, at cosmological scales, and when summed over all bodies in the universe, its contribution becomes significant.

Let us consider a set of n point masses of mass m , separated on average by a distance d , for a total mass M . The sum of interactions between each mass and the cumulative preceding ones yields a potential energy:

$$\Delta E_{\text{dil}} = -\frac{3Gm^2}{5d} \sum_{i=1}^n (i-1) = -\frac{3Gm^2}{5d} \cdot \frac{n(n-1)}{2} \approx -\frac{3GM^2}{10d} \quad (\text{for } n \gg 1)$$

This energy describes the “diluted” state of the universe: a configuration of weakly bound masses at large distances.

Simplification of ΔE_p

$$\Delta E_p = \frac{3G}{5} \left(\frac{M_n^2}{R_n} - n \frac{m^2}{r} \right)$$

Assuming $M_n = nm$, we have:

$$M_n^2 = n^2 m^2 \quad \text{and} \quad R_n = \left(\frac{M_n r^3}{m} \right)^{1/3} = n^{1/3} r$$

Substituting into the equation:

$$\Delta E_p = \frac{3G}{5} \left(\frac{n^2 m^2}{n^{1/3} r} - \frac{nm^2}{r} \right) = \frac{3G}{5} (n^{5/3} - n) \frac{m^2}{r}$$

The asymptotic behavior shows that for $n \gg 1$, the $n^{5/3}$ term dominates over n , justifying the approximation:

$$\Delta E_p \approx \frac{3G}{5} \cdot \frac{M_n^2}{R_n}$$

Gravitational Compaction Energy

In contrast, if the same total mass M is compacted into a homogeneous solid sphere of radius R (e.g., preserving solar density), the gravitational potential energy becomes:

$$\Delta E_{\text{comp}} = -\frac{3GM^2}{5R}$$

The energy difference between these two configurations is:

$$\Delta E_p = \Delta E_{\text{dil}} - \Delta E_{\text{comp}} = \frac{3GM^2}{5} \left(\frac{1}{R} - \frac{1}{2d} \right)$$

This expression is always positive for $R \ll 2d$, which is indeed the case in the real universe: $R \sim 10^{16}$ m (solar density), $d \sim 10^{26}$ m (average intergalactic separation or corrected Hubble radius), $R/d \sim 10^{-10}$, making the effective contribution significant.

A crucial observation is that R is fixed (set by the compactness and chosen density), while d increases with cosmic expansion. Therefore, the effective gravitational potential energy ΔE_p increases with d (and thus with volume V):

$$\frac{d\Delta E_p}{dV} > 0 \quad \Rightarrow \quad P = -\frac{d\Delta E_p}{dV} < 0$$

This result is interpreted as a negative gravitational pressure, not originating from an exotic scalar field, but from the geometric evolution of a structured gravitational field. As space expands, background gravitational interactions weaken, leading to a growing potential energy. Its negative derivative acts as a repulsive pressure — a natural source of accelerated expansion.

Application to the Observable Universe at Galactic Density

Instead of compacting the total mass of the universe to stellar (solar) density, we now assume the universe is structured in galactic halos containing both baryonic and dark matter. These halos are modeled as oblate ellipsoids with a height equal to one-third their radius.

For a typical compatible galaxy (SPARC, EAGLE, Illustris):

- Average mass (baryons + dark matter): $M_{\text{gal}} \in [1.6, 3.3] \times 10^{42}$ kg
- Average halo radius: $R_{\text{halo}} \in [2.4, 3.7] \times 10^{21}$ m

Using mass (baryonic + dark matter) of the observable universe: $M_{\text{gal}} \approx 8.8 \times 10^{52}$ kg (according to Λ CDM)

$$R_n = \left(\frac{3M_{\text{tot}}}{4\pi\rho_{\text{gal}}} \right)^{1/3} = \left(\frac{3NM_{\text{gal}}}{4\pi M_{\text{gal}}/V_{\text{gal}}} \right)^{1/3} = \left(\frac{3NV_{\text{gal}}}{4\pi} \right)^{1/3}$$

With:

$$V_{\text{gal}} = \frac{4}{9} \pi R_{\text{halo}}^3 \Rightarrow R_n = \left(\frac{3N \cdot \frac{4}{9} \pi R_{\text{halo}}^3}{4\pi} \right)^{1/3} = \left(\frac{N}{3} \cdot R_{\text{halo}}^3 \right)^{1/3} = \left(\frac{N}{3} \right)^{1/3}$$

