

Interaction of Complex Scalar Fields and Electromagnetic Fields in Klein-Gordon-Maxwell Theory in Cosmological Inertial Frame

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ABSTRACT

We found equations of complex scalar fields and electromagnetic fields on interaction of complex scalar fields and electromagnetic fields in Klein-Gordon-Maxwell theory from Type A of wave function and Type B of expanded distance in cosmological inertial frame.

PACS Number:03.30.+p,03.65

Key words: Klein-Gordon-Maxwell Theory;

Cosmological Inertial Frame;

Complex Scalar fields;

Electromagnetic fields

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1. Introduction

The Lagrangian L of complex scalar fields ϕ, ϕ^* and Electromagnetic fields $F^{\mu\nu}, F_{\mu\nu}$ is Klein-Gordon-Maxwell theory in special relativity theory,

$$L = (\partial_\mu \phi + ieA_\mu \phi)(\partial^\mu \phi^* - ieA^\mu \phi^*) - \frac{m^2 c^2}{\hbar^2} \phi \phi^* - \frac{1}{4} F^{\mu\nu} F_{\mu\nu}$$

ϕ^* is ϕ 's adjoint scalar, m is the mass of scalar fields ϕ, ϕ^*

$$F^{\mu\nu} = \partial^\mu A^\nu - \partial^\nu A^\mu, F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \quad (1)$$

2. Equations of Interaction of Complex Scalar Fields and Electromagnetic Fields in Cosmological Inertial Frame

The Lagrangian L of interaction of complex scalar fields and Electromagnetic fields is Klein-Gordon-Maxwell theory in cosmological inertial frame,

$$L = (\bar{\partial}_\mu \phi + ie\bar{A}_\mu' \phi)(\bar{\partial}^\mu \phi^* - ie\bar{A}^\mu' \phi^*) - \frac{m^2 c^2}{\hbar^2} \phi \phi^* - \frac{1}{4} F^{\mu\nu'} F_{\mu\nu'} \quad (2-1)$$

We consider Type A of wave function and Type B of expanded distance,[1],[2],[3],[4]

$$\text{Type A of wave function: } r \rightarrow r\sqrt{\Omega(t_0)}, \quad t \rightarrow \frac{t}{\sqrt{\Omega(t_0)}},$$

Type B of expanded distance: $r \rightarrow r\Omega(t_0), t \rightarrow t$

$$\bar{\partial}_\mu = (\sqrt{\Omega(t_0)} \frac{\partial}{c\partial t}, \frac{1}{\sqrt{\Omega(t_0)}} \vec{\nabla}), \bar{\partial}^\mu = (\sqrt{\Omega(t_0)} \frac{\partial}{c\partial t}, -\frac{1}{\sqrt{\Omega(t_0)}} \vec{\nabla})$$

$$\bar{A}_\mu' = (\phi, \bar{A}\Omega(t_0)), \bar{A}^\mu' = (\phi, -\bar{A}\Omega(t_0)), \bar{F}_{\mu\nu}' = F_{\mu\nu}\Omega(t_0), \bar{F}^{\mu\nu'} = F^{\mu\nu}\Omega(t_0)$$

$$t_0 \text{ is the cosmological time. } \Omega(t_0) \text{ is the expanding ratio of universe in the cosmological time } t_0. \quad (2-2)$$

Complex scalar field equations are in Klein-Gordon-Maxwell theory in cosmological inertial frame,

$$\bar{\partial}_\mu \left(\frac{\partial L}{\partial(\bar{\partial}_\mu \phi)} \right) - \frac{\partial L}{\partial \phi} = (\bar{\partial}_\mu - ie\bar{A}_\mu')(\bar{\partial}^\mu \phi^* - ie\bar{A}^\mu' \phi^*) + \frac{m^2 c^2}{\hbar^2} \phi^* = 0 \quad (3)$$

The other equation is in Klein-Gordon-Maxwell theory in cosmological inertial frame,

$$\bar{\partial}_\mu \left(\frac{\partial L}{\partial(\bar{\partial}_\mu \phi^*)} \right) - \frac{\partial L}{\partial \phi^*} = (\bar{\partial}^\mu + ie\bar{A}^\mu')(\bar{\partial}_\mu \phi + ie\bar{A}_\mu' \phi) + \frac{m^2 c^2}{\hbar^2} \phi = 0 \quad (4)$$

