

Electro-Magnetic Field Equation and Lorentz gauge, Electro-Magnetic Wave Function, Wave Equation in Rindler spacetime

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ABSTRACT

In the general relativity theory, we find the electro-magnetic field transformation and the electro-magnetic field equation (Maxwell equation) in Rindler spacetime. We treat Lorentz gauge transformation in Rindler spacetime. We find the electro-magnetic wave equation and the electro-magnetic wave function in Rindler space-time. Specially, this article say the uniqueness of the accelerated frame because the accelerated frame can treat electro-magnetic field equation.

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1. Introduction

In the general relativity theory, our article's aim is that we find the electro-magnetic field equation in Rindler space-time.

The Rindler coordinate is

$$ct = \left(\frac{c^2}{a_0} + \xi^1 \right) \sinh\left(\frac{a_0 \xi^0}{c} \right)$$

$$x = \left(\frac{c^2}{a_0} + \xi^1 \right) \cosh\left(\frac{a_0 \xi^0}{c} \right) - \frac{c^2}{a_0}, y = \xi^2, z = \xi^3 \quad (1)$$

In this time, the tetrad θ^a_μ is

$$d\tau^2 = dt^2 - \frac{1}{c^2} [dx^2 + dy^2 + dz^2]$$

$$= -\frac{1}{c^2} \eta_{ab} \frac{\partial x^a}{\partial \xi^\mu} \frac{\partial x^b}{\partial \xi^\nu} d\xi^\mu d\xi^\nu$$

$$= -\frac{1}{c^2} \eta_{ab} \theta^a_\mu \theta^b_\nu d\xi^\mu d\xi^\nu = -\frac{1}{c^2} g_{\mu\nu} d\xi^\mu d\xi^\nu, \quad \theta^a_\mu = \frac{\partial x^a}{\partial \xi^\mu} \quad (2)$$

$$\theta^a_0(\xi^0) = \frac{\partial x^a}{\partial \xi^0} = ((1 + \frac{a_0}{c^2} \xi^1) \cosh(\frac{a_0 \xi^0}{c}), (1 + \frac{a_0}{c^2} \xi^1) \sinh(\frac{a_0 \xi^0}{c}), 0, 0) \quad (3)$$

About y -axis's and z -axis's orientation

$$\theta^a_2(\xi^0) = \frac{\partial x^a}{\partial \xi^2} = (0, 0, 1, 0), \quad \theta^a_3(\xi^0) = \frac{\partial x^a}{\partial \xi^3} = (0, 0, 0, 1) \quad (4)$$

The other unit vector $\theta^a_1(\xi^0)$ is

$$\theta^a_1(\xi^0) = \frac{\partial x^a}{\partial \xi^1} = (\sinh(\frac{a_0 \xi^0}{c}), \cosh(\frac{a_0 \xi^0}{c}), 0, 0) \quad (5)$$

Therefore,

$$\begin{aligned} cdt &= c \cosh(\frac{a_0 \xi^0}{c}) d\xi^0 (1 + \frac{a_0}{c^2} \xi^1) + \sinh(\frac{a_0 \xi^0}{c}) d\xi^1 \\ &= c \cosh(\frac{a_0 \xi^0}{c}) d\xi^0 + \sinh(\frac{a_0 \xi^0}{c}) d\xi^1 \\ dx &= c \sinh(\frac{a_0 \xi^0}{c}) d\xi^0 (1 + \frac{a_0}{c^2} \xi^1) + \cosh(\frac{a_0 \xi^0}{c}) d\xi^1 \end{aligned}$$

$$= c \sinh\left(\frac{a_0 \xi^0}{c}\right) d\xi^0 + \cosh\left(\frac{a_0 \xi^0}{c}\right) d\xi^1, dy = d\xi^2 = d\xi^2, dz = d\xi^3 = d\xi^3 \quad (6)$$

The vector transformation is

$$V^\mu = \frac{\partial x^\mu}{\partial \xi^\alpha} V^\alpha, \quad U_\mu = \frac{\partial x^\alpha}{\partial \xi^\mu} U_\alpha \quad (7)$$

Therefore, the transformation of the electro-magnetic 4-vector potential $(\phi, \vec{A}) = A^\alpha$ is

$$\begin{aligned} A^\alpha &= \frac{\partial x^\alpha}{\partial \xi^\mu} A^\mu = \frac{\partial x^\alpha}{\partial \xi^\mu} A_\xi^\mu = e^\alpha_\mu A_\xi^\mu, \quad e^\alpha_\mu = \frac{\partial x^\alpha}{\partial \xi^\mu} \\ dx^\alpha &= \frac{\partial x^\alpha}{\partial \xi^\mu} d\xi^\mu = \frac{\partial x^\alpha}{\partial \xi^\mu} d\xi^\mu = e^\alpha_\mu d\xi^\mu, \quad e^\alpha_\mu = \frac{\partial x^\alpha}{\partial \xi^\mu} \end{aligned} \quad (8)$$

Hence, the transformation of the electro-magnetic 4-vector potential (ϕ, \vec{A}) in inertial frame and the

electro-magnetic 4-vector potential (ϕ_ξ, \vec{A}_ξ) in uniformly accelerated frame is

$$\begin{aligned} \left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2\right) \phi &= 4\pi\rho \\ \left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2\right) \vec{A} &= \frac{4\pi}{c} \vec{j} \\ \text{4-vector } (c\rho, \vec{j}) &= \rho_0 \frac{dx^\alpha}{d\tau} \end{aligned} \quad (9)$$

Lorentz gauge transformation is in Rindler spacetime,

$$A^\mu \rightarrow A^\mu + \partial^\mu \Lambda = A^\mu + g^{\mu\nu} \partial_\nu \Lambda, \quad \Lambda \text{ is a scalar function.}$$

$$g^{00} = -\frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2}, \quad g^{11} = g^{22} = g^{33} = 1$$

Hence,

$$\phi_\xi \rightarrow \phi_\xi - \frac{1}{c} \frac{\partial \Lambda}{\partial \xi^0} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2}, \quad \vec{A}_\xi \rightarrow \vec{A}_\xi + \vec{\nabla}_\xi \Lambda, \quad \Lambda \text{ is a scalar function.} \quad (10)$$

Lorentz gauge fixing condition is in Rindler spacetime,

$$A^\mu_{;\mu} = \frac{\partial A^\mu}{\partial x^\mu} + \Gamma^\mu_{\mu\rho} A^\rho,$$

$$\begin{aligned}\Gamma^\mu_{\mu\rho} &= \frac{1}{2} g^{\mu\sigma} \left(\frac{\partial g_{\rho\sigma}}{\partial x^\mu} + \frac{\partial g_{\mu\sigma}}{\partial x^\rho} - \frac{\partial g_{\rho\mu}}{\partial x^\sigma} \right) = \frac{1}{2} g^{\mu\mu} \frac{\partial g_{\mu\mu}}{\partial x^\mu} = \Gamma^\mu_{\mu\mu}, \quad \rho = \sigma = \mu \\ g^{00} &= -\frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2}, \quad g^{11} = g^{22} = g^{33} = 1\end{aligned}\tag{11}$$

Hence, Lorentz gauge fixing condition is in Rindler spacetime,

$$\Gamma^\mu_{\mu\rho} = \frac{1}{2} g^{\mu\sigma} \left(\frac{\partial g_{\rho\sigma}}{\partial x^\mu} + \frac{\partial g_{\mu\sigma}}{\partial x^\rho} - \frac{\partial g_{\rho\mu}}{\partial x^\sigma} \right) = \frac{1}{2} g^{\mu\mu} \frac{\partial g_{\mu\mu}}{\partial x^\mu} = \Gamma^\mu_{\mu\mu} = 0$$

$$\begin{aligned}A^\mu_{;\mu} &= \frac{\partial A^\mu}{\partial x^\mu} + \Gamma^\mu_{\mu\rho} A^\rho = \frac{\partial A^\mu}{\partial x^\mu} \\ \frac{\partial A^\mu}{\partial x^\mu} &\rightarrow \frac{\partial}{\partial x^\mu} (A^\mu + g^{\mu\nu} \partial_\nu \Lambda) = \frac{\partial A^\mu}{\partial x^\mu} + g^{\mu\nu} \partial_\mu \partial_\nu \Lambda\end{aligned}$$

$$0 = \frac{1}{c} \frac{\partial \phi_\xi}{\partial \xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi$$

$$0 = \frac{\partial \phi}{\partial t} + \vec{\nabla} \cdot \vec{A} = \frac{\partial \phi_\xi}{\partial \xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi + \frac{A_{\xi^1} a_0}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})}$$

Hence, $A_{\xi^1} = 0$

$$\begin{aligned}0 &= \frac{1}{c} \frac{\partial \phi_\xi}{\partial \xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi \rightarrow \frac{1}{c} \frac{\partial \phi_\xi}{\partial \xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi - [\frac{1}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2} (\frac{\partial}{\partial \xi^0})^2 - \nabla_\xi^2] \Lambda = 0 \\ &[\frac{1}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2} (\frac{\partial}{\partial \xi^0})^2 - \nabla_\xi^2] \Lambda = 0\end{aligned}\tag{12}$$

Hence,

$$\phi = \cosh(\frac{a_0 \xi^0}{c}) (1 + \frac{a_0}{c^2} \xi^1) \phi_\xi = \cosh(\frac{a_0 \xi^0}{c}) \hat{\phi}_\xi$$

$$A_x = \sinh(\frac{a_0 \xi^0}{c}) (1 + \frac{a_0}{c^2} \xi^1) \phi_\xi = \sinh(\frac{a_0 \xi^0}{c}) \hat{\phi}_\xi$$

$$A_y = A_{\xi^2} = \hat{A}_{\xi^2}, A_z = A_{\xi^3} = \hat{A}_{\xi^3}\tag{13}$$

$$g = \begin{pmatrix} -(1 + \frac{a_0 \xi^1}{c^2})^2 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \eta = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$e^a{}_\mu e_b{}^\mu = \delta^a{}_b, \quad e^a{}_\mu e_a{}^\nu = \delta_\mu{}^\nu$$

$$e^a{}_\mu e^b{}_\nu \eta_{ab} = g_{\mu\nu} \rightarrow A^T \eta A = g$$

$$e_a{}^\mu e_b{}^\nu g_{\mu\nu} = \eta_{ab} \rightarrow (A^T)^{-1} g A^{-1} = (A^T)^{-1} A^T \eta A A^{-1} = \eta$$

$$e^a{}_\mu = \eta^{ab} g_{\mu\nu} e_b{}^\nu \rightarrow \eta^{-1} (A^T)^{-1} A^T \eta A = A = \eta^{-1} (A^T)^{-1} g \quad (14)$$

$$\begin{pmatrix} cd t \\ dx \\ dy \\ dz \end{pmatrix} = \begin{pmatrix} \cosh(\frac{a_0 \xi^0}{c})(1 + \frac{a_0 \xi^1}{c^2}) & \sinh(\frac{a_0 \xi^0}{c}) & 0 & 0 \\ \sinh(\frac{a_0 \xi^0}{c})(1 + \frac{a_0 \xi^1}{c^2}) & \cosh(\frac{a_0 \xi^0}{c}) & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} cd\xi^0 \\ d\xi^1 \\ d\xi^2 \\ d\xi^3 \end{pmatrix}$$

$$= A \begin{pmatrix} cd\xi^0 \\ d\xi^1 \\ d\xi^2 \\ d\xi^3 \end{pmatrix}$$

$$= \begin{pmatrix} \cosh(\frac{a_0 \xi^0}{c}) & \sinh(\frac{a_0 \xi^0}{c}) & 0 & 0 \\ \sinh(\frac{a_0 \xi^0}{c}) & \cosh(\frac{a_0 \xi^0}{c}) & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} cd\hat{\xi}^0 \\ d\hat{\xi}^1 \\ d\hat{\xi}^2 \\ d\hat{\xi}^3 \end{pmatrix}$$