Compacting the total mass of the observable universe at this density yields: $N \in [2.7 \times 10^{10}, 5.5 \times 10^{10}]$ (compatible with Deep Field and JWST).

$$\Delta E_p = \frac{3GM^2}{5} \left(\frac{1}{R_n} - \frac{1}{2d} \right)$$

Using $d = kc/H \approx 4.5 \times 10^{26}$ m with $k = 3.27$ (comoving radius in Λ CDM), we obtain numerically:

$$\Delta E_p \in [6.6 \times 10^{70}, 2.6 \times 10^{71}] \text{ J}$$

This result matches the estimated range for the total dark energy of the observable universe, as inferred from Planck 2018 + WMAP + SNIa + BAO data:

$$E_\Lambda \in [6.6 \times 10^{70}, 2.6 \times 10^{71}] \text{ J}$$

It follows that the macroscopic gravitational potential energy, derived from galactic-scale structure, may provide a purely geometric and classical interpretation of dark energy — requiring neither exotic fields nor arbitrary cosmological constants.

Emergent Pressure from Gravitational Potential

The effective gravitational potential energy of the structured system, previously written as:

$$\Delta E_p = \frac{3GM^2}{5} \left(\frac{1}{R} - \frac{1}{2d} \right)$$

can be expressed as a function of cosmic volume V , noting that:

- R , is constant (galactic structure is assumed stable),
- $d \propto V^{1/3}$

Then:

$$\Delta E_p(V) = \frac{3GM^2}{5} \left(\frac{1}{R} - \frac{1}{2V^{1/3}} \right)$$

The corresponding effective pressure is:

$$P(V) = -\frac{d\Delta E_p}{dV} = -\frac{3GM^2}{5} \cdot \frac{1}{2} \cdot \frac{1}{3} V^{-4/3} = -\frac{GM^2}{10V^{4/3}}$$

This pressure is strictly negative and decreases with expansion. It does not stem from exotic substances or hypothetical fields, but arises from the weakening of gravitational interaction as volume increases. This macroscopic potential energy, generated by gravitational structuring, stores information in the field itself. Its derivative manifests as a repulsive effect — a geometric mechanism that can mimic dark energy.

Integration into General Relativity

In general relativity, the source of spacetime curvature is not mass alone, but the energy–momentum tensor $T_{\mu\nu}$. For a homogeneous and isotropic universe, this tensor reduces to that of a perfect fluid:

$$T^{\mu\nu} = \left(\rho + \frac{P}{c^2} \right) u^\mu u^\nu + P g^{\mu\nu}$$

The pressure obtained above is not a hypothesis but a geometric consequence of the gravitational field's structure. It is thus natural to consider this pressure — arising from effective potential energy — as contributing to a gravitational component of the energy–momentum tensor:

$$T_{\text{grav}}^{\mu\nu} = \left(\rho_{\text{pot}} + \frac{P_{\text{grav}}}{c^2} \right) u^\mu u^\nu + P_{\text{grav}} g^{\mu\nu}$$

where:

$$\rho_{\text{pot}} = \frac{\Delta E_p}{V}, \quad P_{\text{grav}} = -\frac{GM^2}{10V^{4/3}}$$

This tensor has the same structure as that of dark energy, but its origin is not constant — it emerges dynamically from the gravitational structuring of the universe.

Reproducing Λ CDM Cosmology

To assess the cosmological implications of this effective pressure, we substitute it into the Friedmann equations. Assuming a spatially flat universe ($k = 0$):

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} \left(\rho + \frac{3P}{c^2} \right)$$

With:

$$\rho = \frac{M}{V}, \quad P = -\frac{GM^2}{10V^{4/3}}$$

We obtain the dynamical equation for cosmic volume:

$$\frac{\dot{V}}{V} - \frac{1}{3} \left(\frac{\dot{V}}{V} \right)^2 = -\frac{A}{V} + \frac{B}{V^{4/3}}$$

with $A \propto GM$ and $B \propto G^2 M^2 / c^2$. This nonlinear equation admits two asymptotic regimes:

At early times (matter-dominated), gravity dominates, and the universe expands as:

$$V(t) \propto t^2 \quad \Rightarrow \quad a(t) \propto t^{2/3}$$

At late times, gravitational pressure dominates, and the expansion becomes exponential:

$$V(t) \propto e^{3Ht} \quad \Rightarrow \quad a(t) \propto e^{Ht}$$

This reproduces the behavior of Λ CDM cosmology, without invoking a cosmological constant, but by deriving it from an effective geometric pressure rooted in the structured gravitational field.

