A blue geometric graphic consisting of a solid blue trapezoid on the left and a lighter blue triangle on the right, both pointing towards the right.

**Negative Mass and Negative Refractive
Index in Atom Nuclei**

Nuclear Wave Equation

Gravitational and Inertial Control

Part 5

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University Degree in Electronics Engineering

Negative Mass and Negative Refractive Index in Atom Nuclei

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Part-5

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The present study is divided into six parts – [Index of Part-5](#)

In This Paper

Following the analysis made in Part-3, the third remaining external force is evaluated in this study.

The **nuclear response to external forces** is analyzed with the aim to observe any **changes in the nuclear mass** and study the **behavior of the refractive index** under such changes.

The analysis will be performed in the time domain as well as in the frequency domain by means of the Fast Fourier Transform (FFT) method. The external forces applied to the nucleus were classified into three types:

- **The force originated by a polarized transverse electromagnetic wave (TEM)** (*see Part-3*)
- **The force originated by a polarized TEM plus a static electric field** (*see Part-4*)
- **The force originated from a signal plus a static electric field**

Abstract

Some efforts have been made to prove negative mass behavior through some experiments performed in mechanics [1], and other disciplines [9], as well as some theories in electrostatics [2,3,4,5,6,7,8], but I haven't found research about similar effects in atomic level, where the most elementary mass given by the atomic nucleus is to be found.

- Is the second Newton's law still valid with negative mass?
- What could happen if we make the atom behave in a negative mass regime?
- Is the negative refractive index related to negative mass?
- Are we able to control the magnitude of mass?
- Are we able to control the sign of mass?

The answers to these questions are given through this series of papers, with results that are coincident with experimental data, except for the negative mass regime. Experiments must be done to confirm or invalidate the theory developed in these articles. Needless to say, if experiments validate this theory, then a significant change in mankind is going to happen. In that

case, I strongly ask scientists to cooperate by making use of the derived technologies for good and refrain from doing it for evil.

Introduction

The theory presented in these papers, as described in Part-1 is based on three fundamental aspects that have proved to be extremely effective to describe physical phenomena and predicting results that agree with experimental data [10, 11, 12]:

- **Spinning Ring Model of Elementary Particles (toroidal ring of continuous charge)**
- **New Atomic Model**
- **The Universal Electrodynamic Force**

Based on the new atomic model, a shell arrangement of the nuclear particles has been assumed in Part 1, as shown in Fig. 1.

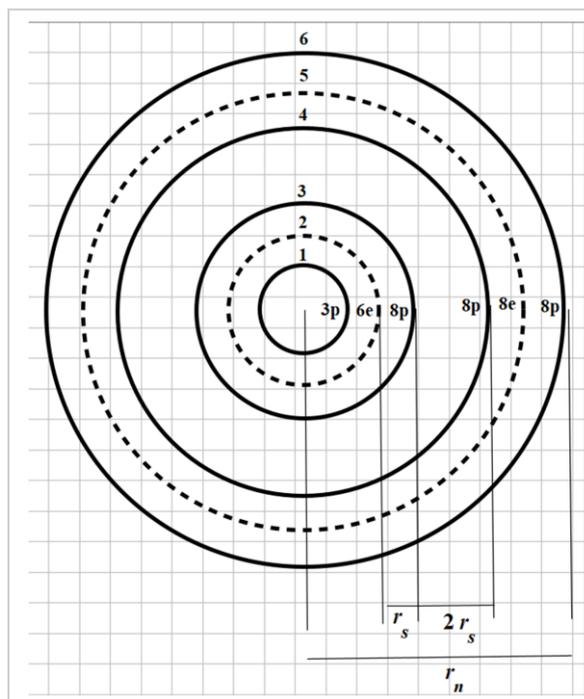


Figure 1

Assumed shell arrangement for Aluminum atomic nucleus

This sandwich configuration keeps the particles very tightly bound together. Note that at three shells in from the outermost shell, there are always two proton shells in a row for the larger nuclides.

This weak binding allows the outermost sandwich of shells to have liquid-like properties and forms the proper justification for a Liquid Drop Model of the nucleus.

As we already know, the torus ring model of the particles has an associated electric field as well as a magnetic field. However, due to the very tight packing configuration of the particles, we may safely assume that the distance among shells is extremely tiny and that the predominant force in the nucleus is of electrostatic origin, while the weaker magnetic forces will add some contribution to the equilibrium distance between each shell.

As demonstrated in Part 1, mass is an intrinsic property of the atomic nucleus. Under natural circumstances, it has a constant universal magnitude and is always positive. However, with some proper external agents, **we might be able to manipulate the intrinsic mass by changing its magnitude and sign.**

III. Nuclear Response to Force Caused by an Electric Signal plus a Static Electric Field

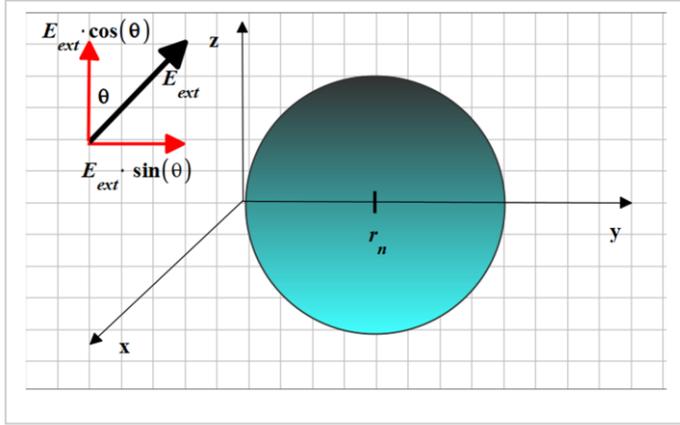


Figure 2

External signal and static electric field acting on the nucleus

Let's take an external excitation composed of a uniform electric field E_f plus a signal with amplitude E_s and frequency ω , which is overlaid with the static field.

Assume also that the composed field is given as in Fig. 2 by $E_{ext} = E_f + E_s \cos(\omega t)$, where E_f is the magnitude of the static electric field.

Assume that the external field E_{ext} is parallel to the z-y plane, making an angle θ with the z-axis.

The components of the field in cartesian coordinates are: $\vec{E}_x = 0\hat{i} + E_{ext}\hat{j} + E_{ext}\hat{k}$, and in spherical coordinates we have,

$$\vec{E}_x = -((\sin(\phi) - 1)\cos(\theta)^2 - \sin(\phi))(E_f + E_s \cos(\omega t))\hat{r} + \sin(\theta)\cos(\theta)(\sin(\phi) - 1)(E_f + E_s \cos(\omega t))\hat{\theta} + \cos(\phi)\sin(\theta)(E_f + E_s \cos(\omega t))\hat{\phi} \quad (1)$$

Consider that the effect of the force caused by this external field is distributed in the entire nucleus volume. Then, a reasonable approach will be to average the external field over the nuclear volume, that is,

$$\vec{E}_{xav} = \frac{1}{V_n} \int_0^{2\pi} \int_0^\pi \int_0^{r_n} \vec{E}_x (r - r_n)^2 \sin(\theta) dr d\theta d\phi$$

$$\vec{E}_{xav} = \frac{3}{4\pi r_n^3} \int_0^{2\pi} \int_0^\pi \int_0^{r_n} \vec{E}_x (r - r_n)^2 \sin(\theta) dr d\theta d\phi$$

The result of this integral gives us the averaged field in the nucleus,

$$\vec{E}_{xav} = \frac{(E_f + E_s \cos(\omega t))}{3} \hat{r} \quad (2)$$

The electric field \vec{E}_f gives the direction of the external force. The total force exerted by the field on the nucleus is the sum of the forces on every proton and electron, which form the nuclear shells. According to our 6 shells structure (Fig. 1), we have:

$F_{ext} = QE_{xav}$, where Q is the total nuclear charge. Based on the shell structure of Fig. 1 and taking the sign of the charges for each shell, we get

$$\vec{F}_{ext} = \vec{E}_{xav} (3q - 6q + 8q + 8q - 8q + 8q) = 13q \vec{E}_{xav} \quad (3)$$

By replacing Eq. (2) in (3), we obtain the (rough) final expression of the external force acting on the nucleus:

$$\vec{F}_{ext} = \frac{13q}{3} (E_f + E_s \cos(\omega t)) \hat{r} \quad (4)$$

Now that we have the external force exerted on the nucleus caused by an electric signal plus a static electric field, let's evaluate the nuclear response related to mass and refractive index behaviors.