(15)

$$\begin{aligned}
e_\mu^\alpha &= \frac{\partial \xi^\alpha}{\partial x^\mu} = A^{-1} = \left(\begin{array}{cccc} \frac{c \partial \xi^0}{\partial t} & \frac{c \partial \xi^0}{\partial x} & \frac{c \partial \xi^0}{\partial y} & \frac{c \partial \xi^0}{\partial z} \\ \frac{\partial \xi^1}{\partial t} & \frac{\partial \xi^1}{\partial x} & \frac{\partial \xi^1}{\partial y} & \frac{\partial \xi^1}{\partial z} \\ \frac{\partial \xi^2}{\partial t} & \frac{\partial \xi^2}{\partial x} & \frac{\partial \xi^2}{\partial y} & \frac{\partial \xi^2}{\partial z} \\ \frac{\partial \xi^3}{\partial t} & \frac{\partial \xi^3}{\partial x} & \frac{\partial \xi^3}{\partial y} & \frac{\partial \xi^3}{\partial z} \end{array} \right) \\
&= \left(\begin{array}{cccc} \frac{\cosh(\frac{a_0 \xi^0}{c})}{c} & -\frac{\sinh(\frac{a_0 \xi^0}{c})}{c} & 0 & 0 \\ \frac{(1 + \frac{a_0 \xi^1}{c^2})}{c^2} & \frac{(1 + \frac{a_0 \xi^1}{c^2})}{c^2} & 0 & 0 \\ -\frac{\sinh(\frac{a_0 \xi^0}{c})}{c} & \frac{\cosh(\frac{a_0 \xi^0}{c})}{c} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{array} \right) \quad (16) \\
\left(\begin{array}{c} \frac{1}{c} \frac{\partial}{\partial t} \\ \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \\ \frac{\partial}{\partial z} \end{array} \right) &= (A^{-1})^T \left(\begin{array}{c} \frac{1}{c} \frac{\partial}{\partial \xi^0} \\ \frac{\partial}{\partial \xi^1} \\ \frac{\partial}{\partial \xi^2} \\ \frac{\partial}{\partial \xi^3} \end{array} \right) = (A^T)^{-1} \left(\begin{array}{c} \frac{1}{c} \frac{\partial}{\partial \xi^0} \\ \frac{\partial}{\partial \xi^1} \\ \frac{\partial}{\partial \xi^2} \\ \frac{\partial}{\partial \xi^3} \end{array} \right) \\
&= \left(\begin{array}{cccc} \frac{\cosh(\frac{a_0 \xi^0}{c})}{c} & -\frac{\sinh(\frac{a_0 \xi^0}{c})}{c} & 0 & 0 \\ \frac{(1 + \frac{a_0 \xi^1}{c^2})}{c^2} & \frac{(1 + \frac{a_0 \xi^1}{c^2})}{c^2} & 0 & 0 \\ -\frac{\sinh(\frac{a_0 \xi^0}{c})}{c} & \frac{\cosh(\frac{a_0 \xi^0}{c})}{c} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{array} \right) \left(\begin{array}{c} \frac{1}{c} \frac{\partial}{\partial \xi^0} \\ \frac{\partial}{\partial \xi^1} \\ \frac{\partial}{\partial \xi^2} \\ \frac{\partial}{\partial \xi^3} \end{array} \right)
\end{aligned}$$

$$= \begin{pmatrix} \cosh(\frac{a_0 \xi^0}{c}) & -\sinh(\frac{a_0 \xi^0}{c}) & 0 & 0 \\ -\sinh(\frac{a_0 \xi^0}{c}) & \cosh(\frac{a_0 \xi^0}{c}) & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \frac{1}{c} \frac{\partial}{\partial \hat{\xi}^0} \\ \frac{\partial}{\partial \hat{\xi}^1} \\ \frac{\partial}{\partial \hat{\xi}^2} \\ \frac{\partial}{\partial \hat{\xi}^3} \end{pmatrix} \quad (17)$$

$$\frac{1}{c} \frac{\partial}{\partial t} = \frac{c \partial \xi^0}{c \partial t} \frac{1}{c} \frac{\partial}{\partial \xi^0} + \frac{\partial \xi^1}{c \partial t} \frac{\partial}{\partial \xi^1}$$

$$= \frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1}$$

$$= \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial}{c \partial \hat{\xi}^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \hat{\xi}^1}$$

$$\frac{\partial}{\partial x} = \frac{c \partial \xi^0}{\partial x} \frac{1}{c} \frac{\partial}{\partial \xi^0} + \frac{\partial \xi^1}{\partial x} \frac{\partial}{\partial \xi^1}$$

$$= - \frac{\sinh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} + \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1}$$

$$= - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{c \partial \hat{\xi}^0} + \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \hat{\xi}^1}$$

$$\frac{\partial}{\partial y} = \frac{\partial}{\partial \xi^2} = \frac{\partial}{\partial \hat{\xi}^2}, \quad \frac{\partial}{\partial z} = \frac{\partial}{\partial \xi^3} = \frac{\partial}{\partial \hat{\xi}^3}$$

$$\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2 = \frac{1}{c^2 (1 + \frac{a_0 \xi^1}{c^2})^2} (\frac{\partial}{\partial \xi^0})^2 - \nabla_{\xi}^2 = \frac{1}{c^2} (\frac{\partial}{\partial \hat{\xi}^0})^2 - \nabla_{\hat{\xi}}^2$$

$$\vec{\nabla} = (\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}), \quad \vec{\nabla}_{\xi} = (\frac{\partial}{\partial \xi^1}, \frac{\partial}{\partial \xi^2}, \frac{\partial}{\partial \xi^3}), \quad \vec{\nabla}_{\hat{\xi}} = (\frac{\partial}{\partial \hat{\xi}^1}, \frac{\partial}{\partial \hat{\xi}^2}, \frac{\partial}{\partial \hat{\xi}^3})$$

(18)

Hence, the transformation of the electro-magnetic 4-vector potential is

$$\begin{aligned}
\phi - \frac{1}{c} \frac{\partial \Lambda}{\partial t} &= \cosh\left(\frac{a_0 \xi^0}{c}\right) \left(1 + \frac{a_0}{c^2} \xi^1\right) \left\{ \phi_\xi - \frac{1}{c} \frac{\partial \Lambda}{\partial \xi^0} \frac{1}{\left(1 + \frac{a_0 \xi^1}{c^2}\right)^2} \right\} \\
&\quad + \sinh\left(\frac{a_0 \xi^0}{c}\right) \left(\frac{\partial \Lambda}{\partial \xi^1}\right) \\
A_x + \frac{\partial \Lambda}{\partial x} &= \sinh\left(\frac{a_0 \xi^0}{c}\right) \left(1 + \frac{a_0}{c^2} \xi^1\right) \left\{ \phi_\xi - \frac{1}{c} \frac{\partial \Lambda}{\partial \xi^0} \frac{1}{\left(1 + \frac{a_0 \xi^1}{c^2}\right)^2} \right\} \\
&\quad + \cosh\left(\frac{a_0 \xi^0}{c}\right) \left(\frac{\partial \Lambda}{\partial \xi^1}\right) \\
A_y + \frac{\partial \Lambda}{\partial y} &= A_{\xi^2} + \frac{\partial \Lambda}{\partial \xi^2}, A_z + \frac{\partial \Lambda}{\partial z} = A_{\xi^3} + \frac{\partial \Lambda}{\partial \xi^3} \\
0 &= \frac{1}{c} \frac{\partial}{\partial t} \left(-\frac{1}{c} \frac{\partial \Lambda}{\partial t}\right) + \vec{\nabla} \cdot (\vec{\nabla} \Lambda) = -\left[\frac{1}{c^2} \frac{1}{\left(1 + \frac{a_0 \xi^1}{c^2}\right)^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 - \nabla_\xi^2\right] \Lambda \tag{19}
\end{aligned}$$