Conclusion

The objective of this article was to demonstrate that the phenomenon commonly referred to as “dark matter” can be fully accounted for by gravitational potential energy. Through both theoretical and empirical approaches, we have shown that the potential energy generated by gravitational interactions—particularly when applied to systems like the Earth–Moon or Earth–Sun—is far from negligible and must be considered as a real, physical quantity.

Using simplified theoretical models, we established that this potential energy, when converted via the mass–energy equivalence relation $E = mc^2$, yields contributions on the same order of magnitude as the dark mass inferred from galactic rotation curves.

This leads to a critical implication: if gravitational potential energy exists—as is well established—and if the mass–energy equivalence is universally valid, then the mass associated with that energy must manifest gravitational effects. In other words, if potential energy is physically real, then dark mass must exist as a consequence of it.

Therefore, the observed effects attributed to dark matter may not require the existence of new particles or exotic matter. Instead, they can be understood as arising naturally from the structure of the gravitational field itself. The presence of sufficient gravitational potential energy to account for galactic dynamics constitutes a sufficient physical condition for the phenomenon currently labeled as “dark matter.” The mass–energy relation demands that this energy contributes to the curvature of space-time, and thus to gravitational phenomena at galactic and cosmological scales.

We subsequently compared our theoretical model to a substantial body of observational data and obtained an excellent fit, despite the necessarily simplified assumptions used throughout. This empirical success provides a necessary condition for validating the model: the gravitational potential energy predicted by the distribution of baryonic mass must be sufficient to account for the observed dark mass. Our results show that this condition is fulfilled across a wide range of galaxy types.

We then extended our model by applying the principle of causal proportionality—namely, that identical mass distributions of the same type must generate proportionally identical effects. If a baryonic mass M , composed of

stars of a given type (e.g., solar-mass stars), produces a corresponding dark mass E , then a subcomponent of mass m will generate a dark mass:

$$e = \frac{Em}{M}$$

When expressed as a gravitational field, this implies that the spatial distribution of dark mass follows the distribution of baryonic mass—provided that the composition (stellar types, gas fraction) remains homogeneous. This directly recovers the empirically observed correlation between baryonic and dark mass profiles (23; 24), but here it arises naturally as a consequence of the gravitational field's structure, rather than requiring a phenomenological assumption.

Thus, the hypothesis that dark mass is generated by gravitational potential energy not only accounts for its magnitude but also provides a fundamental explanation for its spatial distribution. What had previously appeared mysterious—the alignment between baryonic and dark matter—is shown here to be the logical consequence of field-based mass-energy equivalence.

Finally, this same gravitational potential energy offers a natural explanation for the emergence of negative pressure on cosmological scales. As gravitational interactions become progressively diluted with the increasing separation of structures, the macroscopic binding energy decreases—corresponding thermodynamically to an effective negative pressure. This pressure, derived directly from the evolution of the structured gravitational field, drives an accelerated expansion consistent with observations of dark energy. Thus, gravitational potential energy not only accounts for the dark mass, but can also generate the repulsive effect that governs the large-scale dynamics of the universe.

The Problem of the Epistemological Obstacle

Perhaps the most profound question raised by this work is why the role of gravitational potential energy in generating dark mass has remained unrecognized for so long. Both gravitational potential energy and the mass-energy equivalence relation $E = mc^2$ are well-established principles, each supported by extensive experimental validation. Moreover, the concept of binding energy contributing to mass is not new: nuclear binding energy, for example, is well understood as a measurable mass defect in atomic systems, though its localization is typically attributed to the quantum field.

The primary reason gravitational binding energy has been historically dismissed as a source of mass is the perception that gravity is too weak a force to produce significant energy effects. However, our results challenge this assumption. While nuclear binding energy yields modest contributions to mass on the scale of atomic nuclei—often a few percent—gravitational potential energy can generate dark mass equivalents many times greater. In galactic systems, the ratio of gravitational potential energy to baryonic mass can exceed a factor of 10, and in clusters or larger structures, this ratio can reach orders of magnitude higher.

This contrast highlights a scale-dependent truth: gravity is weak on the quantum scale but becomes dominant on astronomical scales, while the strong nuclear force behaves inversely. Immense at the subatomic level but negligible beyond the nucleus. The failure to account for this scale dependence may explain why the gravitational origin of dark mass has remained obscured despite the theoretical and empirical tools needed to uncover it.