If operator $\bar{\partial}_\mu', \bar{\partial}^\mu'$ are in cosmological inertial frame,[1],[2],[3],[4]

$$\bar{\partial}_\mu' = \left(\frac{\partial}{c\partial t}, \frac{1}{\Omega(t_0)} \vec{\nabla} \right), \bar{\partial}^\mu' = \left(\frac{\partial}{c\partial t}, -\frac{1}{\Omega(t_0)} \vec{\nabla} \right)$$

$$\bar{F}^{\mu\nu}' = \bar{\partial}^\mu' \bar{A}^\nu' - \bar{\partial}^\nu' \bar{A}^\mu', \bar{F}_{\mu\nu}' = \bar{\partial}_\mu' \bar{A}_\nu' - \bar{\partial}_\nu' \bar{A}_\mu' \quad (5)$$

Electromagnetic field equations are in Klein-Gordon-Maxwell theory in cosmological inertial frame,

$$\begin{aligned} & \bar{\partial}_\nu' \left(\frac{\partial L}{\partial (\bar{\partial}_\nu' \bar{A}_\mu')} \right) - \frac{\partial L}{\partial \bar{A}_\mu'} \\ &= \frac{1}{4} \bar{\partial}_\nu' (\bar{\partial}^\mu' \bar{A}^\nu' - \bar{\partial}^\nu' \bar{A}^\mu') - ie\phi(\bar{\partial}^\mu' \phi^* - ie\bar{A}^\mu' \phi^*) + ie\phi^*(\bar{\partial}^\mu' \phi + ie\bar{A}^\mu' \phi) \\ &= \frac{1}{4} \bar{\partial}_\nu' \bar{F}^{\mu\nu}' - ie\phi(\bar{\partial}^\mu' \phi^* - ie\bar{A}^\mu' \phi^*) + ie\phi^*(\bar{\partial}^\mu' \phi + ie\bar{A}^\mu' \phi) = 0 \end{aligned} \quad (6)$$

Hence,[5],

$$\begin{aligned} \bar{\partial}_\nu' \bar{F}^{\mu\nu}' &= \frac{4\pi}{c} \bar{J}^\mu' = 4ie[\phi^* \bar{\partial}_\nu' \bar{A}_\mu' - ie\bar{A}_\nu' \phi^* - ie\bar{A}_\mu' \phi] \\ \bar{J}^\mu' &= \frac{c}{\pi} ie[\phi(\bar{\partial}^\mu' \phi^* - ie\bar{A}^\mu' \phi^*) - \phi^*(\bar{\partial}^\mu' \phi + ie\bar{A}^\mu' \phi)] \end{aligned} \quad (7)$$

The other equation is in Klein-Gordon-Maxwell theory in cosmological inertial frame,

$$\begin{aligned} & \bar{\partial}^\nu' \left(\frac{\partial L}{\partial (\bar{\partial}^\nu' \bar{A}^\mu)} \right) - \frac{\partial L}{\partial \bar{A}^\mu} \\ &= \frac{1}{4} \bar{\partial}^\nu' (\bar{\partial}_\mu' \bar{A}_\nu' - \bar{\partial}_\nu' \bar{A}_\mu') + ie\phi^*(\bar{\partial}_\mu' \phi + ie\bar{A}_\mu' \phi) - ie\phi(\bar{\partial}_\mu' \phi^* - ie\bar{A}_\mu' \phi^*) \\ &= \frac{1}{4} \bar{\partial}^\nu' \bar{F}_{\mu\nu}' + ie\phi^*(\bar{\partial}_\mu' \phi + ie\bar{A}_\mu' \phi) - ie\phi(\bar{\partial}_\mu' \phi^* - ie\bar{A}_\mu' \phi^*) = 0 \end{aligned} \quad (8)$$

Hence,[5],

$$\begin{aligned} \bar{\partial}^\nu' \bar{F}_{\mu\nu}' &= \frac{4\pi}{c} \bar{J}_\mu' = -4ie[\phi^*(\bar{\partial}_\mu' \phi + ie\bar{A}_\mu' \phi) - \phi(\bar{\partial}_\mu' \phi^* - ie\bar{A}_\mu' \phi^*)] \\ \bar{J}_\mu' &= -\frac{c}{\pi} ie[\phi^*(\bar{\partial}_\mu' \phi + ie\bar{A}_\mu' \phi) - \phi(\bar{\partial}_\mu' \phi^* - ie\bar{A}_\mu' \phi^*)] \\ &= \frac{c}{\pi} i \bar{\partial}_\mu' \phi^* - ie\bar{A}_\mu' \phi^* - ie\bar{A}_\mu' \phi + ie\bar{A}_\mu' \phi \end{aligned} \quad (9)$$

3. Conclusion

We found equations of complex scalar fields and electromagnetic fields on interaction of complex scalar fields and electromagnetic fields in Klein-Gordon-Maxwell theory from Type A of wave function and Type B of expanded distance in cosmological inertial frame.

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