III.a Nuclear Mass Analysis due to a Force caused by a Static Electric Field plus an Electric Signal – Partial or Total Energy Absorption

The intrinsic net force in the nucleus was already defined with Eq. (23) in Part-1. Now we have the action of an external force acting on the nucleus that will interact with the internal force. By applying Newton's second law, we have

$$\sum F = m_n \cdot a_{ep} = F_{net} + F_{ext} \quad \Rightarrow \quad m_n \cdot a_{ep} = F_{net} + F_{ext}$$

$$m_n = \frac{1}{a_{ep}} \cdot (F_{net} + F_{ext}) \quad (7)$$

By replacing the forces in (7), we obtain the expression of the nuclear mass for this case:

$$m_n = \frac{1}{a_{ep}(t)} \cdot \left(- \frac{3.410^{12} \cdot q^2 \left(1 - \frac{v_{ep}^2(t)}{c^2} + \frac{v_{ep}^2(t) r_{ep}(t) a_{ep}(t)}{c^4} + \frac{v_{ep}^4(t)}{c^4} + \frac{2 r_{ep}(t) a_{ep}(t)}{c^2} \right)}{r_{ep}^2(t)} + \frac{2.0510^{13} \cdot q^2}{r_n^2} + \frac{13 q}{3} (E_f + E_s \cos(\omega t)) \right) \quad (8)$$

Recall that:

$$r_{ep}(t) = (0.37 r_n + A_e \cos(\omega_e t) - A_p \cos(\omega_p t))$$

$$v_{ep}(t) = (-A_e \omega_e \sin(\omega_e t) + A_p \sin(\omega_p t) \omega_p)$$

$$a_{ep}(t) = (-A_e \omega_e^2 \cos(\omega_e t) + A_p \omega_p^2 \cos(\omega_p t))$$

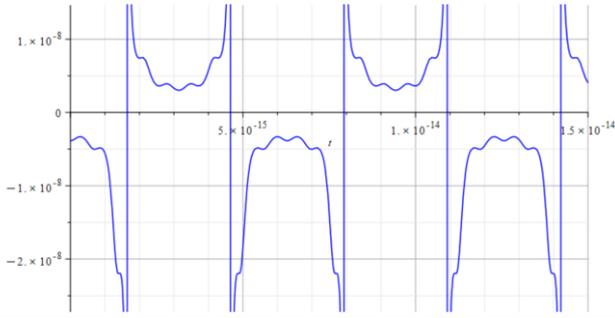
Some graphs as examples are shown below to have a perception of what could be done to modify the nuclear mass magnitude and sign.

The main parameters used for the net force are:

$$r_n = 3.5 \cdot 10^{-15} [m]; \quad A_e = 2 \cdot 10^{-16} [m]; \quad A_p = 10^{-3} A_e [m]; \quad N_0 = 378; \quad N_s = 312$$

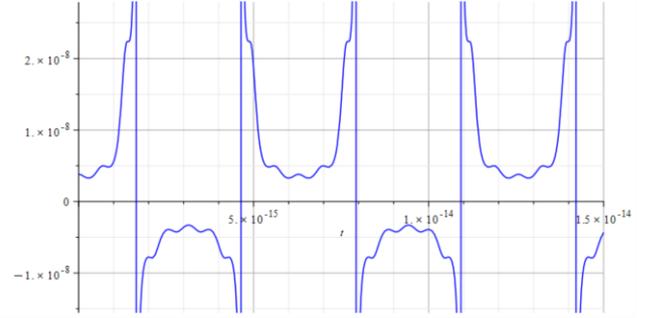
$$\omega_e = 10^{15} \left[\frac{1}{s} \right]; \quad \omega_p = 10^{16} \left[\frac{1}{s} \right]$$

Time Analysis of the Nuclear Mass



$+E_f$

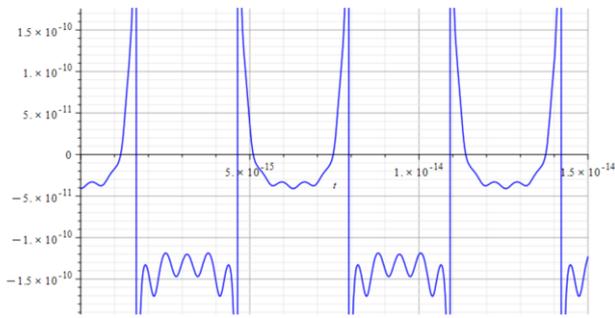
Mass vs. time for the following parameters:
 $\omega = 10^{15} \left[\frac{1}{s}\right]$ $E_s = 10^{18} \left[\frac{V}{m}\right]$ $E_f = 10^{24} \left[\frac{V}{m}\right]$



$-E_f$

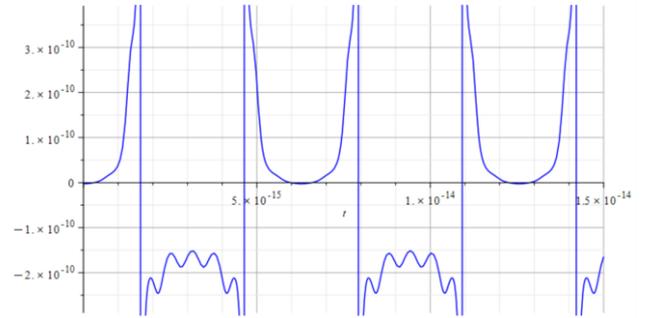
Mass vs. time for the following parameters:
 $\omega = 10^{15} \left[\frac{1}{s}\right]$ $E_s = 10^{18} \left[\frac{V}{m}\right]$ $E_f = -10^{24} \left[\frac{V}{m}\right]$

Figure 3



$+E_f$

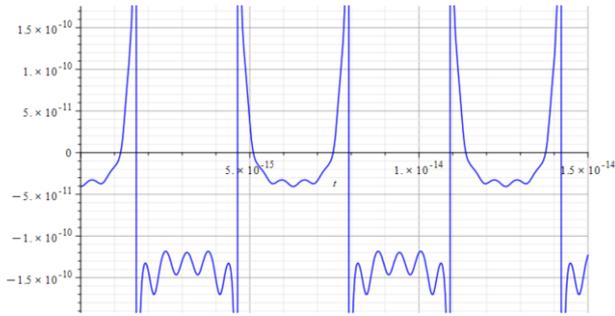
Mass vs. time for the following parameters:
 $\omega = 10^{15} \left[\frac{1}{s}\right]$ $E_s = 10^{18} \left[\frac{V}{m}\right]$ $E_f = 5 \cdot 10^{21} \left[\frac{V}{m}\right]$



$-E_f$

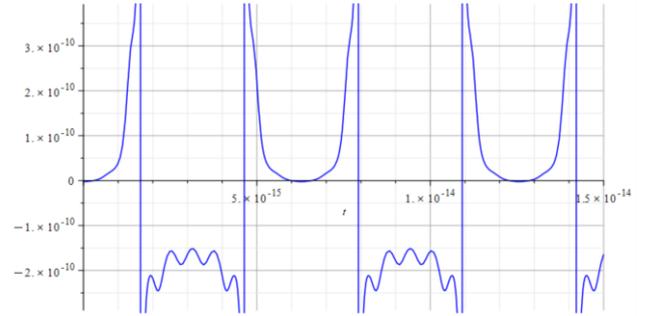
Mass vs. time for the following parameters:
 $\omega = 10^{15} \left[\frac{1}{s}\right]$ $E_s = 10^{18} \left[\frac{V}{m}\right]$ $E_f = -5 \cdot 10^{21} \left[\frac{V}{m}\right]$

Figure 4



$+E_f$

Mass vs. time for the following parameters:
 $\omega = 10^{10} \left[\frac{1}{s}\right]$ $E_s = 10^{18} \left[\frac{V}{m}\right]$ $E_f = 5 \cdot 10^{21} \left[\frac{V}{m}\right]$



$-E_f$

Mass vs. time for the following parameters:
 $\omega = 10^{10} \left[\frac{1}{s}\right]$ $E_s = 10^{18} \left[\frac{V}{m}\right]$ $E_f = -5 \cdot 10^{21} \left[\frac{V}{m}\right]$

Figure 5

From the period of the mass plot, we determine that the oscillation frequency is approximately:

$$f = 1.6 \cdot 10^{14} \text{ [Hz]}$$

Frequency Analysis of the Nuclear Mass with FFT

Total number of samples $N = 2^{14}$, sampling frequency $f_s = 2^3 f_p$ (proton frequency), which gives a frequency resolution $\Delta f = \frac{f_s}{N} [Hz]$ and a total acquisition time of $T = \frac{N}{f_s} [s]$. The frequency at the i -sample number on the plot is determined by $f = \frac{N(i)}{T} [Hz]$.