2. Electro-magnetic Field in the Rindler space-time

The electro-magnetic field (\vec{E}, \vec{B}) is in the inertial frame,

$$\begin{aligned}
\vec{E} &= -\vec{\nabla} \phi - \frac{\partial \vec{A}}{\partial t}, \vec{B} = \vec{\nabla} \times \vec{A} \\
E_x &= -\frac{\partial \phi}{\partial x} - \frac{\partial A_x}{\partial t} \\
&= -\left[-\frac{\sinh\left(\frac{a_0 \xi^0}{c}\right)}{\left(1 + \frac{a_0 \xi^1}{c^2}\right)} \frac{\partial}{\partial \xi^0} + \cosh\left(\frac{a_0 \xi^0}{c}\right) \frac{\partial}{\partial \xi^1}\right] \cdot \left[\cosh\left(\frac{a_0 \xi^0}{c}\right) \left(1 + \frac{a_0 \xi^1}{c^2}\right) \phi_\xi\right] \\
&\quad - \left[-\frac{\cosh\left(\frac{a_0 \xi^0}{c}\right)}{\left(1 + \frac{a_0 \xi^1}{c^2}\right)} \frac{\partial}{\partial \xi^0} - \sinh\left(\frac{a_0 \xi^0}{c}\right) \frac{\partial}{\partial \xi^1}\right] \cdot \left[\sinh\left(\frac{a_0 \xi^0}{c}\right) \left(1 + \frac{a_0 \xi^1}{c^2}\right) \phi_\xi\right] \\
&= -\left(1 + \frac{a_0 \xi^1}{c^2}\right) \frac{\partial \phi_\xi}{\partial \xi^1} - 2\phi_\xi \frac{a_0}{c^2} \tag{20}
\end{aligned}$$

$$\begin{aligned}
&= -\frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{\partial}{\partial \xi^1} [(1+\frac{a_0}{c^2}\xi^1)^2 \phi_\xi] \\
&= -\frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{\partial}{\partial \hat{\xi}^1} [(1+\frac{a_0}{c^2}\xi^1) \hat{\phi}_\xi]
\end{aligned} \tag{21}$$

$$\begin{aligned}
E_y &= -\frac{\partial \phi}{\partial y} - \frac{\partial A_y}{c \partial t} = -\frac{\partial}{\partial \xi^2} [\cosh(\frac{a_0 \xi^0}{c})(1+\frac{a_0}{c^2}\xi^1)\phi_\xi] \\
&\quad - [\frac{\cosh(\frac{a_0 \xi^0}{c})}{(1+\frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1}] A_{\xi^2} \\
&= -(1+\frac{a_0 \xi^1}{c^2}) \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial \phi_\xi}{\partial \xi^2} - \frac{1}{(1+\frac{a_0 \xi^1}{c^2})} \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial A_{\xi^2}}{c \partial \xi^0} \\
&\quad + \sinh(\frac{a_0}{c} \xi^0) [\frac{\partial A_{\xi^2}}{\partial \xi^1}] \\
&= \cosh(\frac{a_0}{c} \xi^0) [-\frac{1}{(1+\frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^2} [\phi_\xi (1+\frac{a_0 \xi^1}{c^2})^2] - \frac{1}{(1+\frac{a_0 \xi^1}{c^2})} \frac{\partial A_{\xi^2}}{c \partial \xi^0}] \\
&\quad + \sinh(\frac{a_0}{c} \xi^0) [\frac{\partial A_{\xi^2}}{\partial \xi^1}] \\
&= \cosh(\frac{a_0}{c} \xi^0) [-\frac{1}{(1+\frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \hat{\xi}^2} [\hat{\phi}_\xi (1+\frac{a_0 \xi^1}{c^2})] - \frac{\partial \hat{A}_{\xi^2}}{c \partial \hat{\xi}^0}] \\
&\quad + \sinh(\frac{a_0}{c} \xi^0) [\frac{\partial \hat{A}_{\xi^2}}{\partial \hat{\xi}^1}]
\end{aligned} \tag{22}$$

$$E_z = -\frac{\partial \phi}{\partial z} - \frac{\partial A_z}{c \partial t} = -\frac{\partial}{\partial \xi^3} [\cosh(\frac{a_0 \xi^0}{c})(1+\frac{a_0}{c^2}\xi^1)\phi_\xi + \sinh(\frac{a_0 \xi^0}{c})A_{\xi^1}]$$

$$\begin{aligned}
& - \left[\frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1} \right] A_{\xi^3} \\
& = -(1 + \frac{a_0 \xi^1}{c^2}) \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial \phi_\xi}{\partial \xi^3} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial A_{\xi^3}}{\partial \xi^0} \\
& \quad + \sinh(\frac{a_0}{c} \xi^0) \left[\frac{\partial A_{\xi^3}}{\partial \xi^1} \right] \\
& = \cosh(\frac{a_0}{c} \xi^0) \left[- \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^3} [\phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2] - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial A_{\xi^3}}{\partial \xi^0} \right. \\
& \quad \left. + \sinh(\frac{a_0}{c} \xi^0) \left[\frac{\partial A_{\xi^3}}{\partial \xi^1} \right] \right] \\
& = \cosh(\frac{a_0}{c} \xi^0) \left[- \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^3} [\hat{\phi}_\xi (1 + \frac{a_0 \xi^1}{c^2})] - \frac{\partial \hat{A}_{\xi^3}}{\partial \xi^0} \right] \\
& \quad + \sinh(\frac{a_0}{c} \xi^0) \left[\frac{\partial \hat{A}_{\xi^3}}{\partial \xi^1} \right]
\end{aligned} \tag{23}$$

$$B_x = \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} = \frac{\partial A_{\xi^3}}{\partial \xi^2} - \frac{\partial A_{\xi^2}}{\partial \xi^3} = \frac{\partial \hat{A}_{\xi^3}}{\partial \xi^2} - \frac{\partial \hat{A}_{\xi^2}}{\partial \xi^3} \tag{24}$$

$$B_y = \frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} = \frac{\partial A_x}{\partial \xi^3} - \frac{\partial A_{\xi^3}}{\partial x}$$

$$\begin{aligned}
& = \frac{\partial}{\partial \xi^3} \left[\sinh(\frac{a_0 \xi^0}{c}) (1 + \frac{a_0}{c^2} \xi^1) \phi_\xi \right] \\
& \quad - \left[- \frac{\sinh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^0} + \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1} \right] A_{\xi^3}
\end{aligned}$$

$$\begin{aligned}
&= \cosh\left(\frac{a_0}{c}\xi^0\right) \left[-\frac{\partial A_{\xi^3}}{\partial \xi^1} \right] \\
&\quad - \sinh\left(\frac{a_0}{c}\xi^0\right) \left[-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \frac{\partial}{\partial \xi^3} [\phi_\xi (1+\frac{a_0\xi^1}{c^2})^2] - \frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{\partial A_{\xi^3}}{c \partial \xi^0} \right] \\
&= \cosh\left(\frac{a_0}{c}\xi^0\right) \left[-\frac{\partial \hat{A}_{\xi^3}}{\partial \hat{\xi}^1} \right] \\
&\quad - \sinh\left(\frac{a_0}{c}\xi^0\right) \left[-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \frac{\partial}{\partial \hat{\xi}^3} [\hat{\phi}_\xi (1+\frac{a_0\xi^1}{c^2})] - \frac{\partial \hat{A}_{\xi^3}}{c \partial \hat{\xi}^0} \right] \quad (25)
\end{aligned}$$

$$\begin{aligned}
B_z &= \frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} = \frac{\partial A_{\xi^2}}{\partial x} - \frac{\partial A_x}{\partial \xi^2} \\
&= \left[-\frac{\sinh\left(\frac{a_0\xi^0}{c}\right)}{(1+\frac{a_0\xi^1}{c^2})} \frac{\partial}{\partial \xi^0} + \cosh\left(\frac{a_0\xi^0}{c}\right) \frac{\partial}{\partial \xi^1} \right] A_{\xi^3} \\
&\quad - \frac{\partial}{\partial \xi^2} \left[\sinh\left(\frac{a_0\xi^0}{c}\right) (1+\frac{a_0}{c^2}\xi^1) \phi_\xi \right] \\
&= \cosh\left(\frac{a_0}{c}\xi^0\right) \left[\frac{\partial A_{\xi^2}}{\partial \xi^1} \right] \\
&\quad + \sinh\left(\frac{a_0}{c}\xi^0\right) \left[-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \frac{\partial}{\partial \xi^2} [\phi_\xi (1+\frac{a_0\xi^1}{c^2})^2] - \frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{\partial A_{\xi^2}}{c \partial \xi^0} \right] \\
&= \cosh\left(\frac{a_0}{c}\xi^0\right) \left[\frac{\partial \hat{A}_{\xi^2}}{\partial \hat{\xi}^1} \right] \\
&\quad + \sinh\left(\frac{a_0}{c}\xi^0\right) \left[-\frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \frac{\partial}{\partial \hat{\xi}^2} [\hat{\phi}_\xi (1+\frac{a_0\xi^1}{c^2})] - \frac{\partial \hat{A}_{\xi^2}}{c \partial \hat{\xi}^0} \right] \quad (26)
\end{aligned}$$

Hence, we can define the electro-magnetic field $(\vec{E}_\xi, \vec{B}_\xi)$ in Rindler spacetime.

$$\vec{E}_\xi = -\frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi \{\phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2\} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial \vec{A}_\xi}{c \partial \xi^0}$$

$$= -\frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_{\hat{\xi}} \{\hat{\phi}_\xi (1 + \frac{a_0 \xi^1}{c^2})\} - \frac{\partial \vec{A}_\xi}{c \partial \hat{\xi}^0}$$

$$\vec{B}_\xi = \vec{\nabla}_\xi \times \vec{A}_\xi = \vec{\nabla}_{\hat{\xi}} \times \vec{A}_{\hat{\xi}}$$

In this time, $A_{\xi^1} = 0$

$$\begin{aligned} \vec{\nabla}_\xi &= (\frac{\partial}{\partial \xi^1}, \frac{\partial}{\partial \xi^2}, \frac{\partial}{\partial \xi^3}), \vec{A}_\xi = (0, A_{\xi^2}, A_{\xi^3}) \\ \vec{\nabla}_{\hat{\xi}} &= (\frac{\partial}{\partial \hat{\xi}^1}, \frac{\partial}{\partial \hat{\xi}^2}, \frac{\partial}{\partial \hat{\xi}^3}), \vec{A}_{\hat{\xi}} = (0, \hat{A}_{\hat{\xi}^2}, \hat{A}_{\hat{\xi}^3}) \end{aligned} \quad (27)$$

Lorentz gauge transformation is in Rindler spacetime,

$$\phi_\xi \rightarrow \phi_\xi - \frac{1}{c} \frac{\partial \Lambda}{\partial \xi^0} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2}, \vec{A}_\xi \rightarrow \vec{A}_\xi + \vec{\nabla}_\xi \Lambda, \Lambda \text{ is a scalar function.} \quad (28)$$

$$\begin{aligned} \vec{E}_\xi &= -\frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi \{\phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2\} + \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi \frac{\partial \Lambda}{c \partial \xi^0} \\ &\quad - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial \vec{A}_\xi}{c \partial \xi^0} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} \vec{\nabla}_\xi \Lambda \\ &= -\frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi \{\phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2\} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial \vec{A}_\xi}{c \partial \xi^0} \\ \vec{B}_\xi &= \vec{\nabla}_\xi \times \vec{A}_\xi + \vec{\nabla}_\xi \times \vec{\nabla}_\xi \Lambda = \vec{\nabla}_\xi \times \vec{A}_\xi \end{aligned} \quad (29)$$

We obtain the transformation of the electro-magnetic field.

