

# An atypical case of round-off in *Mathematica*<sup>®</sup>

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## Abstract

A *Mathematica* round-off case generated by a real function of a real variable, not elementary expressible.

## 1 Introduction

In semiconductor physics, the calculation of the *chemical potential*  $\mu(T)$  [1], as a function of the thermodynamic equilibrium temperature  $T$ , is fundamental. This quantity solves a transcendent functional equation

$$F(T, \mu) = 0 \quad (1)$$

which can be made algebraic with the change of variable  $z = k_B T \ln \mu$ , where  $k_B$  is the Boltzmann constant, and the new quantity  $z$  is the *fugacity*. So

$$G(T, z) = 0 \quad (2)$$

Despite its algebraic character, (2) cannot be solved in closed form; solving numerically with *Mathematica* [2] we obtain the graph in Figure 1. Violent oscillations can be rarefied but not damped, using the `MaxRecursion` instruction as shown in Figure 1, while in Figure 3 we have reduced the thermal range to  $[0, 10]$ . From this last graph we see that the *Mathematica* kernel is unable to graph  $\mu(T)$  for  $T < 2$ . However, it seems to be

$$\lim_{T \rightarrow 0} \mu(T) = 0$$

On the contrary, it is well known that the aforementioned limit is  $< 0$ .

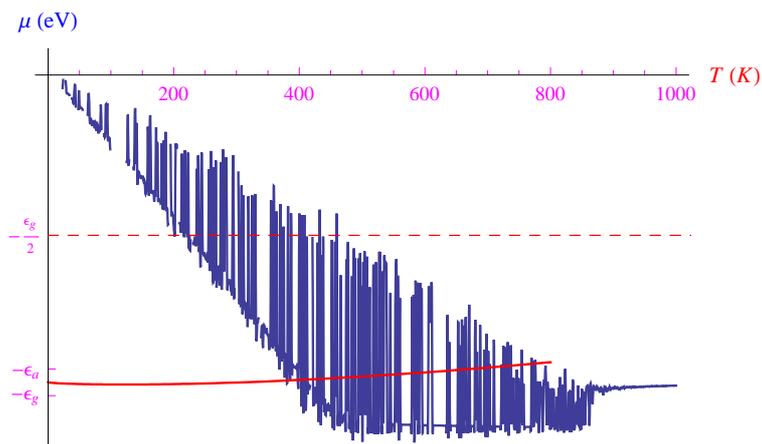


Figure 1: The curve in red is the solution obtained in a right neighborhood of  $T = 0$ . The rapid swings are due to the *Mathematica* round-off.

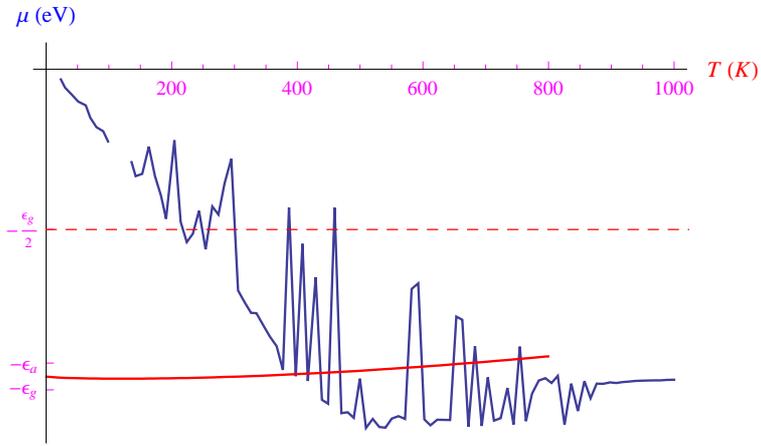


Figure 2: Reduction of oscillations by placing MaxRecursion  $\rightarrow 1$ .

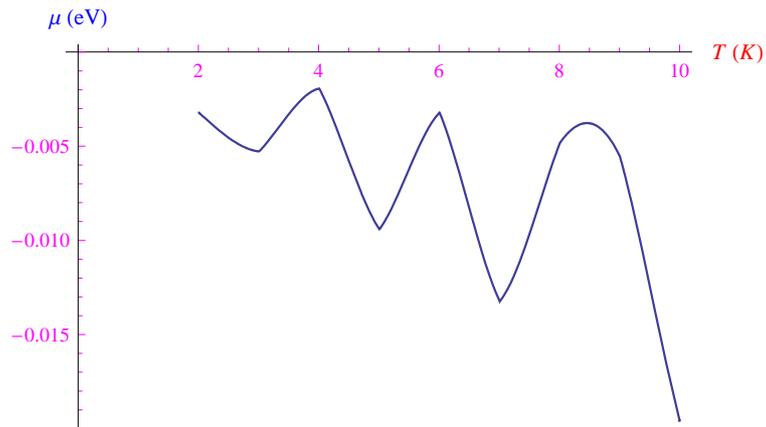


Figure 3: Trend for  $T \in [0, 10]$ .