For $+E_f$

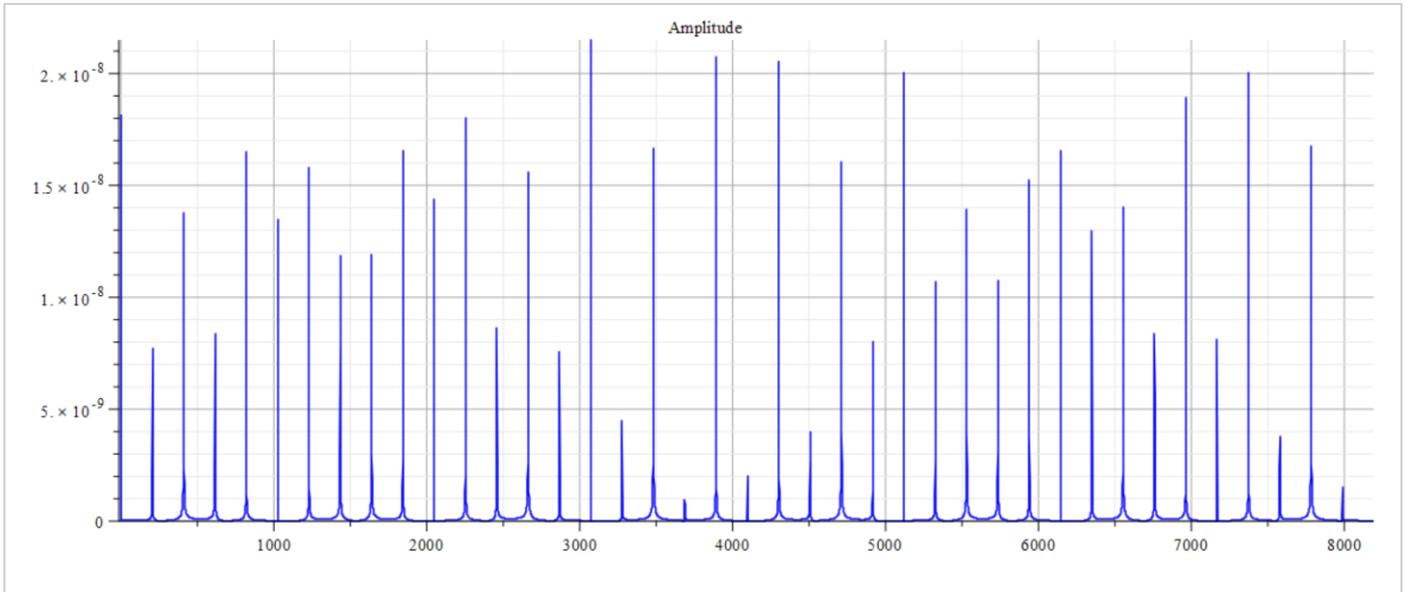


Figure 6

Frequency spectrum for the following parameters: $\omega = 10^{15} [\frac{1}{s}]$ $E_s = 10^{18} [\frac{V}{m}]$ $E_f = 10^{24} [\frac{V}{m}]$

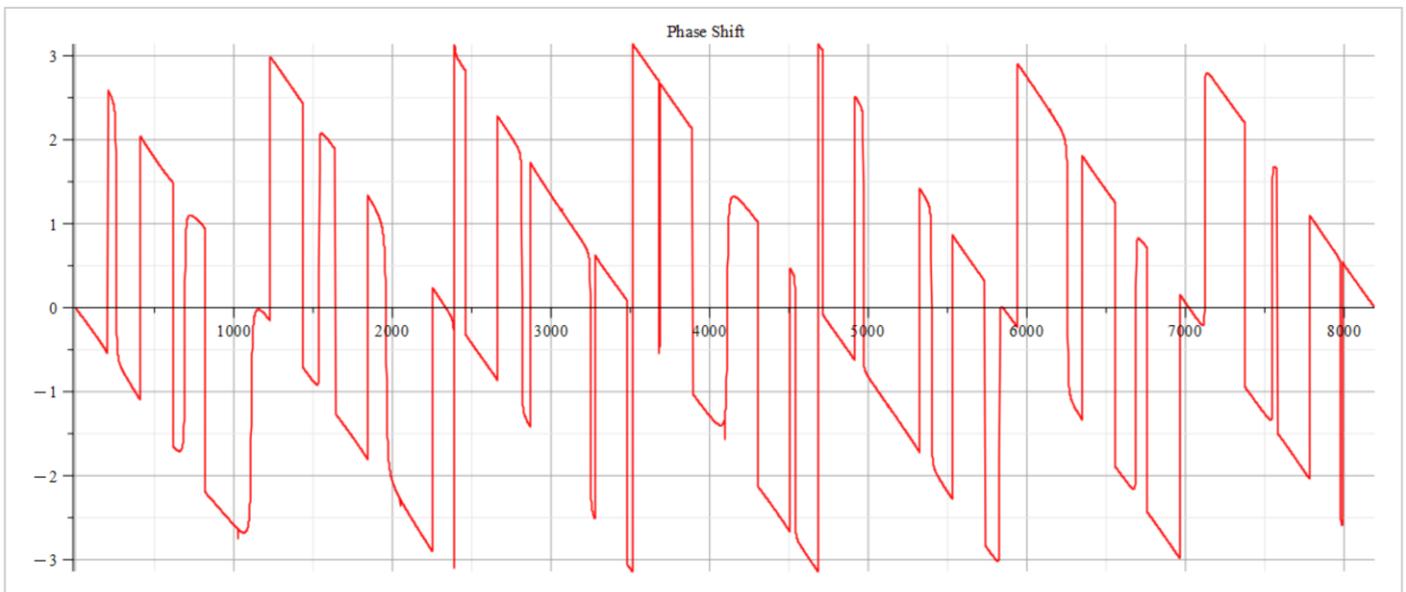


Figure 7

Phase shift for for the following parameters: $\omega = 10^{15} [\frac{1}{s}]$ $E_s = 10^{18} [\frac{V}{m}]$ $E_f = 10^{24} [\frac{V}{m}]$

For $-E_f$

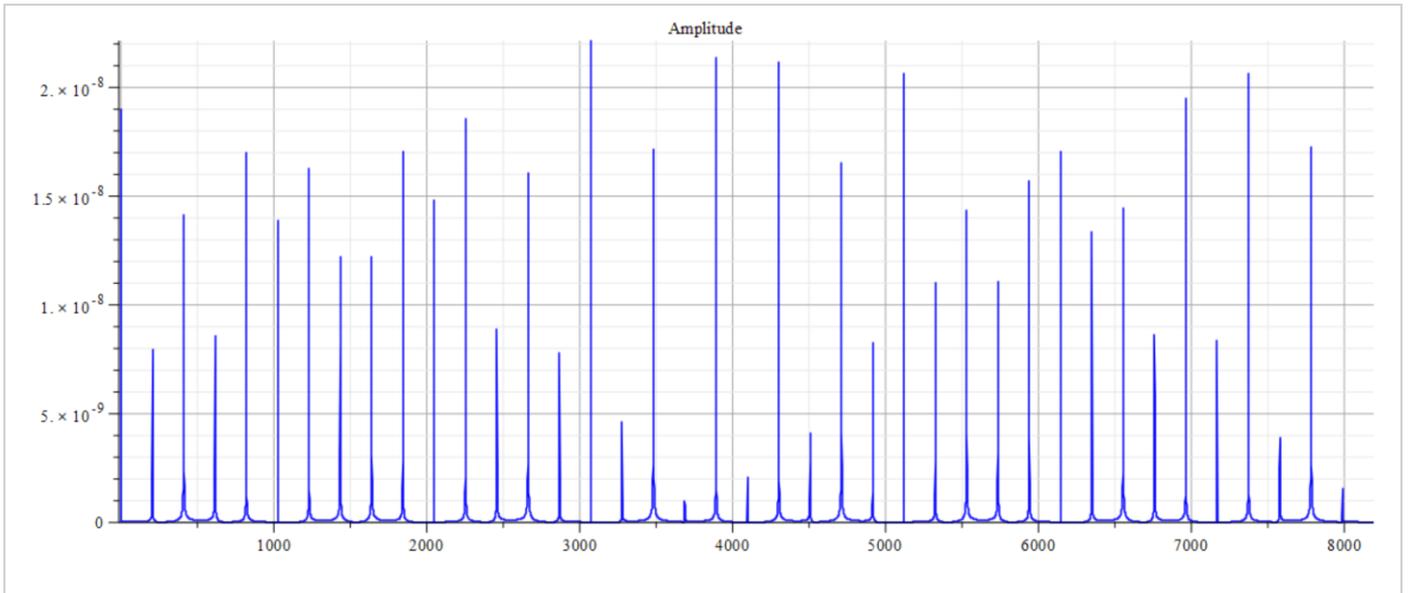


Figure 8

Frequency spectrum for the following parameters: $\omega = 10^{15} [\frac{1}{s}]$ $E_s = 10^{18} [\frac{V}{m}]$ $E_f = -10^{24} [\frac{V}{m}]$