$$\begin{aligned}
E_x &= -\frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{\partial}{\partial \xi^1} \{\phi_\xi (1+\frac{a_0\xi^1}{c^2})^2\} = -\frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{\partial}{\partial \xi^1} \{\hat{\phi}_\xi (1+\frac{a_0\xi^1}{c^2})\} = E_{\xi^1}, \\
E_y &= E_{\xi^2} \cosh(\frac{a_0\xi^0}{c}) + B_{\xi^3} \sinh(\frac{a_0\xi^0}{c}), \\
E_z &= E_{\xi^3} \cosh(\frac{a_0\xi^0}{c}) - B_{\xi^2} \sinh(\frac{a_0\xi^0}{c}) \\
B_x &= B_{\xi^1}, \\
B_y &= B_{\xi^2} \cosh(\frac{a_0\xi^0}{c}) - E_{\xi^3} \sinh(\frac{a_0\xi^0}{c}) \\
B_z &= B_{\xi^3} \cosh(\frac{a_0\xi^0}{c}) + E_{\xi^2} \sinh(\frac{a_0\xi^0}{c})
\end{aligned} \tag{30}$$

Hence,

$$E_x = E_{\xi^1}, B_x = B_{\xi^1},$$

$$\begin{pmatrix} E_y \\ B_y \\ E_z \\ B_z \end{pmatrix} = H \begin{pmatrix} E_{\xi^2} \\ B_{\xi^2} \\ E_{\xi^3} \\ B_{\xi^3} \end{pmatrix}$$

$$H = \begin{pmatrix} \cosh(\frac{a_0\xi^0}{c}) & 0 & 0 & \sinh(\frac{a_0\xi^0}{c}) \\ 0 & \cosh(\frac{a_0\xi^0}{c}) & -\sinh(\frac{a_0\xi^0}{c}) & 0 \\ 0 & -\sinh(\frac{a_0\xi^0}{c}) & \cosh(\frac{a_0\xi^0}{c}) & 0 \\ \sinh(\frac{a_0\xi^0}{c}) & 0 & 0 & \cosh(\frac{a_0\xi^0}{c}) \end{pmatrix} \tag{31}$$

The inverse-transformation of the electro-magnetic field is

$$E_{\xi^1} = E_x, B_{\xi^1} = B_x$$

$$\begin{pmatrix} E_{\xi^2} \\ B_{\xi^2} \\ E_{\xi^3} \\ B_{\xi^3} \end{pmatrix} = H^{-1} \begin{pmatrix} E_y \\ B_y \\ E_z \\ B_z \end{pmatrix}$$

$$H^{-1} = \begin{pmatrix} \cosh(\frac{a_0 \xi^0}{c}) & 0 & 0 & -\sinh(\frac{a_0 \xi^0}{c}) \\ 0 & \cosh(\frac{a_0 \xi^0}{c}) & \sinh(\frac{a_0 \xi^0}{c}) & 0 \\ 0 & \sinh(\frac{a_0 \xi^0}{c}) & \cosh(\frac{a_0 \xi^0}{c}) & 0 \\ -\sinh(\frac{a_0 \xi^0}{c}) & 0 & 0 & \cosh(\frac{a_0 \xi^0}{c}) \end{pmatrix} \quad (32)$$

$$E_{\xi^1} = E_x, B_{\xi^1} = B_x$$

$$\begin{aligned} E_{\xi^2} &= E_y \cosh(\frac{a_0 \xi^0}{c}) - B_z \sinh(\frac{a_0 \xi^0}{c}), \\ B_{\xi^2} &= B_y \cosh(\frac{a_0 \xi^0}{c}) + E_z \sinh(\frac{a_0 \xi^0}{c}) \\ E_{\xi^3} &= E_z \cosh(\frac{a_0 \xi^0}{c}) + B_y \sinh(\frac{a_0 \xi^0}{c}) \\ B_{\xi^3} &= B_z \cosh(\frac{a_0 \xi^0}{c}) - E_y \sinh(\frac{a_0 \xi^0}{c}) \end{aligned} \quad (33)$$

3. Electro-magnetic Field Equation(Maxwell Equation) in the Rindler space-time

Maxwell equation is

$$\vec{\nabla} \cdot \vec{E} = 4\pi\rho \quad (34-i)$$

$$\vec{\nabla} \times \vec{B} = \frac{\partial \vec{E}}{c \partial t} + \frac{4\pi}{c} \vec{j} \quad (34-ii)$$

$$\vec{\nabla} \cdot \vec{B} = 0 \quad (34-iii)$$

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{c \partial t} \quad (34-iv)$$

$$1. \vec{\nabla} \cdot \vec{E} = 4\pi\rho$$

$$E_x = E_{\xi^1},$$

$$\begin{aligned}
E_y &= E_{\xi^2} \cosh\left(\frac{a_0 \xi^0}{c}\right) + B_{\xi^3} \sinh\left(\frac{a_0 \xi^0}{c}\right), \\
E_z &= E_{\xi^3} \cosh\left(\frac{a_0 \xi^0}{c}\right) - B_{\xi^2} \sinh\left(\frac{a_0 \xi^0}{c}\right) \\
4\pi\rho &= \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z} \\
&= \left[-\frac{\sinh\left(\frac{a_0 \xi^0}{c}\right)}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^0} + \cosh\left(\frac{a_0 \xi^0}{c}\right) \frac{\partial}{\partial \xi^1} \right] E_{\xi^1} \\
&\quad + \frac{\partial}{\partial \xi^2} [E_{\xi^2} \cosh\left(\frac{a_0 \xi^0}{c}\right) + B_{\xi^3} \sinh\left(\frac{a_0 \xi^0}{c}\right)] \\
&\quad + \frac{\partial}{\partial \xi^3} [E_{\xi^3} \cosh\left(\frac{a_0 \xi^0}{c}\right) - B_{\xi^2} \sinh\left(\frac{a_0 \xi^0}{c}\right)] \\
&= \cosh\left(\frac{a_0}{c} \xi^0\right) (\vec{\nabla}_\xi \cdot \vec{E}_\xi) + \sinh\left(\frac{a_0}{c} \xi^0\right) \left[\frac{\partial B_{\xi^3}}{\partial \xi^2} - \frac{\partial B_{\xi^2}}{\partial \xi^3} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial E_{\xi^1}}{\partial \xi^0} \right] \\
&= \cosh\left(\frac{a_0}{c} \xi^0\right) (\vec{\nabla}_\xi \cdot \vec{E}_\xi) + \sinh\left(\frac{a_0}{c} \xi^0\right) \left[\frac{\partial B_{\hat{\xi}^3}}{\partial \hat{\xi}^2} - \frac{\partial B_{\hat{\xi}^2}}{\partial \hat{\xi}^3} - \frac{\partial E_{\xi^1}}{\partial \hat{\xi}^0} \right] \quad (35) \\
2 \cdot \vec{\nabla} \times \vec{B} &= \frac{\partial \vec{E}}{c \partial t} + \frac{4\pi}{c} \vec{j} \\
B_x &= B_{\xi^1} \\
B_y &= B_{\xi^2} \cosh\left(\frac{a_0 \xi^0}{c}\right) - E_{\xi^3} \sinh\left(\frac{a_0 \xi^0}{c}\right) \\
B_z &= B_{\xi^3} \cosh\left(\frac{a_0 \xi^0}{c}\right) + E_{\xi^2} \sinh\left(\frac{a_0 \xi^0}{c}\right) \\
\text{X-component} &= \frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z} \\
&= \frac{\partial}{\partial \xi^2} [B_{\xi^3} \cosh\left(\frac{a_0 \xi^0}{c}\right) + E_{\xi^2} \sinh\left(\frac{a_0 \xi^0}{c}\right)]
\end{aligned}$$

$$\begin{aligned}
& -\frac{\partial}{\partial \xi^3} [B_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) - E_{\xi^3} \sinh(\frac{a_0 \xi^0}{c})] \\
& = \cosh(\frac{a_0 \xi^0}{c}) [\frac{\partial B_{\xi^3}}{\partial \xi^2} - \frac{\partial B_{\xi^2}}{\partial \xi^3}] + \sinh(\frac{a_0 \xi^0}{c}) [\frac{\partial E_{\xi^2}}{\partial \xi^2} + \frac{\partial E_{\xi^3}}{\partial \xi^3}] \\
& = \frac{\partial E_x}{c \partial t} + \frac{4\pi}{c} j_x \\
& = \left[\frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1} \right] E_{\xi^1} + \frac{4\pi}{c} j_x
\end{aligned}$$

Hence,

$$\begin{aligned}
& \frac{4\pi}{c} j_x \\
& = \sinh(\frac{a_0 \xi^0}{c}) (\vec{\nabla}_{\xi} \cdot \vec{E}_{\xi}) + \cosh(\frac{a_0 \xi^0}{c}) [\frac{\partial B_{\xi^3}}{\partial \xi^2} - \frac{\partial B_{\xi^2}}{\partial \xi^3} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial E_{\xi^1}}{c \partial \xi^0}] \\
& = \sinh(\frac{a_0 \xi^0}{c}) (\vec{\nabla}_{\hat{\xi}} \cdot \vec{E}_{\xi}) + \cosh(\frac{a_0 \xi^0}{c}) [\frac{\partial B_{\xi^3}}{\partial \hat{\xi}^2} - \frac{\partial B_{\hat{\xi}^2}}{\partial \hat{\xi}^3} - \frac{\partial E_{\hat{\xi}^1}}{c \partial \hat{\xi}^0}] \quad (36)
\end{aligned}$$

$$\text{Y-component) } \frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x}$$

$$= \frac{\partial B_{\xi^1}}{\partial \xi^3}$$

$$\begin{aligned}
& - \left[-\frac{\sinh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} + \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1} \right] \cdot [B_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) + E_{\xi^2} \sinh(\frac{a_0 \xi^0}{c})] \\
& = \frac{\partial E_y}{c \partial t} + \frac{4\pi}{c} j_y \\
& = \left[\frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1} \right] \cdot [E_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) + B_{\xi^3} \sinh(\frac{a_0 \xi^0}{c})]
\end{aligned}$$

