

# On pressure of fation gas

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## 1. Introduction.

According to the classical kinetic theory [1], pressure of gas on a closed container wall is equals to

$$p = \frac{2}{3} \varepsilon, \quad (1)$$

where  $\varepsilon$  is the gas energy density.

Here we give a derivation of the formula for the pressure of the perfect gas filling the infinite space. I have used this formula in [2] for pressure of the fation gas in the theory of the shadow-gravity. It is supposed that faion gas is the perfect gas, which consists of neutral sub-particles, free pass of which in the gas is very large. Ibid I also used the notion about fundamental sub-particles (FSP), from which all substance consists. I suggest to consider as fundamental such sub-particles, which are *absolutely impenetrable* for fations.

## 2. Derivation of a new expression

Let the fation gas flow, directed within limits of the element  $d\Omega$  of the solid angle  $\Omega$ , falls on the infinity small surface element,  $ds$ , of FSP (Fig. 1). The normal component of its momentum is equal to

$$dP = (2 - \delta)(\varepsilon^* / c)V \cos \Omega_p d\Omega = (2 - \delta)(\varepsilon^* / c) \cos^2 \Omega_p a d\Omega ds, \quad (2)$$

where  $\Omega_p$  is the plane angle,  $\varepsilon^*$  is the energy density of the uniformly directed flow of fations,  $\delta$  is the probability of fation absorption by the body surface.  $V = ads \cos \Omega_p$  is the volume of the oblique cylinder,  $a = ct$ ,  $t$  is the time interval. Since  $dP/t$  is the element of force,  $dF$ , and  $dF/ds = dp$  is the element of pressure, we obtain from (2)

$$dp = (2 - \delta) \varepsilon^* \cos^2 \Omega_p d\Omega. \quad (3)$$

Solid angle  $\Omega$  is equal to ratio of the spherical segment area,  $S = 2\pi a^2 (1 - \cos \Omega_p)$ , to  $a^2$ , i.e.

$$\Omega = \frac{S}{a^2} = 2\pi (1 - \cos \Omega_p). \quad (4)$$

By differentiating (4) with respect to  $\Omega_p$  we obtain

$$d\Omega = 2\pi \sin \Omega_p d\Omega_p. \quad (5)$$

After substituting (5) in (3) we find

$$dp = 2\pi(2 - \delta) \varepsilon^* \sin \Omega_p \cos^2 \Omega_p d\Omega_p. \quad (6)$$

Finally, after integrating (6), we obtain expression for pressure as

$$p = 2\pi(2 - \delta) \varepsilon^* \int_0^{\pi/2} \sin \Omega_p \cos^2 \Omega_p d\Omega_p = \frac{(2 - \delta) \varepsilon^*}{2} \left( -\frac{\cos^3 \Omega_p}{3} \right) \Bigg|_0^{\pi/2} \cong 1/3 \varepsilon, \quad (7)$$

where  $\varepsilon = 4\pi \varepsilon^*$  is the general energy density of the omnidirectional flows of fations. We also have taken into account that  $\delta \sim 10^{-42} \ll 2$  [2].

As is seen, the new expression differs from that (1) obtained for the closed container.

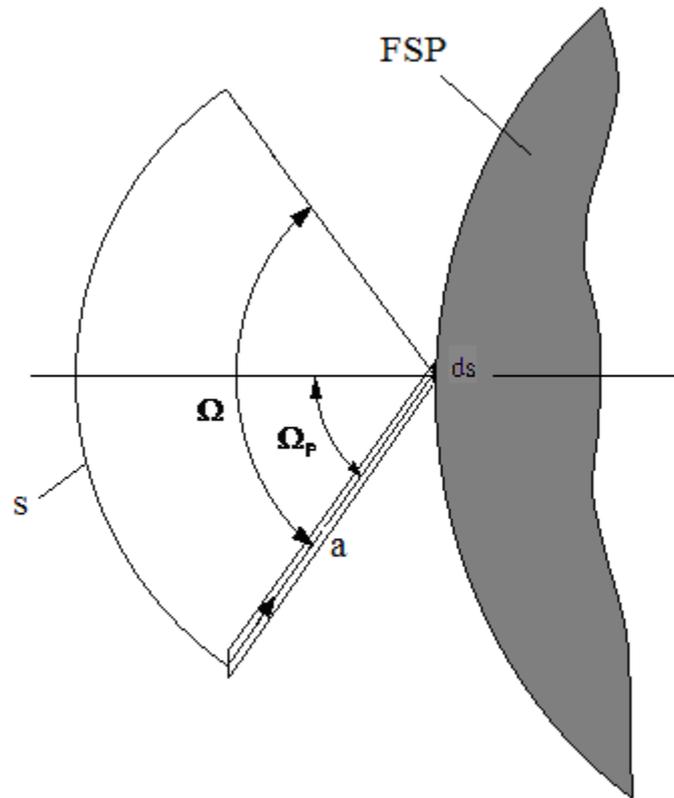


Fig. 1. Diagram of fations falling on the surface element,  $ds$ , of FSP.

### References

- [1] Reif, F. *Statistical Physics* (Berkeley Physics Course, New York, 1967), Vol. 5
- [2] Nikolay V. Dibrov, "Exact Formula for Shadow-Gravity, Strong Gravity" (2013), [viXra.org e-Print archive, viXra:1309.0175](https://arxiv.org/abs/1309.0175), Exact Formula [for ..](#)