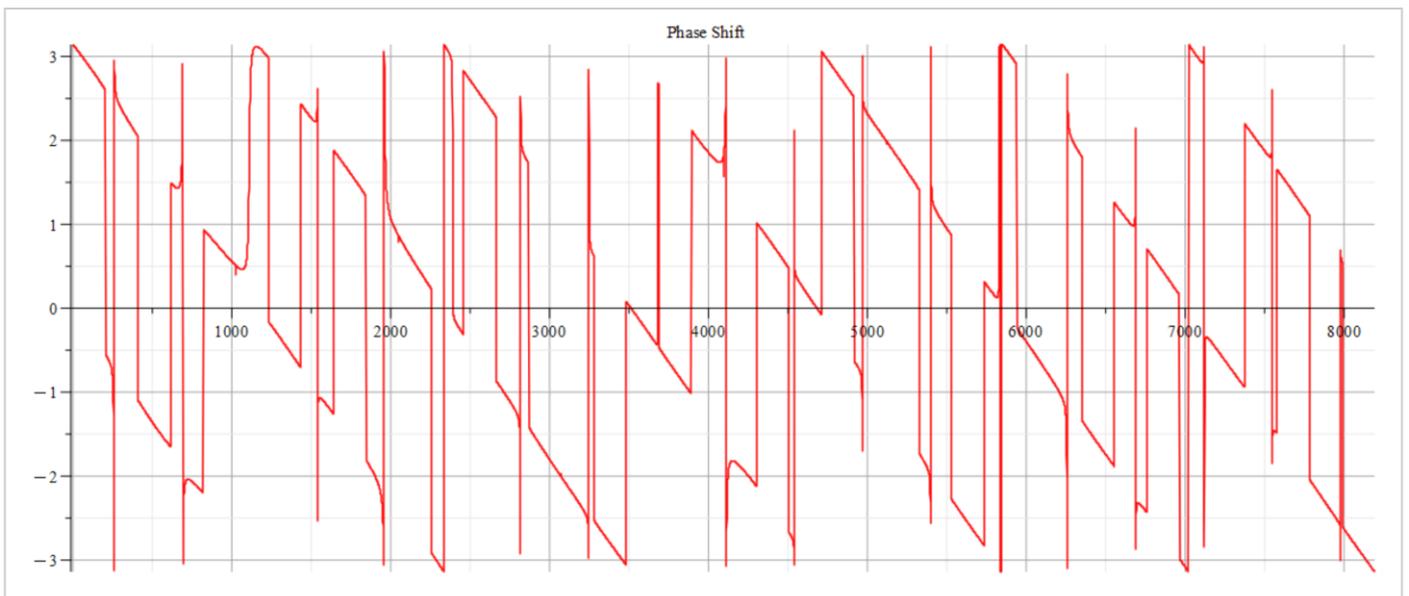


Figure 9

Phase shift for the following parameters : $\omega = 10^{15} [\frac{1}{s}]$ $E_s = 10^{18} [\frac{V}{m}]$ $E_f = -10^{24} [\frac{V}{m}]$

For $+E_f$

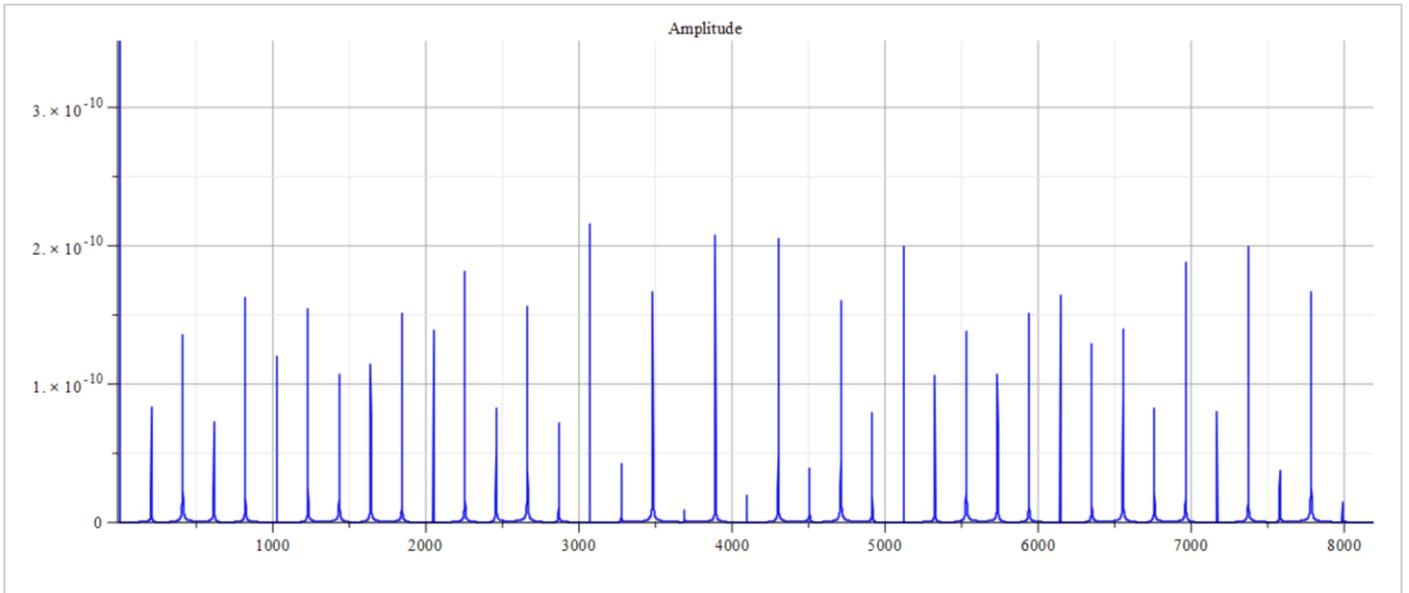


Figure 10

Frequency spectrum for the following parameters: $\omega = 10^{15} \left[\frac{1}{s} \right]$ $E_s = 10^{18} \left[\frac{V}{m} \right]$ $E_f = 5 \cdot 10^{21} \left[\frac{V}{m} \right]$

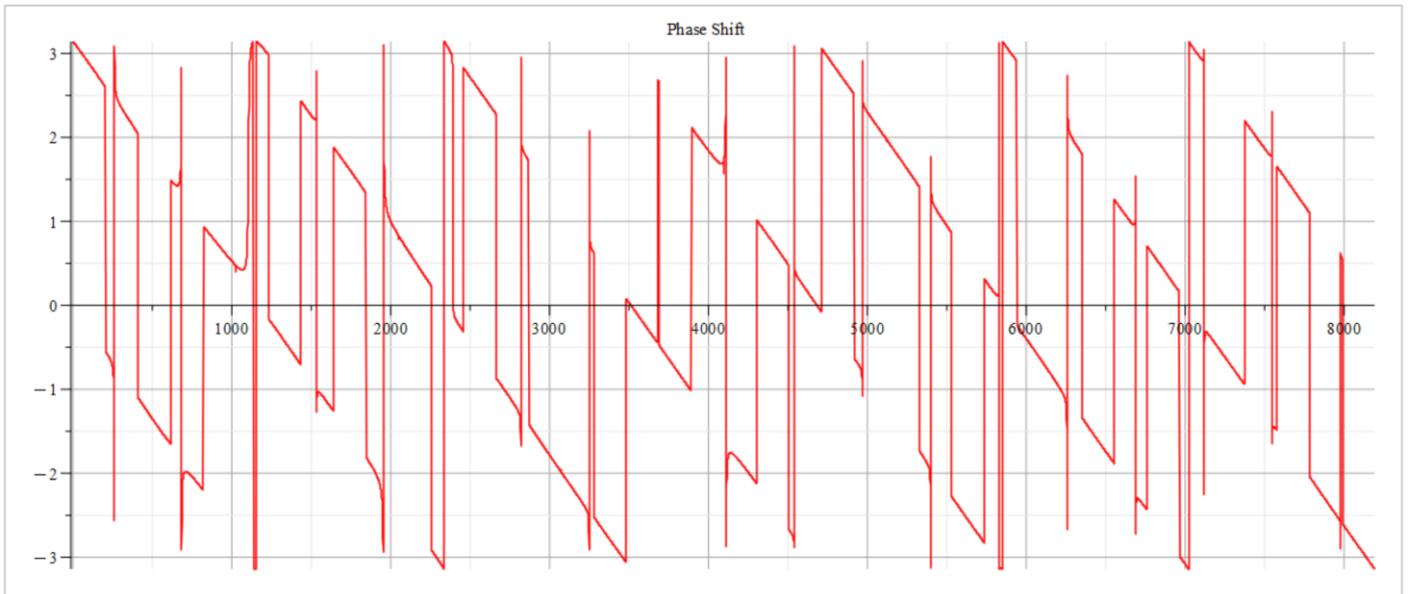


Figure 11

Phase shift for the following parameters: $\omega = 10^{15} \left[\frac{1}{s} \right]$ $E_s = 10^{18} \left[\frac{V}{m} \right]$ $E_f = 5 \cdot 10^{21} \left[\frac{V}{m} \right]$

