

# Geodesic Curve of a Gravitational Plane Wave Pulse and Curve of Particle

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## Abstract

We consider a system of a gravitational wave coming from infinity that collides with a mass  $M$ . The metric of the system approaches the metric of a gravitational plane wave pulse as the mass of  $M$  goes to zero. The metric of the gravitational plane wave pulse having a specific form. We show there is a limiting curve of  $M$ , as mass and size of  $M$  go to zero, that is not a geodesic curve. We show conservation of energy-momentum does not hold and there is no solution to the Einstein field equations for this system.

## 1 Gravitational plane wave pulse metric

Define  $u = t - x$  and let the metric  $g_{\mu\nu}(u)$  be [1]

$$g_{00}(u) = -1 \tag{1}$$

$$g_{11}(u) = 1 \tag{2}$$

$$g_{22}(u) = [L(u)]^2 e^{2\beta(u)} \tag{3}$$

$$g_{33}(u) = [L(u)]^2 e^{-2\beta(u)} \tag{4}$$

$$g_{01}(u) = g_{02}(u) = g_{03}(u) = g_{12}(u) = g_{13}(u) = g_{23}(u) = 0 \tag{5}$$

having  $g_{\mu\nu}(u) = \eta_{\mu\nu}$  for  $u < 0$  and

$$\frac{d^2 L}{du^2}(u) + \left[ \frac{d\beta}{du}(u) \right]^2 L(u) = 0 \tag{6}$$

This metric will satisfy  $R_{\mu\nu} = 0$ . It is the metric of a gravitational plane wave pulse. Let  $L(0) = 1$  and  $\beta \neq 0$ . We then have by (6) that  $L(u)$  will decrease and become zero at some point  $u_0 > 0$ . Consequently  $g_{22}(u) > 0$  for  $u < u_0$ .

## 2 Proper Lorentz transformation

Consider a coordinate transformation from  $t, x, y, z$  to  $t', x', y', z'$  coordinates that is a composition of a rotation by  $\theta$  about the  $z$  axis followed by a boost by  $2 \cos \theta / (1 + \cos^2 \theta)$  in the  $x$  direction followed by a rotation by  $\theta + \pi$  about the  $z$  axis. For  $\theta/\pi$  not an integer this is a proper Lorentz transformation such that

$$t = t'(1 + 2 \cot^2 \theta) - 2x' \cot^2 \theta + 2y' \cot \theta \tag{7}$$

$$x = 2t' \cot^2 \theta + x'(1 - 2 \cot^2 \theta) + 2y' \cot \theta \tag{8}$$

$$y = 2t' \cot \theta - 2x' \cot \theta + y' \tag{9}$$

$$z = z' \tag{10}$$

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By (7) and (8) we have  $t - x = t' - x'$ . For (7)-(10) define the metric  $g'_{\mu\nu}(u)$  by

$$g'_{\mu\nu}(u) = \frac{\partial x^\alpha}{\partial x'^\mu} \frac{\partial x^\beta}{\partial x'^\nu} g_{\alpha\beta}(u) \quad (11)$$

hence for the metric (1)-(5) we get

$$g'_{00}(u) = -1 - 4[1 - g_{22}(u)] \cot^2 \theta \quad (12)$$

$$g'_{01}(u) = 4[1 - g_{22}(u)] \cot^2 \theta \quad (13)$$

$$g'_{11}(u) = 1 - 4[1 - g_{22}(u)] \cot^2 \theta \quad (14)$$

$$g'_{02}(u) = -g'_{12}(u) = -2[1 - g_{22}(u)] \cot \theta \quad (15)$$

$$g'_{22}(u) = g_{22}(u) \quad (16)$$

$$g'_{33}(u) = g_{33}(u) \quad g'_{03}(u) = g'_{13}(u) = g'_{23}(u) = 0 \quad (17)$$

The metric  $g'_{\mu\nu}(u)$  satisfying  $R_{\mu\nu} = 0$  and  $g'_{\mu\nu}(u) = \eta_{\mu\nu}$  for  $u < 0$  is then also the metric of a gravitational plane wave pulse.