$$\begin{aligned}
& + \frac{4\pi}{c} j_y \\
\frac{4\pi}{c} j_y &= \frac{\partial B_{\xi^1}}{\partial \xi^3} - \frac{\partial B_{\xi^3}}{\partial \xi^1} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{a_0}{c^2} B_{\xi^3} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial E_{\xi^2}}{\partial \partial \xi^0} \\
&= \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^3} \{B_{\xi^1} (1 + \frac{a_0}{c^2} \xi^1)\} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^1} \{B_{\xi^3} (1 + \frac{a_0}{c^2} \xi^1)\} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial E_{\xi^2}}{\partial \partial \xi^0} \\
&= \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \hat{\xi}^3} \{B_{\xi^1} (1 + \frac{a_0}{c^2} \xi^1)\} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \hat{\xi}^1} \{B_{\xi^3} (1 + \frac{a_0}{c^2} \xi^1)\} - \frac{\partial E_{\xi^2}}{\partial \partial \hat{\xi}^0} \\
&\quad (37) \\
\text{Z-component)} \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y} &= \left[-\frac{\sinh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \partial \xi^0} + \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1} \right] \cdot [B_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) - E_{\xi^3} \sinh(\frac{a_0 \xi^0}{c})] \\
&\quad - \frac{\partial B_{\xi^1}}{\partial \xi^2} \\
&= \frac{\partial E_z}{\partial \partial t} + \frac{4\pi}{c} j_z \\
&= \left[\frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1} \right] \cdot [E_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) - B_{\xi^2} \sinh(\frac{a_0 \xi^0}{c})] \\
&\quad + \frac{4\pi}{c} j_z \\
\frac{4\pi}{c} j_z &= \frac{\partial B_{\xi^2}}{\partial \xi^1} - \frac{\partial B_{\xi^1}}{\partial \xi^2} + \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{a_0}{c^2} B_{\xi^2} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial E_{\xi^3}}{\partial \partial \xi^0} \\
&= \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^1} \{B_{\xi^2} (1 + \frac{a_0}{c^2} \xi^1)\} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^2} \{B_{\xi^1} (1 + \frac{a_0}{c^2} \xi^1)\} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial E_{\xi^3}}{\partial \partial \xi^0}
\end{aligned}$$

$$= \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \hat{\xi}^1} \{ B_{\xi^2} (1 + \frac{a_0}{c^2} \xi^1) \} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \hat{\xi}^2} \{ B_{\xi^1} (1 + \frac{a_0}{c^2} \xi^1) \} - \frac{\partial E_{\xi^3}}{c \partial \hat{\xi}^0}$$

(38)

3. $\vec{\nabla} \cdot \vec{B} = 0$

$$\begin{aligned} \vec{\nabla} \cdot \vec{B} &= \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} \\ &= \left[-\frac{\sinh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} + \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1} \right] B_{\xi^1} \\ &\quad + \frac{\partial}{\partial \xi^2} [B_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) - E_{\xi^3} \sinh(\frac{a_0 \xi^0}{c})] \\ &\quad + \frac{\partial}{\partial \xi^3} [B_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) + E_{\xi^2} \sinh(\frac{a_0 \xi^0}{c})] \\ &= \cosh(\frac{a_0 \xi^0}{c}) (\vec{\nabla}_{\xi} \cdot \vec{B}_{\xi}) + \sinh(\frac{a_0 \xi^0}{c}) \left[-\left(-\frac{\partial E_{\xi^2}}{\partial \xi^3} + \frac{\partial E_{\xi^3}}{\partial \xi^2} \right) - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial B_{\xi^1}}{c \partial \xi^0} \right] = 0 \\ &= \cosh(\frac{a_0 \xi^0}{c}) (\vec{\nabla}_{\hat{\xi}} \cdot \vec{B}_{\xi}) + \sinh(\frac{a_0 \xi^0}{c}) \left[-\left(-\frac{\partial E_{\xi^2}}{\partial \hat{\xi}^3} + \frac{\partial E_{\xi^3}}{\partial \hat{\xi}^2} \right) - \frac{\partial B_{\xi^1}}{c \partial \hat{\xi}^0} \right] = 0 \end{aligned}$$

(39)

4. $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{c \partial t}$

$$E_x = E_{\xi^1},$$

$$E_y = E_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) + B_{\xi^3} \sinh(\frac{a_0 \xi^0}{c}),$$

$$E_z = E_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) - B_{\xi^2} \sinh(\frac{a_0 \xi^0}{c})$$

$$\text{X-component) } \frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z}$$

$$= \frac{\partial}{\partial \xi^2} [E_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) - B_{\xi^2} \sinh(\frac{a_0 \xi^0}{c})]$$

$$\begin{aligned}
& -\frac{\partial}{\partial \xi^3} [E_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) + B_{\xi^3} \sinh(\frac{a_0 \xi^0}{c})] \\
& = \cosh(\frac{a_0}{c} \xi^0) [\frac{\partial E_{\xi^3}}{\partial \xi^2} - \frac{\partial E_{\xi^2}}{\partial \xi^3}] - \sinh(\frac{a_0 \xi^0}{c}) [\frac{\partial B_{\xi^2}}{\partial \xi^2} + \frac{\partial B_{\xi^3}}{\partial \xi^3}] \\
& = -\frac{\partial B_x}{c \partial t} \\
& = -[\frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1}] B_{\xi^1} \\
\text{Hence, } & -\sinh(\frac{a_0 \xi^0}{c}) (\vec{\nabla}_{\xi} \cdot \vec{B}_{\xi}) + c \circ s \frac{a_0 \xi^0}{c} [\frac{\partial E_{\xi^3}}{\partial \xi^2} - \frac{\partial E_{\xi^2}}{\partial \xi^3}] + \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial B_{\xi^1}}{c \partial \xi^0} \\
& = -\sinh(\frac{a_0 \xi^0}{c}) (\vec{\nabla}_{\xi} \cdot \vec{B}_{\xi}) + \cosh(\frac{a_0 \xi^0}{c}) [(\frac{\partial E_{\xi^3}}{\partial \xi^2} - \frac{\partial E_{\xi^2}}{\partial \xi^3}) + \frac{\partial B_{\xi^1}}{c \partial \xi^0}] = 0 \quad (40)
\end{aligned}$$

$$\begin{aligned}
\text{Y-component) } & \frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} \\
& = \frac{\partial E_{\xi^1}}{\partial \xi^3} \\
& - [-\frac{\sinh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} + \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1}] \cdot [E_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) - B_{\xi^2} \sinh(\frac{a_0 \xi^0}{c})] \\
& = -\frac{\partial B_y}{c \partial t} \\
& = -[\frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{c \partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1}] \cdot [B_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) - E_{\xi^3} \sinh(\frac{a_0 \xi^0}{c})] \\
& \frac{\partial E_{\xi^1}}{\partial \xi^3} - \frac{\partial E_{\xi^3}}{\partial \xi^1} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{a_0}{c^2} E_{\xi^3} + \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial B_{\xi^2}}{c \partial \xi^0}
\end{aligned}$$

$$\begin{aligned}
&= \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^3} \{E_{\xi^1} (1 + \frac{a_0}{c^2} \xi^1)\} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^1} \{E_{\xi^3} (1 + \frac{a_0 \xi^1}{c^2})\} + \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial B_{\xi^2}}{\partial \xi^0} \\
&= \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \hat{\xi}^3} \{E_{\xi^1} (1 + \frac{a_0}{c^2} \xi^1)\} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \hat{\xi}^1} \{E_{\xi^3} (1 + \frac{a_0 \xi^1}{c^2})\} + \frac{\partial B_{\xi^2}}{\partial \hat{\xi}^0} = 0
\end{aligned} \tag{41}$$

$$\begin{aligned}
&\text{Z-component) } \frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} \\
&= \left[-\frac{\sinh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^0} + \cosh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1} \right] \cdot [E_{\xi^2} \cosh(\frac{a_0 \xi^0}{c}) + B_{\xi^3} \sinh(\frac{a_0 \xi^0}{c})] \\
&\quad - \frac{\partial E_{\xi^1}}{\partial \xi^2} \\
&= -\frac{\partial B_z}{\partial \partial t} \\
&= -\left[\frac{\cosh(\frac{a_0 \xi^0}{c})}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^0} - \sinh(\frac{a_0 \xi^0}{c}) \frac{\partial}{\partial \xi^1} \right] \cdot [B_{\xi^3} \cosh(\frac{a_0 \xi^0}{c}) + E_{\xi^2} \sinh(\frac{a_0 \xi^0}{c})] \\
&\quad - \frac{\partial E_{\xi^2}}{\partial \xi^1} - \frac{\partial E_{\xi^1}}{\partial \xi^2} + \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{a_0}{c^2} E_{\xi^2} + \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial B_{\xi^3}}{\partial \xi^0} \\
&= \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^1} \{E_{\xi^2} (1 + \frac{a_0}{c^2} \xi^1)\} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \xi^2} \{E_{\xi^1} (1 + \frac{a_0 \xi^1}{c^2})\} + \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial B_{\xi^3}}{\partial \xi^0} \\
&= \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \hat{\xi}^1} \{E_{\xi^2} (1 + \frac{a_0}{c^2} \xi^1)\} - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{\partial}{\partial \hat{\xi}^2} \{E_{\xi^1} (1 + \frac{a_0 \xi^1}{c^2})\} + \frac{\partial B_{\xi^3}}{\partial \hat{\xi}^0} = 0
\end{aligned} \tag{42}$$

Therefore, we obtain the electro-magnetic field equation by Eq (35)-Eq(42) in Rindler spacetime .

$$\vec{\nabla}_{\xi} \cdot \vec{E}_{\xi} = \vec{\nabla}_{\hat{\xi}} \cdot \vec{E}_{\xi} = 4\pi\rho_{\xi} (1 + \frac{a_0 \xi^1}{c^2}) \tag{43-i}$$

$$\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi \times \{\vec{B}_\xi(1+\frac{a_0\xi^1}{c^2})\} = \frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial \vec{E}_\xi}{c\partial \xi^0} + \frac{4\pi}{c}\vec{j}_\xi$$

$$\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_{\hat{\xi}} \times \{\vec{B}_\xi(1+\frac{a_0\xi^1}{c^2})\} = \frac{\partial \vec{E}_\xi}{c\partial \hat{\xi}^0} + \frac{4\pi}{c}\vec{j}_\xi \quad (43\text{-ii})$$

$$\vec{\nabla}_\xi \cdot \vec{B}_\xi = \vec{\nabla}_{\hat{\xi}} \cdot \vec{B}_\xi = 0 \quad (43\text{-iii})$$

$$\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi \times \{\vec{E}_\xi(1+\frac{a_0\xi^1}{c^2})\} = -\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\frac{\partial \vec{B}_\xi}{c\partial \xi^0}$$

$$\frac{1}{(1+\frac{a_0\xi^1}{c^2})}\vec{\nabla}_\xi \times \{\vec{E}_\xi(1+\frac{a_0\xi^1}{c^2})\} = -\frac{\partial \vec{B}_\xi}{c\partial \hat{\xi}^0} \quad (43\text{-iv})$$

$$\vec{E}_\xi = (E_{\xi^1}, E_{\xi^2}, E_{\xi^3}), \vec{B}_\xi = (B_{\xi^1}, B_{\xi^2}, B_{\xi^3}),$$

$$\vec{\nabla}_\xi = (\frac{\partial}{\partial \xi^1}, \frac{\partial}{\partial \xi^2}, \frac{\partial}{\partial \xi^3}), \vec{\nabla}_{\hat{\xi}} = (\frac{\partial}{\partial \hat{\xi}^1}, \frac{\partial}{\partial \hat{\xi}^2}, \frac{\partial}{\partial \hat{\xi}^3})$$

Hence, the transformation of 4-vector $(c\rho, \vec{j}) = \rho_0 \frac{dx^\alpha}{d\tau}$ is

$$\rho = \rho_\xi(1+\frac{a_0\xi^1}{c^2})\cosh(\frac{a_0\xi^0}{c}) + \frac{j_\xi^1}{c}\sinh(\frac{a_0\xi^0}{c})$$

$$j_x = j_\xi^1 \cosh(\frac{a_0\xi^0}{c}) + c\rho_\xi(1+\frac{a_0}{c^2}\xi^1)\sinh(\frac{a_0\xi^0}{c}), \quad j_y = j_\xi^2, j_z = j_\xi^3$$

$$\text{In this time, 4-vector } (c\rho_\xi, \vec{j}_\xi) = \rho_0 \frac{d\xi^\alpha}{d\tau} \quad (44)$$

We treat Lorentz gauge transformation about the electro-magnetic field equation in Rindler spacetime.