For $-E_f$

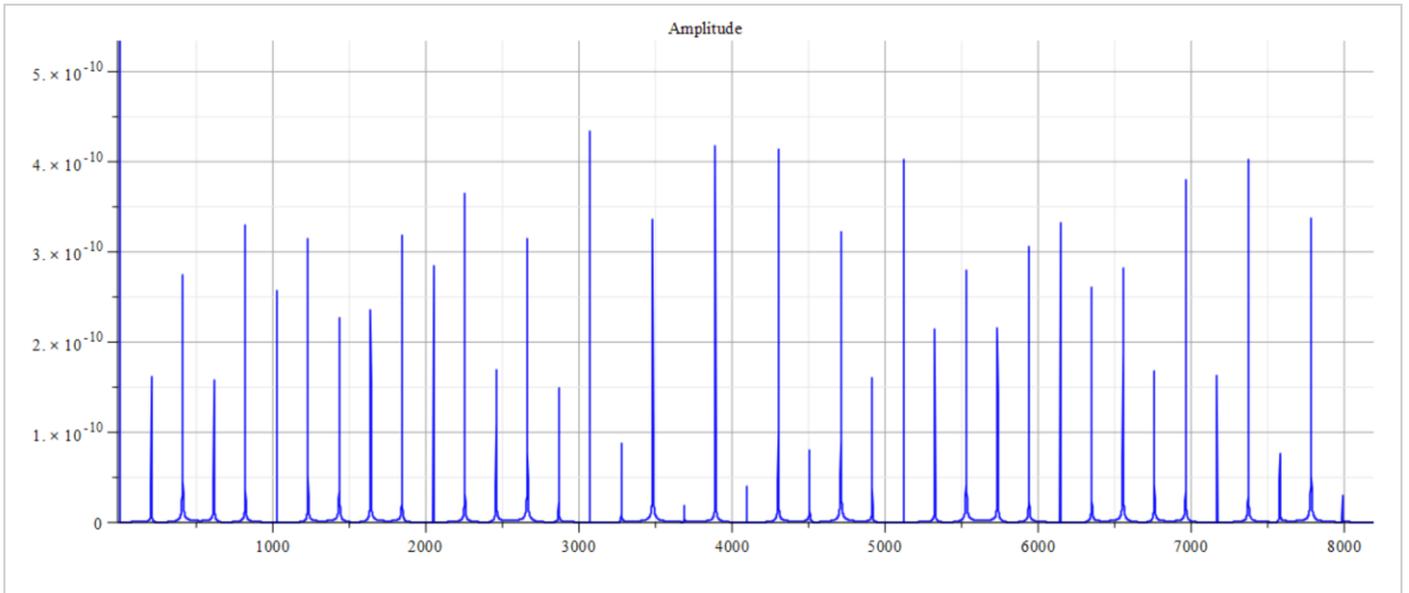


Figure 12

Frequency spectrum for the following parameters: $\omega = 10^{15} \left[\frac{1}{s} \right]$ $E_s = 10^{18} \left[\frac{V}{m} \right]$ $E_f = -5 \cdot 10^{21} \left[\frac{V}{m} \right]$

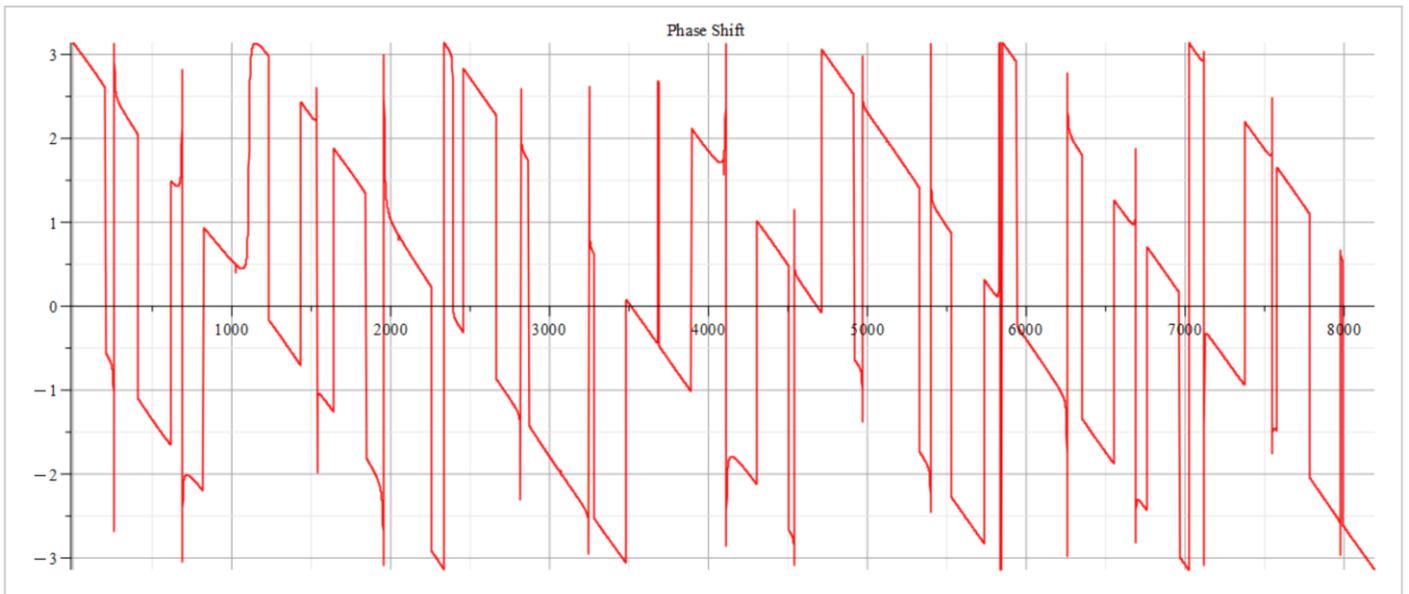


Figure 13

Phase shift for the following parameters: $\omega = 10^{15} \left[\frac{1}{s} \right]$ $E_s = 10^{18} \left[\frac{V}{m} \right]$ $E_f = -5 \cdot 10^{21} \left[\frac{V}{m} \right]$