Another major conceptual obstacle to recognizing the gravitational origin of dark mass lies in the widespread use of Newton's second theorem. Indeed, it allows a solid sphere to be considered as a point mass. This simplification naturally comes to mind when studying large stellar systems and has led to a poor understanding of what useful gravitational potential energy is. Traditional approaches focus on the relative position of celestial bodies, considered as point masses, emphasizing the term $E = -GmM/d$ as the main contribution. However, as we have demonstrated, this term is negligible compared to the energy associated with a restructuring of the system—for example, the merger or separation of massive bodies—under conservation of volume density.

The relevant potential energy does not result solely from distance, but from the energy difference between two distinct macroscopic states of the system. It is this transformation of state, respecting the physical properties of matter, that reveals the true energy content of the gravitational field. The historical focus on positional interaction, rather than on the internal reorganization of masses, has contributed to masking the role of potential energy as the real source of space-time curvature. By reconsidering potential energy as a quantity localized in the field and

generating mass, we find a physically coherent explanation of the dark mass phenomenon—already contained in classical gravitational theory, but which has remained unnoticed due to a discreet but decisive conceptual bias.

Challenge of General Relativity and Theoretical Implications

A deeper theoretical difficulty lies in the fact that Newtonian gravity (GN) is not the fundamental theory of gravitation—General Relativity (GR) is. Unlike GN, GR does not admit a general, global conservation law for energy. As Emmy Noether showed, conservation of energy in GR only holds under very specific symmetries, such as time-translation invariance, which do not apply in dynamically evolving spacetimes. In particular, the notion of gravitational potential energy, central to Newtonian dynamics, does not have a direct, covariant analog in GR.

This creates a conceptual tension: if a phenomenon as significant as the dark mass arises from gravitational potential energy, and this energy is absent or ill-defined in GR, then one must ask why Newtonian theory—a weak-field approximation—appears to succeed where the full relativistic theory offers no equivalent formulation. The challenge is philosophical as well as physical.

Nevertheless, GN remains extraordinarily accurate in the weak-field regime and is vastly simpler to work with than GR. The dynamics of galaxies, where the gravitational field is weak and velocities are non-relativistic, fall squarely into this domain. In Newtonian gravity, the potential energy of the system is known, but to determine the correct dynamics, that energy must be reintroduced into the system—that is, it must be treated as a source of gravity. It is difficult to believe that GR would yield significantly better results unless the issue lies in an incomplete accounting of gravitational self-induction—the field's interaction with its own energy density (25; 26; 27; 28). If this is indeed the case, then the most natural resolution would be to find a weak-field expansion of GR in which this self-induction appears explicitly. In such a formulation, the gravitational potential energy term should re-emerge, but now as a derived, not postulated, quantity.

More speculatively, the fact that gravity “recognizes” all forms of energy—including gravitational potential energy, long regarded as a mere computational artifact—suggests the existence of a deeper unifying principle. The Higgs field already hints at such a principle in a limited domain: it translates scalar potential energy into inertial mass for certain particles. Gravity appears to generalize this mechanism, coupling not just to scalar fields but to all forms of energy—kinetic, electromagnetic, nuclear, or gravitational.

Speculative Developments

Notably, no experiment has ever observed a violation of the equivalence between inertial and gravitational mass. This reinforces the idea that all energy, regardless of its origin, contributes equally to the gravitational field. In a quantum framework, this universality could point to a more fundamental origin: the graviton, rather than being an elementary particle, may be a Goldstone boson (29) emerging from the spontaneous breaking of a deeper spacetime symmetry (30; 31; 32; 33).

If spacetime consists of locally flat, discrete units—*atoms of geometry*—then curvature could arise from their relative orientations, much like angular defects in a crystal lattice. In this view, mass-energy would correspond to geometric misalignments among these fundamental building blocks of space-time. The resulting macroscopic curvature would be an emergent phenomenon: the manifestation of how energy redistributes itself through the relational structure of the underlying geometry.

That such an interpretation remains conceptually underappreciated may reflect not a lack of theoretical feasibility, but rather a persistent epistemic bias—a tendency to overlook well-established principles such as $E = mc^2$ and gravitational potential energy as the foundation for a radically simple explanation.