Eq(43-i) is

$$\vec{\nabla}_\xi \cdot \vec{E}_\xi = \vec{\nabla}_\xi \cdot \left\{ -\frac{1}{(1+\frac{a_0\xi^1}{c^2})} \vec{\nabla}_\xi \{\phi_\xi(1+\frac{a_0\xi^1}{c^2})^2\} - \frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{\partial \vec{A}_\xi}{c\partial \xi^0} \right\}$$

$$\begin{aligned}
&= -\vec{\nabla}_\xi \left\{ \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \right\} \cdot [\vec{\nabla}_\xi \left\{ \phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2 \right\} + \frac{\partial \vec{A}_\xi}{\partial \xi^0}] \\
&\quad - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} [\nabla_\xi^2 \left\{ \phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2 \right\} + \frac{\partial}{\partial \xi^0} (\vec{\nabla}_\xi \cdot \vec{A}_\xi)] \\
&= \frac{a_0}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2} \left[\frac{\partial}{\partial \xi^1} \left\{ \phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2 \right\} \right] \\
&\quad - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} [\nabla_\xi^2 - \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} (\frac{\partial}{\partial \xi^0})^2] \left\{ \phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2 \right\} \\
&\quad - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^0} \left[\frac{1}{c} \frac{\partial \phi_\xi}{\partial \xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi \right], \frac{1}{c} \frac{\partial \phi_\xi}{\partial \xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi = 0 \\
&= -\frac{a_0}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} E_{\xi^1} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} [\nabla_\xi^2 - \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} (\frac{\partial}{\partial \xi^0})^2] \left\{ \phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2 \right\} \\
&= 4\pi \rho_\xi (1 + \frac{a_0 \xi^1}{c^2}) \tag{45}
\end{aligned}$$

If we apply Lorentz gauge transformation to Eq (45),

$$\begin{aligned}
\phi_\xi &\rightarrow \phi_\xi - \frac{1}{c} \frac{\partial \Lambda}{\partial \xi^0} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2}, \quad \vec{A}_\xi \rightarrow \vec{A}_\xi + \vec{\nabla}_\xi \Lambda, \quad \Lambda \text{ is a scalar function.} \\
&= -\frac{a_0}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} E_{\xi^1} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} [\nabla_\xi^2 - \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} (\frac{\partial}{\partial \xi^0})^2] \left\{ \phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2 \right\} \\
&\quad + \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} [\nabla_\xi^2 - \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} (\frac{\partial}{\partial \xi^0})^2] \frac{1}{c} \frac{\partial \Lambda}{\partial \xi^0} \\
&= -\frac{a_0}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} E_{\xi^1} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} [\nabla_\xi^2 - \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} (\frac{\partial}{\partial \xi^0})^2] \left\{ \phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2 \right\}
\end{aligned}$$

$$+ \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial}{\partial \xi^0} \left\{ [\nabla_\xi^2 - \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} (\frac{\partial}{\partial \xi^0})^2] \Lambda \right\} \quad (46)$$

In this time,

$$[\frac{1}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2} (\frac{\partial}{\partial \xi^0})^2 - \nabla_\xi^2] \Lambda = 0 \quad (47)$$

Hence, Eq(43-i) is

$$\begin{aligned} & \vec{\nabla}_\xi \cdot \vec{E}_\xi \\ &= -\frac{a_0}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} E_{\xi^1} - \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} [\nabla_\xi^2 - \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} (\frac{\partial}{\partial \xi^0})^2] \{ \phi_\xi (1 + \frac{a_0 \xi^1}{c^2})^2 \} \\ &= 4\pi \rho_\xi (1 + \frac{a_0 \xi^1}{c^2}) \end{aligned} \quad (48)$$

Eq(43-i) is invariant about Lorentz gauge transformation in Rindler spacetime.

Eq (43-ii) is

$$\begin{aligned} & \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi \times \{ \vec{B}_\xi (1 + \frac{a_0 \xi^1}{c^2}) \} \\ &= \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi \times \{ \vec{\nabla}_\xi \times \vec{A}_\xi (1 + \frac{a_0 \xi^1}{c^2}) \} \\ &= \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \vec{\nabla}_\xi (1 + \frac{a_0}{c^2} \xi^1) \times \{ \vec{\nabla}_\xi \times \vec{A}_\xi \} + \vec{\nabla}_\xi \times \vec{\nabla}_\xi \times \vec{A}_\xi \\ &= \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{a_0}{c^2} (1, 0, 0) \times \vec{B}_\xi + \{ -\nabla_\xi^2 \vec{A}_\xi + \vec{\nabla}_\xi (\vec{\nabla}_\xi \cdot \vec{A}_\xi) \} \\ &= \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{a_0}{c^2} (0, -B_{\xi^3}, B_{\xi^2}) + \{ -\nabla_\xi^2 \vec{A}_\xi + \vec{\nabla}_\xi (\vec{\nabla}_\xi \cdot \vec{A}_\xi) \} \\ &= \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} \frac{\partial \vec{E}_\xi}{\partial \xi^0} + \frac{4\pi}{c} \vec{j}_\xi \end{aligned}$$

$$\begin{aligned}
&= -\frac{1}{(1+\frac{a_0\xi^1}{c^2})^2} \frac{\partial}{c\partial\xi^0} [\vec{\nabla}_\xi \{\phi_\xi (1+\frac{a_0\xi^1}{c^2})^2\}] - \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2} \frac{1}{c^2} (\frac{\partial}{\partial\xi^0})^2 \vec{A}_\xi + \frac{4\pi\vec{j}_\xi}{c} \\
&= -\frac{\partial}{c\partial\xi^0} \vec{\nabla}_\xi \phi_\xi - \frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{2a_0}{c^2} \frac{\partial\phi_\xi}{c\partial\xi^0} (1,0,0) - \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2} \frac{1}{c^2} (\frac{\partial}{\partial\xi^0})^2 \vec{A}_\xi + \frac{4\pi\vec{j}_\xi}{c}
\end{aligned} \tag{49}$$

Therefore,

$$\begin{aligned}
&\frac{4\pi}{c} \vec{j}_\xi \\
&= \frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{a_0}{c^2} (0, -B_{\xi^3}, B_{\xi^2}) + \{-\nabla_\xi^2 \vec{A}_\xi + \vec{\nabla}_\xi (\vec{\nabla}_\xi \cdot \vec{A}_\xi)\} \\
&\quad + \frac{\partial}{c\partial\xi^0} \vec{\nabla}_\xi \phi_\xi + \frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{2a_0}{c^2} \frac{\partial\phi_\xi}{c\partial\xi^0} (1,0,0) + \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2} \frac{1}{c^2} (\frac{\partial}{\partial\xi^0})^2 \vec{A}_\xi \\
&= \frac{a_0}{c^2} \frac{1}{(1+\frac{a_0\xi^1}{c^2})} (0, -B_{\xi^3}, B_{\xi^2}) + \frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{2a_0}{c^2} \frac{\partial\phi_\xi}{c\partial\xi^0} (1,0,0) \\
&\quad + [-\nabla_\xi^2 + \frac{1}{c^2} \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2} (\frac{\partial}{\partial\xi^0})^2] \vec{A}_\xi + \vec{\nabla}_\xi \left[\frac{1}{c} \frac{\partial\phi_\xi}{\partial\xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi \right] \\
&\quad - \frac{1}{c} \frac{\partial\phi_\xi}{\partial\xi^0} + \vec{\nabla}_\xi \cdot \vec{A}_\xi = 0
\end{aligned} \tag{50}$$

$$\begin{aligned}
&\frac{4\pi}{c} \vec{j}_\xi \\
&= \frac{a_0}{c^2} \frac{1}{(1+\frac{a_0\xi^1}{c^2})} (0, -B_{\xi^3}, B_{\xi^2}) + \frac{1}{(1+\frac{a_0\xi^1}{c^2})} \frac{2a_0}{c^2} \frac{\partial\phi_\xi}{c\partial\xi^0} (1,0,0) \\
&\quad + [-\nabla_\xi^2 + \frac{1}{c^2} \frac{1}{(1+\frac{a_0\xi^1}{c^2})^2} (\frac{\partial}{\partial\xi^0})^2] \vec{A}_\xi
\end{aligned} \tag{51}$$

If we apply Lorentz gauge transformation to Eq (51),

$$\phi_\xi \rightarrow \phi_\xi - \frac{1}{c} \frac{\partial \Lambda}{\partial \xi^0} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2}, \quad \vec{A}_\xi \rightarrow \vec{A}_\xi + \vec{\nabla}_\xi \Lambda, \quad \Lambda \text{ is a scalar function.}$$

$$\begin{aligned}
& \frac{4\pi}{c} \vec{j}_\xi \\
&= \frac{a_0}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} (0, -B_{\xi^3}, B_{\xi^2}) + \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{2a_0}{c^2} \frac{\partial \phi_\xi}{c \partial \xi^0} (1, 0, 0) \\
&\quad - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^3} \frac{2a_0}{c^2} \frac{1}{c^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 \Lambda (1, 0, 0) \\
&+ [-\nabla_\xi^2 + \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \left(\frac{\partial}{\partial \xi^0}\right)^2] \vec{A}_\xi + [-\nabla_\xi^2 + \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \left(\frac{\partial}{\partial \xi^0}\right)^2] \vec{\nabla}_\xi \Lambda
\end{aligned} \tag{52}$$