For + E_f

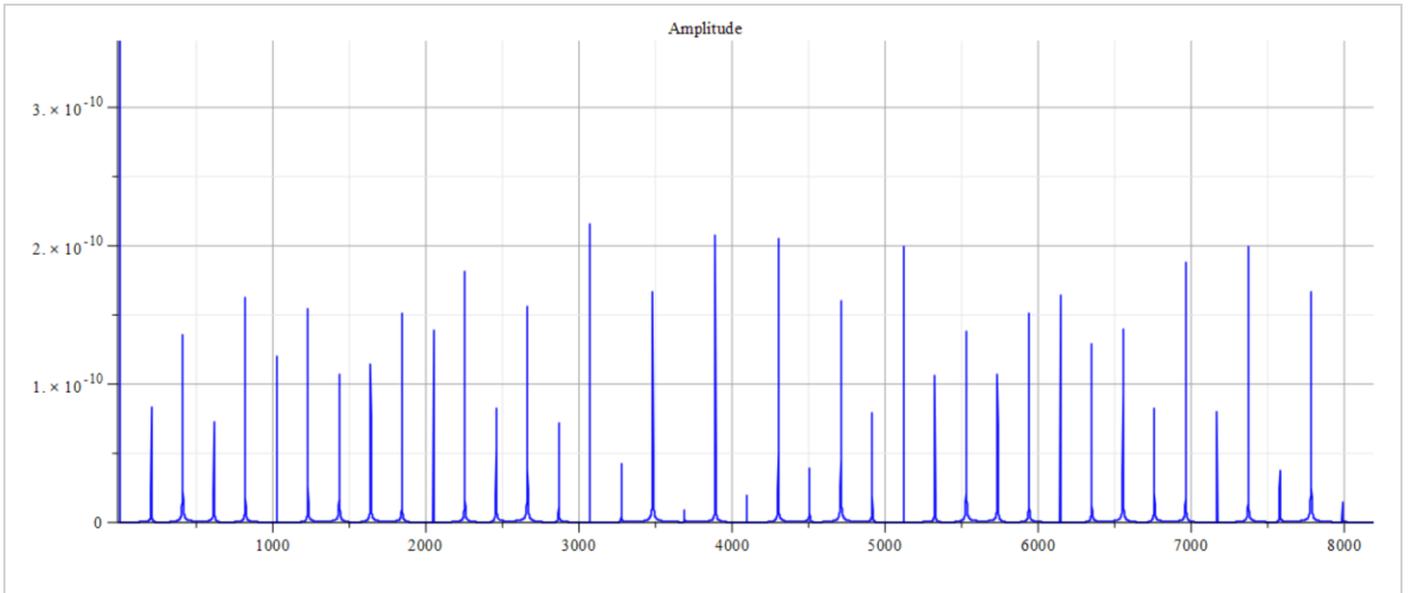


Figure 14

Frequency spectrum for the following parameters: $\omega = 10^{10} \left[\frac{1}{s} \right]$ $E_s = 10^{18} \left[\frac{V}{m} \right]$ $E_f = 5 \cdot 10^{21} \left[\frac{V}{m} \right]$

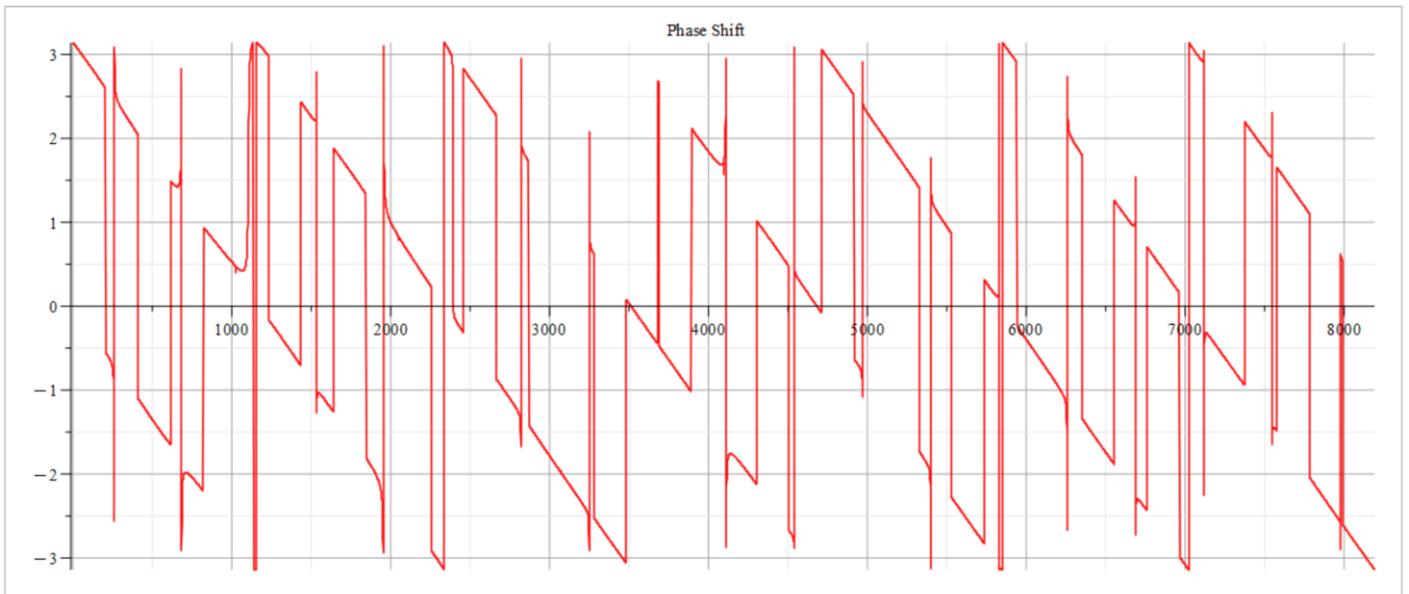


Figure 15

Phase shift for the following parameters: $\omega = 10^{10} \left[\frac{1}{s} \right]$ $E_s = 10^{18} \left[\frac{V}{m} \right]$ $E_f = 5 \cdot 10^{21} \left[\frac{V}{m} \right]$

For - E_f

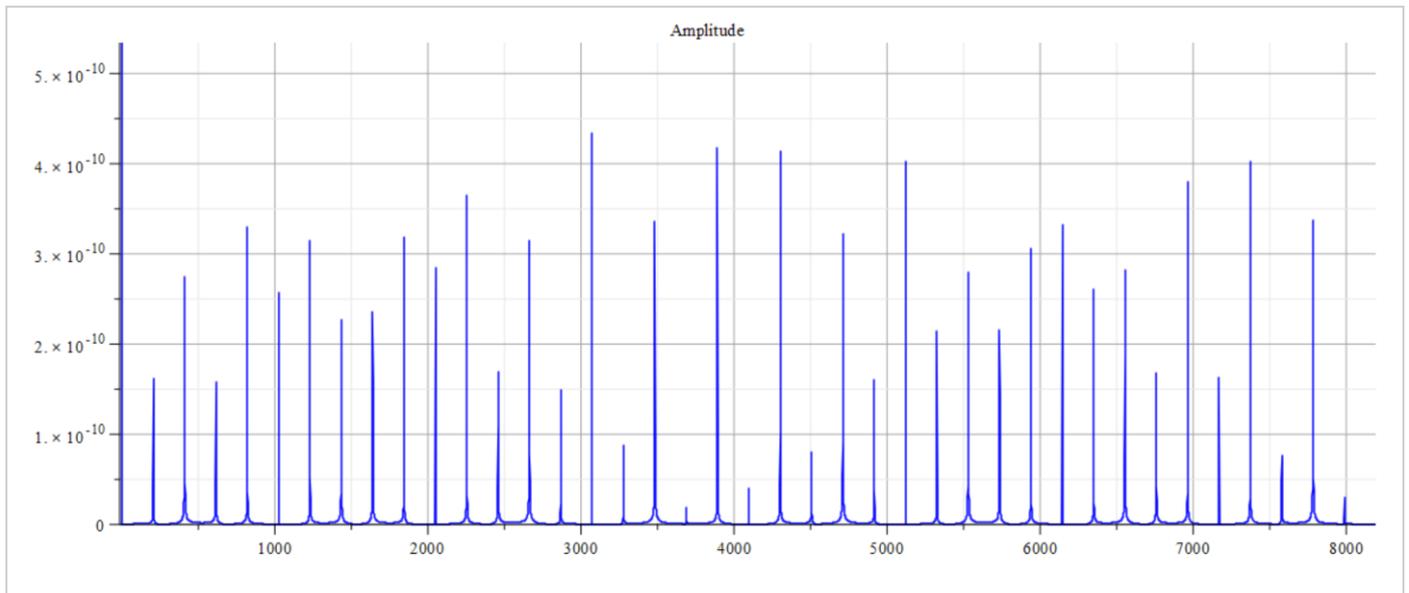


Figure 16

Frequency spectrum for the following parameters: $\omega = 10^{10} \left[\frac{1}{s} \right]$ $E_s = 10^{18} \left[\frac{V}{m} \right]$ $E_f = -5 \cdot 10^{21} \left[\frac{V}{m} \right]$