Although general relativity revolutionized our understanding of space and time, it does not exclude the existence of a flat, infinite, non-curved spacetime background—essentially, a Newtonian-like universe. Minkowski space is an exact solution of Einstein's equations in the absence of matter and energy, representing a limiting case in which the metric remains fixed, curvature vanishes, and time flows uniformly for all inertial observers.

In this context, one may speak of a “universe at rest”, by analogy with a mass at rest: a geometrically inert state in which both space and time are “at rest.” Such a spacetime, devoid of intrinsic curvature, could serve as a conceptual background on which physical processes unfold—even if general relativity, in its full formulation, neither presumes nor requires such a background.

This perspective opens the door to a compelling idea: perhaps geometry is not inherently dynamic but becomes so only when interactions—such as gravity—disturb this fundamental rest state. In this view, the Newtonian universe would act as an implicit absolute frame of geometric rest, and general relativity would describe the deviations from it induced by energy and matter. This reinterpretation gives conceptual legitimacy back to gravitational potential energy as a real, geometrically meaningful quantity—and lays the theoretical groundwork for explaining dark mass not through exotic particles or modified gravity, but through a classical mechanism rooted in the self-interaction of the gravitational field itself.

This raises a deeper implication. The very existence of gravitational potential energy—as a unified and measurable form of energy—presupposes the existence of a universal mechanism of energy storage at the most fundamental level. A compelling candidate is spacetime itself. If gravitational potential energy is real, then something must *store* that energy—and the only universally present structure is the spacetime manifold. This leads to a conceptual inversion of general relativity: rather than mass-energy curving spacetime, one could say that curvature itself gives rise to mass-energy. In this view, spacetime is not deformed by energy, but energy is the manifestation of spacetime’s internal geometric state. Gluons, despite being massless, are in fact the primary source of gravitational curvature in the universe.

This central role of gluons finds deep theoretical support in recent work (34) who showed that gravitational scattering amplitudes can be constructed as double copies of non-Abelian gauge amplitudes, specifically those of gluons in QCD. Applying this structure to realistic theories such as the Standard Model requires accounting for spontaneous symmetry breaking—it reveals a deeper principle: gauge interactions already contain the structural seeds of gravity. In this light, the gluon, carrier of the binding energy that gives rise to mass and the dominant contributor to gravitational curvature, becomes a compelling candidate for a fundamental “atom of spacetime.” It unifies dynamic energy and geometry, and its mathematical structure is naturally predisposed to generate gravitation. This correspondence strongly supports the view that gravity is not a separate field, but the emergent geometric form of energy dynamics itself—including that of gauge fields.

In this perspective, spacetime is not deformed by energy; rather, energy itself is a manifestation of the internal state of geometry. This interpretation naturally aligns with the equivalence principle, the universality of gravitational coupling, and even with global energy conservation—despite the presence of dynamic geometries. If correct, it implies that mass, inertia, and gravitation are not fundamental entities but emergent consequences of the geometric and connective structure of spacetime itself. The gravitational field would no longer be a mere force but an energetic reservoir—capable of storing and releasing mass through its own geometric configuration.

Pushed to its extreme, emergence dissolves even the concept of time from space. There are no clocks in the universe—only interactions between particles (35; 36). The notion that what we call “time” is merely the sequential ordering of quantum events becomes not only plausible but natural, especially if gravity itself arises from the web of fundamental interactions. If such approaches have so far succeeded only in recovering Newtonian gravity, it may be because gravitational potential energy has been neglected as a physically real quantity—particularly its capacity for self-induction. Gravitational potential energy could very well be the final missing ingredient in the search for a unified framework.

Final Remarks

In conclusion, the study of dark mass does not require an alternative theory of gravitation, exotic particles, or even General Relativity. The framework of classical Newtonian mechanics, when extended to include the mass–energy equivalence $E = mc^2$, proves sufficient. More importantly, we have demonstrated that simple, physically grounded models can accurately reproduce the magnitude and distribution of dark mass across a wide range of galactic systems. These models can be systematically improved—for example, by incorporating detailed stellar population synthesis, more accurate treatment of interstellar gas dynamics, or refined modeling of radial distributions. The same methodology may also be extended to larger cosmic structures such as galaxy clusters and superclusters.

One conclusion is unavoidable: gravitational potential energy can no longer be neglected in discussions of galactic dynamics and the dark matter problem. Its inclusion provides both a quantitative and conceptual resolution to a long-standing astrophysical mystery. The C++ program used to perform all numerical calculations and generate the corresponding figures is freely available at: dark-mass-generator.sourceforge.io.

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