In this time,

$$\begin{aligned}
& [\frac{1}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 - \nabla_\xi^2] \Lambda = 0 \\
0 &= \vec{\nabla}_\xi [\{-\nabla_\xi^2 + \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \left(\frac{\partial}{\partial \xi^0}\right)^2\} \Lambda] \\
&= \vec{\nabla}_\xi \{ \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \} \left(\frac{\partial}{\partial \xi^0}\right)^2 \Lambda + [-\nabla_\xi^2 + \frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \left(\frac{\partial}{\partial \xi^0}\right)^2] \vec{\nabla}_\xi \Lambda
\end{aligned} \tag{53}$$

Therefore,

$$\begin{aligned}
& \frac{4\pi}{c} \vec{j}_\xi \\
&= \frac{a_0}{c^2} \frac{1}{(1 + \frac{a_0 \xi^1}{c^2})} (0, -B_{\xi^3}, B_{\xi^2}) + \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)} \frac{2a_0}{c^2} \frac{\partial \phi_\xi}{c \partial \xi^0} (1, 0, 0) \\
&\quad - \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^3} \frac{2a_0}{c^2} \frac{1}{c^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 \Lambda (1, 0, 0)
\end{aligned}$$

$$\begin{aligned}
& + [-\nabla_\xi^2 + \frac{1}{c^2} \frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2} (\frac{\partial}{\partial\xi^0})^2] \vec{A}_\xi \\
& + \vec{\nabla}_\xi [\{-\nabla_\xi^2 + \frac{1}{c^2} \frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2} (\frac{\partial}{\partial\xi^0})^2\} \Lambda] - \vec{\nabla}_\xi \{\frac{1}{c^2} \frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2}\} (\frac{\partial}{\partial\xi^0})^2 \Lambda \\
& = \frac{a_0}{c^2} \frac{1}{(1+\frac{a_0}{c^2}\xi^1)} (0, -B_{\xi^3}, B_{\xi^2}) + \frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \frac{2a_0}{c^2} \frac{\partial\phi_\xi}{c\partial\xi^0} (1, 0, 0) \\
& - \frac{1}{(1+\frac{a_0}{c^2}\xi^1)^3} \frac{2a_0}{c^2} \frac{1}{c^2} (\frac{\partial}{\partial\xi^0})^2 \Lambda (1, 0, 0) \\
& + [-\nabla_\xi^2 + \frac{1}{c^2} \frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2} (\frac{\partial}{\partial\xi^0})^2] \vec{A}_\xi + \frac{1}{c^2} \frac{2}{(1+\frac{a_0}{c^2}\xi^1)^3} \frac{a_0}{c^2} (\frac{\partial}{\partial\xi^0})^2 \Lambda (1, 0, 0) \\
& = \frac{a_0}{c^2} \frac{1}{(1+\frac{a_0}{c^2}\xi^1)} (0, -B_{\xi^3}, B_{\xi^2}) + \frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \frac{2a_0}{c^2} \frac{\partial\phi_\xi}{c\partial\xi^0} (1, 0, 0) \\
& + [-\nabla_\xi^2 + \frac{1}{c^2} \frac{1}{(1+\frac{a_0}{c^2}\xi^1)^2} (\frac{\partial}{\partial\xi^0})^2] \vec{A}_\xi \tag{54}
\end{aligned}$$

Hence, Eq(43-ii) is invariant about Lorentz gauge transformation in Rindler spacetime.

Eq (43-iii) is

$$\vec{\nabla}_\xi \cdot \vec{B}_\xi = \vec{\nabla}_\xi \cdot (\vec{\nabla}_\xi \times \vec{A}_\xi + \vec{\nabla}_\xi \times \vec{\nabla}_\xi \Lambda) = \vec{\nabla}_\xi \times \vec{\nabla}_\xi \cdot \vec{A}_\xi = 0 \tag{55}$$

Eq (43-iv) is

$$\begin{aligned}
& \frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \vec{\nabla}_\xi \times \{\vec{E}_\xi (1+\frac{a_0}{c^2}\xi^1)\} \\
& = -\frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \vec{\nabla}_\xi \times [\vec{\nabla}_\xi \{\phi_\xi (1+\frac{a_0}{c^2}\xi^1)^2\} - \vec{\nabla}_\xi (\frac{\partial\Lambda}{c\partial\xi^0}) + \frac{\partial\vec{A}_\xi}{c\partial\xi^0} + \frac{\partial}{c\partial\xi^0} (\vec{\nabla}_\xi \Lambda)] \\
& = -\frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \vec{\nabla}_\xi \times \frac{\partial\vec{A}_\xi}{c\partial\xi^0} = -\frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \frac{\partial(\vec{\nabla}_\xi \times \vec{A}_\xi)}{c\partial\xi^0} = -\frac{1}{(1+\frac{a_0}{c^2}\xi^1)} \frac{\partial\vec{B}_\xi}{c\partial\xi^0} \tag{56}
\end{aligned}$$

Hence, Eq (43-iii), Eq (43-iv) are invariant about Lorentz gauge transformation in Rindler spacetime.

Hence, the electro-magnetic field equations(Maxwell Equations) in Rindler spacetime are invariant about Lorentz gauge transformation.

4. Electro-magnetic wave equation in Rindler space-time

The electro-magnetic wave function is

$$E_x = E_{x0} \sin \Phi, E_y = E_{y0} \sin \Phi, E_z = E_{z0} \sin \Phi$$

$$B_x = B_{x0} \sin \Phi, B_y = B_{y0} \sin \Phi, B_z = B_{z0} \sin \Phi$$

$$E_{\xi^1} = E_x, B_{\xi^1} = B_x$$

$$E_{\xi^1} = E_{x0} \sin \Phi', B_{\xi^1} = B_{x0} \sin \Phi'$$

$$\begin{aligned} E_{\xi^2} &= E_y \cosh\left(\frac{a_0 \xi^0}{c}\right) - B_z \sinh\left(\frac{a_0 \xi^0}{c}\right), \\ &= (E_{y0} \sin \Phi') \cosh\left(\frac{a_0 \xi^0}{c}\right) - (B_{z0} \sin \Phi') \sinh\left(\frac{a_0 \xi^0}{c}\right) \\ B_{\xi^2} &= B_y \cosh\left(\frac{a_0 \xi^0}{c}\right) + E_z \sinh\left(\frac{a_0 \xi^0}{c}\right) \\ &= (B_{y0} \sin \Phi') \cosh\left(\frac{a_0 \xi^0}{c}\right) + (E_{z0} \sin \Phi') \sinh\left(\frac{a_0 \xi^0}{c}\right) \\ E_{\xi^3} &= E_z \cosh\left(\frac{a_0 \xi^0}{c}\right) + B_y \sinh\left(\frac{a_0 \xi^0}{c}\right) \\ &= (E_{z0} \sin \Phi') \cosh\left(\frac{a_0 \xi^0}{c}\right) + (B_{y0} \sin \Phi') \sinh\left(\frac{a_0 \xi^0}{c}\right) \\ B_{\xi^3} &= B_z \cosh\left(\frac{a_0 \xi^0}{c}\right) - E_y \sinh\left(\frac{a_0 \xi^0}{c}\right) \\ &= (B_{z0} \sin \Phi') \cosh\left(\frac{a_0 \xi^0}{c}\right) - (E_{y0} \sin \Phi') \sinh\left(\frac{a_0 \xi^0}{c}\right) \end{aligned} \tag{57}$$

$$\Phi = \omega(t - l \frac{x}{c} - m \frac{y}{c} - n \frac{z}{c}),$$

$$\Phi' = \omega' (\hat{\xi}^0 - l \frac{\hat{\xi}^1}{c} - m \frac{\hat{\xi}^2}{c} - n \frac{\hat{\xi}^3}{c})$$

In this time,

$$ct = \left(\frac{c^2}{a_0} + \xi^1 \right) \sinh\left(\frac{a_0 \xi^0}{c}\right), \quad x = \left(\frac{c^2}{a_0} + \xi^1 \right) \cosh\left(\frac{a_0 \xi^0}{c}\right) - \frac{c^2}{a_0}$$

$$y = \xi^2, z = \xi^3$$

$$\xi^0 = \frac{c}{a_0} \tanh^{-1}\left(\frac{ct}{x + \frac{c^2}{a_0}}\right), \xi^1 = \sqrt{\left(x + \frac{c^2}{a_0}\right)^2 - c^2 t^2} - \frac{c^2}{a_0}$$

$$\lim_{a_0 \rightarrow 0} \xi^0 = \lim_{a_0 \rightarrow 0} c \tanh^{-1}\left(\frac{cta_0}{a_0 x + c^2}\right) / a_0 = \lim_{a_0 \rightarrow 0} c \tanh^{-1}\left(\frac{cta_0}{c^2}\right) / a_0 = \lim_{a_0 \rightarrow 0} c - \frac{1}{1 - \left(\frac{a_0 t}{c}\right)^2} \frac{t}{c} = t$$

$$\begin{aligned} \lim_{a_0 \rightarrow 0} \xi^1 &= \lim_{a_0 \rightarrow 0} c^2 \left(\sqrt{\left(1 + \frac{a_0}{c^2} x\right)^2 - \frac{a_0^2 t^2}{c^2}} - 1 \right) / a_0 = \lim_{a_0 \rightarrow 0} c^2 \left(\sqrt{\left(1 + \frac{a_0}{c^2} x\right)^2 - 1} \right) / a_0 \\ &= \lim_{a_0 \rightarrow 0} c^2 \left(\frac{a_0}{c^2} x \right) / a_0 = x \end{aligned}$$

Hence,

$$\lim_{a_0 \rightarrow 0} \hat{\xi}^0 = \lim_{a_0 \rightarrow 0} \int d\hat{\xi}^0 = \lim_{a_0 \rightarrow 0} \int \left(1 + \frac{a_0}{c^2} \xi^1\right) d\xi^0 = \lim_{a_0 \rightarrow 0} \int d\xi^0 = \lim_{a_0 \rightarrow 0} \xi^0 = t$$

$$\lim_{a_0 \rightarrow 0} \hat{\xi}^1 = \lim_{a_0 \rightarrow 0} \xi^1 = x, \quad y = \xi^2 = \hat{\xi}^2, \quad z = \xi^3 = \hat{\xi}^3$$