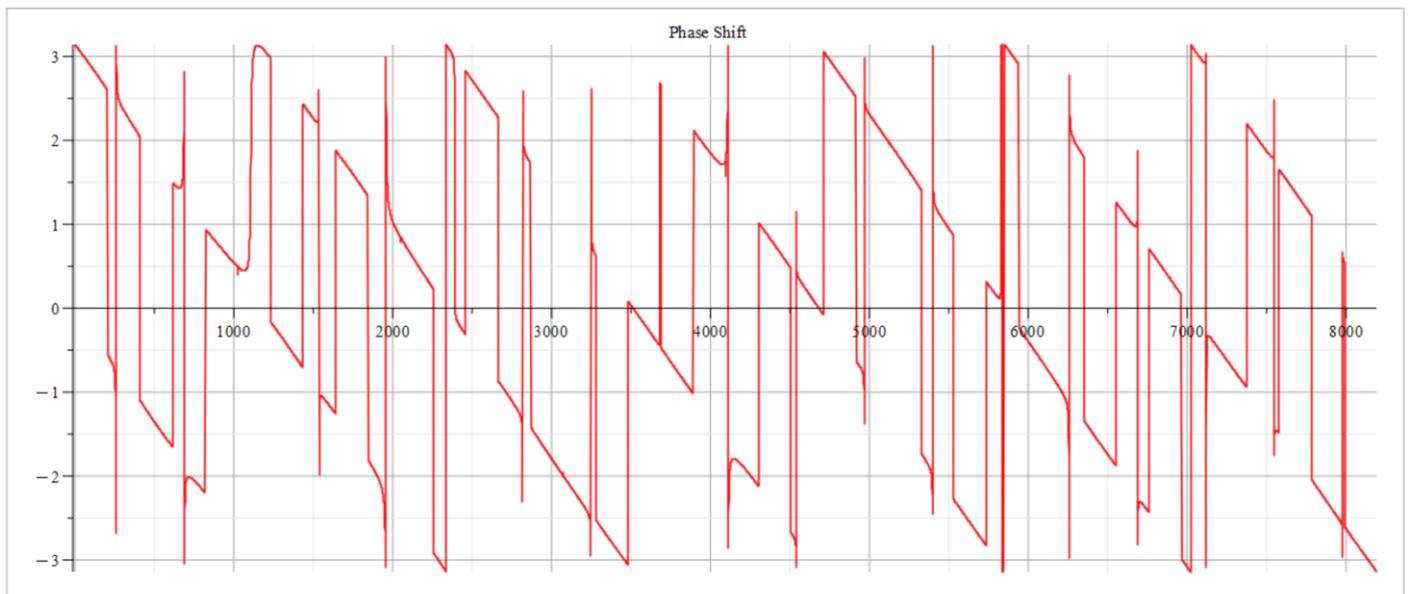


Figure 17

Phase shift for the following parameters: $\omega = 10^{10} \left[\frac{1}{s} \right]$ $E_s = 10^{18} \left[\frac{V}{m} \right]$ $E_f = -5 \cdot 10^{21} \left[\frac{V}{m} \right]$

From the Fourier frequency analysis, we see that the main frequency is: $f_0 = 1.6 \cdot 10^{14} [Hz]$.

In general, the main frequency and harmonics are given by the following formula:

$$f_n = f_0 + n f_0, \quad n = 0, 1, 2, 3, \dots$$

III.b Refractive Index Analysis due to a Force caused by a Static Electric Field plus an Electric Signal – Partial or Total Energy Absorption

When the nucleus is under the action of external forces, and if it doesn't break apart, then we can assume that a dynamic equilibrium state must exist. Under such circumstances, Newton's second law requires that the sum of forces be equal to zero, $\vec{F}_{net} + \vec{F}_{ext} = 0$, that is,

$$F_{net} = -F_{ext} \quad (9)$$

Recall that the net nuclear force has already been written in terms of the index of refraction in Part-1, Eq. (23a):

$$\vec{F}_{net} = -\frac{378k q^2}{r_{ep}^2(t)} \left(1 - \frac{1}{n^2} + \frac{v_{ep}^2(t)r_{ep}(t)a_{ep}(t)}{n^2 c^2} + \frac{1}{n^4} + \frac{2r_{ep}(t)a_{ep}(t)}{c^2} \right) \hat{r} + \frac{2279.035793k q^2 \hat{r}}{r_n^2}$$

Now we can equate the forces according to Eq. (9), then solve for "n",

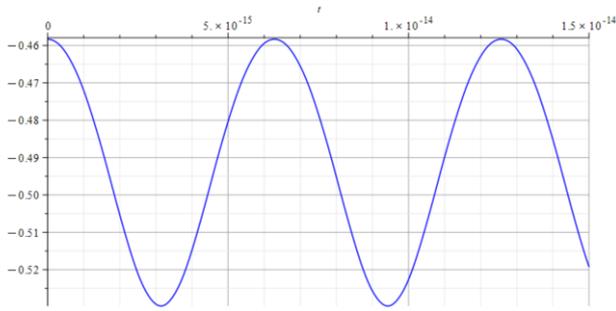
$$-\frac{3.410^{12} \cdot q^2 \left(1 - \frac{1}{n^2} + \frac{r_{ep}(t)a_{ep}(t)}{n^2 c^2} + \frac{1}{n^4} + \frac{2r_{ep}(t)a_{ep}(t)}{c^2} \right)}{r_{ep}^2(t)} + \frac{2.0510^{13} \cdot q^2}{r_n^2} = -\frac{13q}{3} (E_f + E_s \cos(\omega t)) \quad (10)$$

The refractive index "n" is a somewhat long-expression which is nonsense to copy here.

Some plots as examples are shown below, where the main used parameters are:

$$r_n = 3.5 \cdot 10^{-15} [m]; \quad A_e = 2 \cdot 10^{-16} [m]; \quad A_p = 10^{-3} A_e [m]; \quad N_0 = 378; \quad N_s = 312$$

$$\omega_e = 10^{15} \left[\frac{1}{s} \right]; \quad \omega_p = 10^{16} \left[\frac{1}{s} \right]$$



$+E_f$

Refractive Index vs. time for the following parameters: **Figure 18**

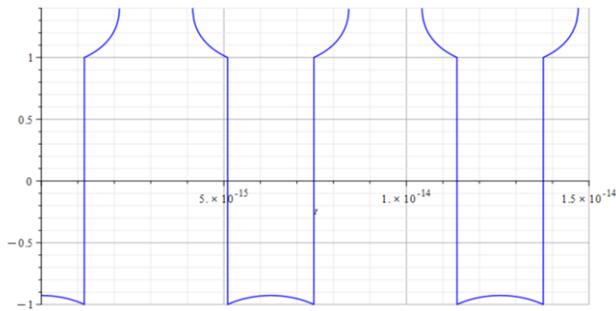
$$\omega = 10^{15} \left[\frac{1}{s} \right] \quad E_s = 10^{18} \left[\frac{V}{m} \right] \quad E_f = 10^{24} \left[\frac{V}{m} \right]$$

No graph was obtained due to complex number result

$-E_f$

Refractive Index vs. time for the following parameters:

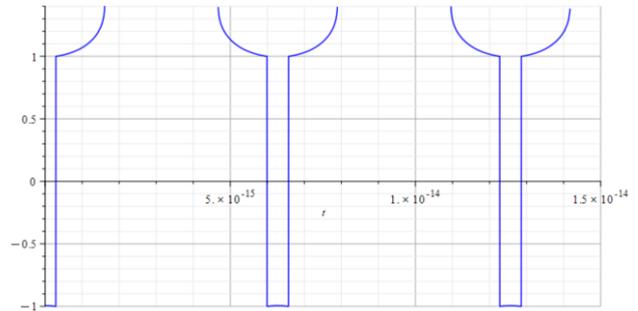
$$\omega = 10^{15} \left[\frac{1}{s} \right] \quad E_s = 10^{18} \left[\frac{V}{m} \right] \quad E_f = -10^{24} \left[\frac{V}{m} \right]$$



$+E_f$

Refractive Index vs. time for the following parameters: **Figure 19**

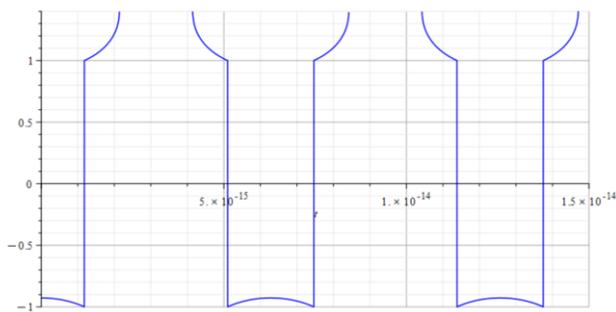
$$\omega = 10^{15} \left[\frac{1}{s} \right] \quad E_s = 10^{18} \left[\frac{V}{m} \right] \quad E_f = 5 \cdot 10^{21} \left[\frac{V}{m} \right]$$



$-E_f$

Refractive Index vs. time for the following parameters:

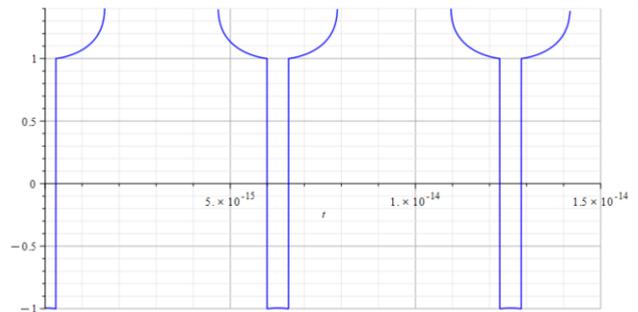
$$\omega = 10^{15} \left[\frac{1}{s} \right] \quad E_s = 10^{18} \left[\frac{V}{m} \right] \quad E_f = -5 \cdot 10^{21} \left[\frac{V}{m} \right]$$