### 3 Geodesic curve

The curve

$$t(\lambda) = (1 + 2 \cot^2 \theta) \lambda - 2 \cot^2 \theta \int_0^\lambda \frac{dw}{g_{22}(w)} \quad (18)$$

$$x(\lambda) = 2 \cot^2 \theta \lambda - 2 \cot^2 \theta \int_0^\lambda \frac{dw}{g_{22}(w)} \quad (19)$$

$$y(\lambda) = -2 \cot \theta \lambda + 2 \cot \theta \int_0^\lambda \frac{dw}{g_{22}(w)} \quad (20)$$

$$z(\lambda) = 0 \quad (21)$$

satisfies the geodesic equation for the metric  $g'_{\mu\nu}(u)$  and so is a geodesic curve. We have for  $\lambda < 0$ , since  $g_{22}(u) = 1$  for  $u < 0$ , that  $x(\lambda) = y(\lambda) = z(\lambda) = 0$ . Choose  $\theta$  so that  $\cot \theta \neq 0$ . We then have by (18), since the integral goes to positive infinity as  $\lambda \rightarrow u_0$ , that  $t(\lambda) \rightarrow -\infty$  as  $\lambda \rightarrow u_0$ .

### 4 Geodesic curve and curve of particle

Consider a system of a gravitational wave pulse coming from negative  $x$  infinity that collides at  $t = 0$  at the origin with a mass  $M$  initially at rest. Let  $\tilde{g}_{\mu\nu}(t, x, y, z)$  be the metric of the system. Require  $\tilde{g}_{\mu\nu}(t, x, y, z) \rightarrow g'_{\mu\nu}(t - x)$  as the mass of  $M$  goes to zero.

For  $t < 0$  as  $t$  gets closer and closer to zero the wave front of the gravitational wave pulse gets closer and closer to  $M$ . For  $t < 0$  the wave front makes no contact with  $M$  hence for  $t < 0$  there is no mass inside the wave.

Let  $M$  have small mass and size and assume the curve of  $M$  is approximately the geodesic curve (18)-(21). For  $\lambda > 0$  points of the geodesic curve are inside the wave. Now  $t(\lambda) \rightarrow -\infty$  as  $\lambda \rightarrow u_0 > 0$ . We can then conclude there is a point with negative  $t$  inside the wave where there is a mass. This contradicts the previous paragraph.

Since the curve of  $M$  is not approximately a geodesic curve we have the limiting curve, as mass and size go to zero, of  $M$  cannot be a geodesic curve.

## 5 Conservation of energy-momentum and existence of solution

At a point  $p$  choose coordinates  $t'', x'', y'', z''$  so that the metric  $\tilde{g}''_{\mu\nu}$  is the Minkowski metric and the first order partial derivatives of the metric are zero.

We have by conservation of energy and momentum that  $T''^{\mu\alpha}_{;\alpha} = 0$  or

$$\frac{\partial T''^{\mu\alpha}}{\partial x''^\alpha} + \Gamma''^{\mu}_{\alpha\beta} T''^{\alpha\beta} + \Gamma''^{\alpha}_{\alpha\beta} T''^{\mu\beta} = 0 \quad (22)$$

where  $T''^{\mu\nu}$  is the energy-momentum tensor of  $M$  in  $t'', x'', y'', z''$  coordinates and  $\Gamma''^{\mu}_{\alpha\beta}$  is calculated using  $\tilde{g}''_{\mu\nu}$ . We then have  $\Gamma''^{\mu}_{\alpha\beta}$  are zero at  $p$  so by (22) we have at  $p$  that  $\partial''_{\alpha} T''^{\mu\alpha} = 0$ . Since the curve  $M$  is not approximately a geodesic there is a  $p$  on the curve of  $M$  so that  $M$  experiences an acceleration in the  $t'', x'', y'', z''$  coordinate system and hence  $\partial''_{\alpha} T''^{\mu\alpha}$  is not zero at  $p$ . Consequently (22) does not hold. Also there is no solution to the Einstein field equations for this system since a solution would have  $T''^{\mu\alpha}_{;\alpha} = 0$ .

## 6 Conclusion

There is a limiting curve of a mass, that a gravitational wave pulse collides with, that is not a geodesic. We showed conservation of energy-momentum does not hold and there does not exist a solution to the Einstein field equations for this system.

## References

- [1] C. M. Misner, K.S. Thorne, and J. A. Wheeler, Gravitation p. 957 (W.H. Freeman, San Francisco, CA, 1973)
- [2] K. De Paepe, Physics Essays, June 2018