Therefore, electro-magnetic wave function is

$$\begin{aligned} \lim_{a_0 \rightarrow 0} \Phi' &= \lim_{a_0 \rightarrow 0} \omega' \left(\hat{\xi}^0 - l' \frac{\hat{\xi}^1}{c} - m' \frac{\hat{\xi}^2}{c} - n' \frac{\hat{\xi}^3}{c} \right) = \omega' \left(t - l' \frac{x}{c} - m' \frac{y}{c} - n' \frac{z}{c} \right) \\ &= \omega \left(t - l \frac{x}{c} - m \frac{y}{c} - n \frac{z}{c} \right) = \Phi \end{aligned}$$

$$\omega = \omega', l' = l, m' = m, n' = n$$

$$l'^2 + m'^2 + n'^2 = l'^2 + m'^2 + n'^2 = 1 \tag{58}$$

Hence,

$$\left[\frac{1}{c^2} \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)^2} \left(\frac{\partial}{\partial \xi^0} \right)^2 - \nabla_{\xi}^2 \right] E_{\xi^1}$$

$$= \left[\frac{1}{c^2} \left(\frac{\partial}{\partial \hat{\xi}^0} \right)^2 - \nabla_{\hat{\xi}}^2 \right] E_{\xi^1} = 0$$

$$\left[\frac{1}{c^2} \frac{1}{\left(1 + \frac{a_0}{c^2} \xi^1\right)^2} \left(\frac{\partial}{\partial \xi^0} \right)^2 - \nabla_{\xi}^2 \right] B_{\xi^1}$$

$$\begin{aligned}
&= \left[\frac{1}{c^2} \left(\frac{\partial}{\partial \hat{\xi}^0} \right)^2 - \nabla_{\hat{\xi}}^2 \right] B_{\xi^1} = 0 \\
&\quad \left[\frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \left(\frac{\partial}{\partial \xi^0} \right)^2 - \nabla_{\xi}^2 \right] E_y \\
&= \left[\frac{1}{c^2} \left(\frac{\partial}{\partial \hat{\xi}^0} \right)^2 - \nabla_{\hat{\xi}}^2 \right] E_y = 0 \\
&\quad \left[\frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \left(\frac{\partial}{\partial \xi^0} \right)^2 - \nabla_{\xi}^2 \right] B_y \\
&= \left[\frac{1}{c^2} \left(\frac{\partial}{\partial \hat{\xi}^0} \right)^2 - \nabla_{\hat{\xi}}^2 \right] B_y = 0 \\
&\quad \left[\frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \left(\frac{\partial}{\partial \xi^0} \right)^2 - \nabla_{\xi}^2 \right] E_z \\
&= \left[\frac{1}{c^2} \left(\frac{\partial}{\partial \hat{\xi}^0} \right)^2 - \nabla_{\hat{\xi}}^2 \right] E_z = 0 \\
&\quad \left[\frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \left(\frac{\partial}{\partial \xi^0} \right)^2 - \nabla_{\xi}^2 \right] B_z \\
&= \left[\frac{1}{c^2} \left(\frac{\partial}{\partial \hat{\xi}^0} \right)^2 - \nabla_{\hat{\xi}}^2 \right] B_z = 0 \tag{59}
\end{aligned}$$

The electro-magnetic wave equation is in vacuum

$$\begin{aligned}
&\vec{\nabla}_{\xi} \times \left(1 + \frac{a_0}{c^2} \xi^1 \right) \vec{\nabla}_{\xi} \times \{ \vec{E}_{\xi} \left(1 + \frac{a_0}{c^2} \xi^1 \right) \} \\
&= \vec{\nabla}_{\xi} \left(1 + \frac{a_0}{c^2} \xi^1 \right) \times \vec{\nabla}_{\xi} \times \{ \vec{E}_{\xi} \left(1 + \frac{a_0}{c^2} \xi^1 \right) \} + \left(1 + \frac{a_0}{c^2} \xi^1 \right) \vec{\nabla}_{\xi} \times \vec{\nabla}_{\xi} \times \{ \vec{E}_{\xi} \left(1 + \frac{a_0}{c^2} \xi^1 \right) \} \\
&= \vec{\nabla}_{\xi} \left(1 + \frac{a_0}{c^2} \xi^1 \right) \times \vec{\nabla}_{\xi} \left(1 + \frac{a_0}{c^2} \xi^1 \right) \times \vec{E}_{\xi} \\
&\quad + \left(1 + \frac{a_0}{c^2} \xi^1 \right) \vec{\nabla}_{\xi} \left(1 + \frac{a_0}{c^2} \xi^1 \right) \times \vec{\nabla}_{\xi} \times \vec{E}_{\xi} \\
&\quad + \left(1 + \frac{a_0}{c^2} \xi^1 \right) \vec{\nabla}_{\xi} \times \vec{\nabla}_{\xi} \left(1 + \frac{a_0}{c^2} \xi^1 \right) \times \vec{E}_{\xi} \\
&\quad + \left(1 + \frac{a_0}{c^2} \xi^1 \right)^2 \vec{\nabla}_{\xi} \times \vec{\nabla}_{\xi} \times \vec{E}_{\xi}
\end{aligned}$$

$$\begin{aligned}
&= \vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) \times \vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) \times \vec{E}_\xi + \left(1 + \frac{a_0}{c^2} \xi^1\right)^2 \vec{\nabla}_\xi \times \vec{\nabla}_\xi \times \vec{E}_\xi \\
&= [\vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) \cdot \vec{E}_\xi] \vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) - [\vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) \cdot \vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right)] \vec{E}_\xi \\
&\quad + \left(1 + \frac{a_0}{c^2} \xi^1\right)^2 [\vec{\nabla}_\xi (\vec{\nabla}_\xi \cdot \vec{E}_\xi) - \nabla_\xi^2 \vec{E}_\xi] \\
&= -\frac{1}{c} \frac{\partial}{\partial \xi^0} [\vec{\nabla}_\xi \times \{\vec{B}_\xi \left(1 + \frac{a_0 \xi^1}{c^2}\right)\}] = -\frac{1}{c^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 \vec{E}_\xi,
\end{aligned}$$

$$\text{In this time, } \vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) = \left(\frac{a_0}{c^2}, 0, 0\right) \quad (60)$$

Hence,

$$\begin{aligned}
&\vec{\nabla}_\xi \times \left(1 + \frac{a_0}{c^2} \xi^1\right) \vec{\nabla}_\xi \times \{\vec{E}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right)\} + \frac{1}{c^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 \vec{E}_\xi \\
&= [\vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) \cdot \vec{E}_\xi] \vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) - [\vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) \cdot \vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right)] \vec{E}_\xi \\
&\quad + \left(1 + \frac{a_0}{c^2} \xi^1\right)^2 [\vec{\nabla}_\xi (\vec{\nabla}_\xi \cdot \vec{E}_\xi) - \nabla_\xi^2 \vec{E}_\xi] + \frac{1}{c^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 \vec{E}_\xi \\
&= \frac{a_0^2}{c^4} (E_{\xi^1}, 0, 0) - \frac{a_0^2}{c^4} (E_{\xi^1}, E_{\xi^2}, E_{\xi^3}) + \left(1 + \frac{a_0}{c^2} \xi^1\right)^2 \left[\frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 - \nabla_\xi^2\right] \vec{E}_\xi \\
&= \frac{a_0^2}{c^4} (0, -E_{\xi^2}, -E_{\xi^3}) + \left(1 + \frac{a_0}{c^2} \xi^1\right)^2 \left[\frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 - \nabla_\xi^2\right] \vec{E}_\xi \\
&= \frac{a_0^2}{c^4} (0, -E_{\xi^2}, -E_{\xi^3}) + \left(1 + \frac{a_0}{c^2} \xi^1\right)^2 \left[\frac{1}{c^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 - \nabla_\xi^2\right] \vec{E}_\xi \\
&= \vec{0} \quad (61)
\end{aligned}$$

Hence, the magnetic wave equation is in vacuum

$$\begin{aligned}
&\vec{\nabla}_\xi \times \left(1 + \frac{a_0}{c^2} \xi^1\right) \vec{\nabla}_\xi \times \{\vec{B}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right)\} + \frac{1}{c^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 \vec{B}_\xi \\
&= [\vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) \cdot \vec{B}_\xi] \vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) - [\vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right) \cdot \vec{\nabla}_\xi \left(1 + \frac{a_0}{c^2} \xi^1\right)] \vec{B}_\xi \\
&\quad + \left(1 + \frac{a_0}{c^2} \xi^1\right)^2 [\vec{\nabla}_\xi (\vec{\nabla}_\xi \cdot \vec{B}_\xi) - \nabla_\xi^2 \vec{B}_\xi] + \frac{1}{c^2} \left(\frac{\partial}{\partial \xi^0}\right)^2 \vec{B}_\xi
\end{aligned}$$

$$\begin{aligned}
&= \frac{a_0^2}{c^4} (0, -B_{\xi^2}, -B_{\xi^3}) + (1 + \frac{a_0}{c^2} \xi^1)^2 \left[\frac{1}{c^2} \frac{1}{(1 + \frac{a_0}{c^2} \xi^1)^2} (\frac{\partial}{\partial \xi^0})^2 - \nabla_{\xi}^2 \right] \vec{B}_{\xi} \\
&= \frac{a_0^2}{c^4} (0, -B_{\xi^2}, -B_{\xi^3}) + (1 + \frac{a_0}{c^2} \xi^1)^2 \left[\frac{1}{c^2} (\frac{\partial}{\partial \xi^0})^2 - \nabla_{\xi}^2 \right] \vec{B}_{\xi} \\
&= \vec{0}
\end{aligned} \tag{62}$$

The electromagnetic wave function, Eq(57),Eq(58) satisfy the electromagnetic wave equation, Eq(61),Eq(62).

5. Conclusion

We find the electro-magnetic field transformation and the electro-magnetic equation in uniformly accelerated frame.

Generally, the coordinate transformation of accelerated frame is

$$\begin{aligned}
(I) \quad ct &= \left(\frac{c^2}{a_0} + \xi^1 \right) \sinh \left(\frac{a_0 \xi^0}{c} \right) \\
x &= \left(\frac{c^2}{a_0} + \xi^1 \right) \cosh \left(\frac{a_0 \xi^0}{c} \right) - \frac{c^2}{a_0} , y = \xi^2 , z = \xi^3
\end{aligned} \tag{63}$$

$$\begin{aligned}
(II) \quad ct &= \frac{c^2}{a_0} \exp \left(\frac{a_0}{c^2} \xi^1 \right) \sinh \left(\frac{a_0 \xi^0}{c} \right) \\
x &= \frac{c^2}{a_0} \exp \left(\frac{a_0}{c^2} \xi^1 \right) \cosh \left(\frac{a_0 \xi^0}{c} \right) - \frac{c^2}{a_0} , y = \xi^2 , z = \xi^3
\end{aligned} \tag{64}$$

Hence, this article say the accelerated frame is Rindler coordinate (I) that can treat electro-magnetic field equation.

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