$+E_f$

Refractive Index vs. time for the following parameters: **Figure 20**

$$\omega = 10^{10} \left[\frac{1}{s} \right] \quad E_s = 10^{18} \left[\frac{V}{m} \right] \quad E_f = 5 \cdot 10^{21} \left[\frac{V}{m} \right]$$



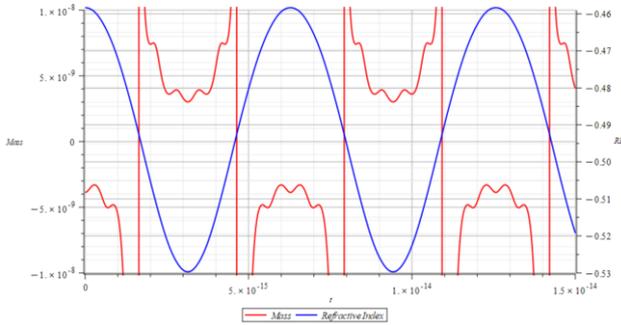
$-E_f$

Refractive Index vs. time for the following parameters:

$$\omega = 10^{10} \left[\frac{1}{s} \right] \quad E_s = 10^{18} \left[\frac{V}{m} \right] \quad E_f = -5 \cdot 10^{21} \left[\frac{V}{m} \right]$$

III.c Comparison of Mass with Refractive Index Behavior due to a Force caused by a Static Electric Field plus an Electric Signal – Partial or Total Energy Absorption

To analyze the changes in the refractive index “n” with respect to changes in nuclear mass, overlaid graphs of both quantities are shown below, which uncover interesting results.



No graph was obtained due to complex number result in refractive index

+E_f

Figure 21

-E_f

Mass & Refractive Index vs. time for the following parameters:

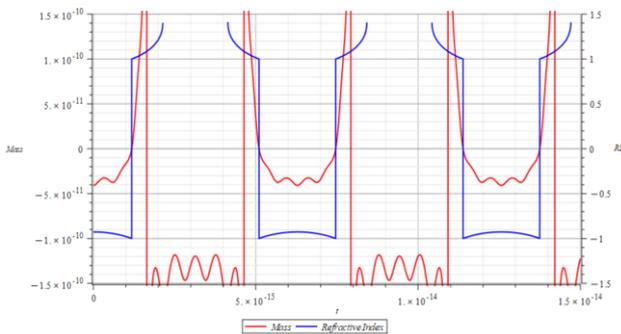
$$\omega = 10^{15} \left[\frac{1}{s} \right] \quad E_s = 10^{18} \left[\frac{V}{m} \right] \quad E_f = 10^{24} \left[\frac{V}{m} \right]$$

Mass & Refractive Index vs. time for the following parameters.:

$$\omega = 10^{15} \left[\frac{1}{s} \right] \quad E_s = 10^{18} \left[\frac{V}{m} \right] \quad E_f = -10^{24} \left[\frac{V}{m} \right]$$

The Refractive Index here is always negative, independent of the mass sign, and oscillates by following the rate of change of the mass. It seems that the refractive index depends on the derivative of the mass, that is

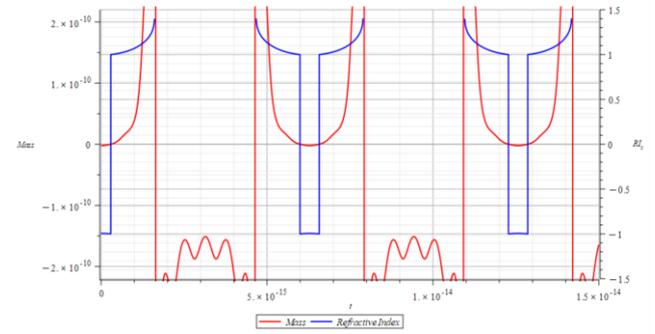
$$n \propto \frac{dm}{dt}$$



+E_f

Figure 22

-E_f



Mass & Refractive Index vs. time for the following parameters:

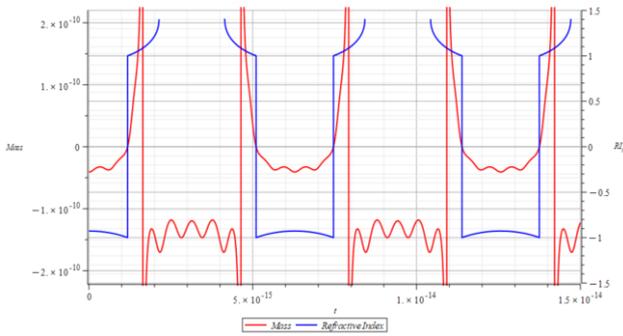
$$\omega = 10^{15} \left[\frac{1}{s} \right] \quad E_s = 10^{18} \left[\frac{V}{m} \right] \quad E_f = 5 \cdot 10^{21} \left[\frac{V}{m} \right]$$

Mass & Refractive Index vs. time for the following parameters.:

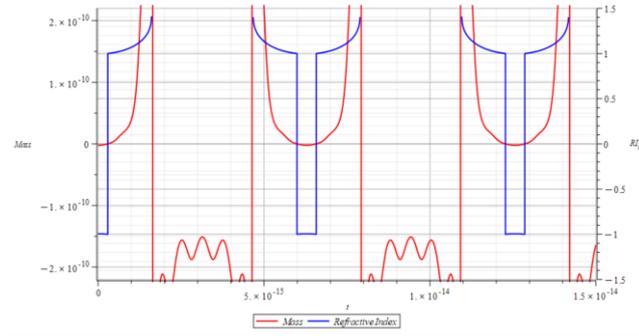
$$\omega = 10^{15} \left[\frac{1}{s} \right] \quad E_s = 10^{18} \left[\frac{V}{m} \right] \quad E_f = -5 \cdot 10^{21} \left[\frac{V}{m} \right]$$

Here, the Refractive Index again depends on the slope of the mass plot. It decreases with negative mass slope, and vice versa, until reaching the m=0 point, where abruptly switches between n=±1. It seems that the refractive index depends on the derivative of the mass, that is

$$n \propto \frac{dm}{dt}$$



+E_f



-E_f

Figure 23

Mass & Refractive Index vs. time for the following parameters:
 $\omega = 10^{10} \left[\frac{1}{s} \right]$ $E_s = 10^{18} \left[\frac{V}{m} \right]$ $E_f = 5 \cdot 10^{21} \left[\frac{V}{m} \right]$

Mass & Refractive Index vs. time for the following parameters.:
 $\omega = 10^{10} \left[\frac{1}{s} \right]$ $E_s = 10^{18} \left[\frac{V}{m} \right]$ $E_f = -5 \cdot 10^{21} \left[\frac{V}{m} \right]$

Here, the Refractive Index again depends on the slope of the mass plot. It decreases with negative mass slope, and vice versa, until reaching the $m=0$ point, where abruptly switches between $n=\pm 1$. It seems that the refractive index depends on the derivative of the mass, that is

$$n \propto \frac{dm}{dt}$$

This is an important result that tells us that the refractive index behavior is like a “beacon”, signaling the zones of the negative mass regime, as well as the points of mass changes.

Conclusions

It has been demonstrated that the application of the Universal Electrodynamic Force to the new Atomic Model predicts important changes in nuclear mass when an external force caused by an Electric Signal with an added Static Electric Field is acting on the atomic nucleus.

As we have demonstrated in previous papers, mass is an electrodynamic quantity and as such, it can be manipulated at will. It was demonstrated in this paper that besides the previously analyzed external agents, one more mean can be used to achieve mass changes.

It has been demonstrated that the magnitude and the sign of the mass can be modified by changing the amplitude and/or frequency of the external signal, within a certain range, and by modifying the magnitude and sign of the static electric field.

The change of the refractive index values during mass change was also clearly demonstrated. There is clear evidence that the refractive index is proportional to the rate of change of the mass, i.e., to the derivative of mass with respect to time.

The refractive index can be used as an aid to search for the negative mass region of the nucleus, as well as in any piece of "macro" material. The refractive index is a "beacon" that signals the exact point of mass sign change and its range, and any mass changes in general.

Fourier's analysis shows in the phase shift graphs many swings of phase between $\pm\pi$, which clearly indicate resonance states in the nucleus at those frequencies, as well interferences with the external agent.

In previous papers, we described that the main force that keeps the nuclear shell in a very tight packing structure in such a tiny space is the electrostatic force. This is really an enormous force. It means that we also need huge external electromagnetic fields to achieve some nuclear interaction, and this may represent a technical limitation in the present time.

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