## 2 Change of variable. Setting the problem

Definendo  $x = k_B T$  ed espressa in eV, dopo aver posto  $f(x) \equiv z \left( \frac{x}{k_B} \right)$  con  $k_B = 8.62 \times 10^{-5} \text{ eV K}^{-1}$ , il nostro problema è:

**Problem 1** *Let  $f(x) > 0$  in  $(0, +\infty)$  and such that*

$$\alpha_e x^{3/2} f(x) + \frac{\lambda f(x)}{f(x) + e^{-\frac{a}{x}}} = \alpha_h x^{3/2} \frac{e^{-\frac{b}{x}}}{f(x)} \quad (3)$$

where  $b > a > 0$ , while  $\lambda \geq 0$  is a free parameter.  $\alpha_e, \alpha_h > 0$  are expressed through Planck's reduced constant. Typical values:  $\alpha_e = 1096.24$ ,  $\alpha_h = 2013.92$ .

Show:

1. For  $x \rightarrow 0^+$ ,  $f(x)$  is an infinitesimal of infinitely large order. Furthermore,  $f \in C^\infty(0, +\infty)$ ,  $f \notin C^\omega(0, +\infty)$ , meaning that  $f(x)$  is continuous together with the derivatives of high order, but is not analytic.

2.  $\lim_{x \rightarrow 0^+} x \ln f(x) = -\frac{a+b}{2}$ .

[Hint: for  $0 < x \ll 1$  neglect  $x^{3/2} f(x)$ ]

**Soluzione**

**Question 1**

For  $\lambda = 0$  the (3) admits the only solution:

$$f(x) = \sqrt{\frac{\alpha_h}{\alpha_e}} e^{-\frac{b}{2x}} \quad (4)$$

which is manifestly an infinitesimal of infinitely large order for  $x \rightarrow 0^+$ :

$$\lim_{x \rightarrow 0^+} x^\alpha f(x) = 0, \quad \forall \alpha > 0$$

The same result is reached for the derivative  $f^{(n)}(x)$ . Extending these functions by continuity at the point  $x = 0$ , we have that  $f \in C^\infty$ , however it is not analytic at  $x = 0$ , since  $f^{(n)}(0) = 0, \forall n$ . It follows that the Taylor expansion of  $f(x)$  with initial point  $x=0$  returns the function identically zero, while in a right neighborhood of  $x = 0$  it is  $f(x) = 0$ . Furthermore

$$g(x) \stackrel{def}{=} x \ln f(x) = -\frac{b}{2} + \frac{x}{2} \ln \left( \frac{\alpha_h}{\alpha_e} \right) \quad (5)$$

which, unlike  $f$ , is class-based  $C^\omega$ .

For  $\lambda > 0$  we observe that the (3) is a third degree equation in  $f(x)$  that cannot be solved in closed form. For  $0 < x \ll 1$  its solutions behave like that of the quadratic equation in  $f(x)$ :

$$\frac{\lambda f(x)}{f(x) + e^{-\frac{a}{x}}} = \alpha_h x^{3/2} \frac{e^{-\frac{b}{x}}}{f(x)} \quad (6)$$

Solving:

$$f_\pm(x) = \frac{\alpha_h}{2\lambda} x^{3/2} e^{-\frac{b}{x}} \left( 1 \pm \sqrt{1 + \frac{4\lambda}{\alpha_h} \frac{e^{-\frac{b-a}{x}}}{x^{3/2}}} \right) \quad (7)$$

The problem (1) requires  $f(x) > 0$ , so

$$f(x) = \frac{\alpha_h}{2\lambda} x^{3/2} e^{-\frac{b}{x}} \left( 1 + \sqrt{1 + \frac{4\lambda}{\alpha_h} \frac{e^{\frac{b-a}{x}}}{x^{3/2}}} \right) \quad (8)$$

For the calculation of  $\lim_{x \rightarrow 0^+} f(x)$  we observe that

$$\lim_{x \rightarrow 0^+} x^{3/2} e^{-\frac{b}{x}} = 0, \quad \lim_{x \rightarrow 0^+} \frac{e^{\frac{b-a}{x}}}{x^{3/2}} = +\infty \quad b > a$$

so that  $\lim_{x \rightarrow 0^+} f(x) = 0 \cdot \infty$ . This indetermination can be removed by observing that in a right neighborhood of  $x = 0$  of arbitrarily small radius it succeeds  $\frac{e^{\frac{b-a}{x}}}{x^{3/2}} \gg 1$  since this ratio diverges positively for  $x \rightarrow 0^+$ . From this it follows that in the radicand of (8) we can neglect 1 with respect to the other term:

$$0 < x \ll 1 \implies f(x) \simeq \frac{\alpha_h}{2\lambda} x^{3/2} e^{-\frac{b}{x}} \left( 1 + \frac{2\lambda^{1/2}}{\alpha_h^{1/2}} \frac{e^{\frac{b-a}{2x}}}{x^{3/4}} \right)$$

In the same way  $\frac{e^{\frac{b-a}{2x}}}{x^{3/4}} \gg 1$  in the same right neighborhood as  $x = 0$ . In this order of approximation we have:

$$0 < x \ll 1 \implies f(x) \simeq \left( \frac{\alpha_h}{\lambda} \right)^{1/2} x^{3/4} e^{-\frac{a+b}{2x}} \quad (9)$$

So

$$\lim_{x \rightarrow 0^+} f(x) = \left( \frac{\alpha_h}{\lambda} \right)^{1/2} \lim_{x \rightarrow 0^+} x^{3/4} e^{-\frac{a+b}{2x}} = 0^+ \quad (10)$$

Since in a right neighborhood of  $x = 0$  with a small radius, the function  $f(x)$  is expressed as the product of an infinitesimal of order 3/4 by an infinitesimal of infinitely large order, we have that  $f(x)$  is in turn an infinitesimal of infinitely large order. By calculating the derivatives of however high order, we arrive at the same result. We conclude that  $f(x)$  is not analytic at  $x = 0$ .

### Question 2

From (9) it follows that in a right neighborhood of  $x = 0$  we have that  $g(x) = x \ln f(x)$  is expressed as

$$g(x) \simeq x \ln \left[ \left( \frac{\alpha_h}{\lambda} \right)^{1/2} x^{3/4} e^{-\frac{a+b}{2x}} \right] = \frac{3}{4} - \frac{a+b}{2} + \frac{x}{2} \ln \left( \frac{\alpha_h}{\lambda} \right)$$

so

$$\lim_{x \rightarrow 0^+} g(x) = -\frac{a+b}{2} \quad (11)$$

To establish the possible presence of the round-off, we plot the graph of  $g(x)$  in  $[x_{\min} = k_B T_{\min}, x_{\max} = k_B T_{\max}]$  where  $T_{\min, \max}$  define the thermal range assigned in the section (1). It follows  $x_{\min} = 8.62 \cdot 10^{-6}$  eV,  $x_{\max} = 0.086$  eV.

In Figure 4 Let's plot the graph of the function  $g(x) = x \ln f(x)$  with  $f(x)$  obtained by numerically solving the equation (3) and compared with the solution for  $x \ll 1$  i.e.  $g(x) = x \ln f(x)$  with  $f(x)$  given by (8). We see, therefore, that the change of variable  $T \rightarrow x = k_B T$  dampened the round-off oscillations. However, the latter is still present because according to the *Mathematica* kernel is  $\lim_{x \rightarrow 0^+} g(x) = 0$  contrary to the (11).

## 3 Conclusions

The round-off is triggered by the parameter  $\lambda > 0$ , since for  $\lambda = 0$  the solution  $f(x)$  is elementary expressible. Values  $\lambda > 0$  destroy the possibility of analytically solving the (3) and the corresponding solution is not elementary expressible.

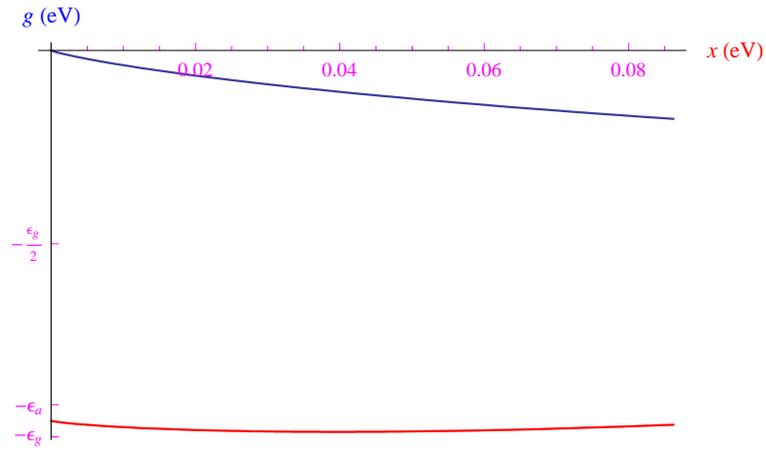


Figure 4: The curve in red is the graph of  $g(x) = x \ln f(x)$  where the function  $f(x)$  is given by (8). The curve in blue is the graph of  $g(x) = x \ln f(x)$  with  $f(x)$  solution of the (3).

## References

- [1] Kittel C. Kroemer H. *Termodinamica statistica*.
- [2] Wolfram S. *An Elementary Introduction to the Wolfram Language*.
- [3] Wagon S. *Mathematica in Action*. Springer