

# **Dynamic Instability of the Standard Model and the Fine-Tuning Problem**

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## *Abstract*

The Standard Model for particle physics (SM) is a nonlinear field theory in which both Yang-Mills and Higgs bosons are self-interacting objects. Their classical or quantum evolution is inevitably sensitive to the transition from order to chaos. With some noteworthy exceptions, the mainstream of theoretical particle physics has ignored the dynamical contribution of chaos in Quantum Field Theory. Here we point out that quantum corrections to the classical interaction of the Higgs with gauge bosons may lower the threshold for the onset of chaos and destabilize the vacuum in the *low* or *intermediate* TeV scale. The inability of the vacuum to survive in this energy region hints to a straightforward solution for the fine-tuning problem. It also implies that perturbative estimates on vacuum stability well above the LHC scale are likely to be invalid.

**Key words:** Quantum Field Theory, Standard Model, vacuum stability, fine-tuning problem, gauge fields, Higgs boson, order-to-chaos transition.

## **1. Introduction**

By construction, the SM represents a *nonlinear field theory* in which Yang-Mills (YM) and Higgs bosons are self-interacting. Nonlinear dynamics of such objects is present at both classical and quantum levels. The chaotic attributes of YM fields have been known and studied since the beginning of the eighties [1-11]<sup>1</sup>. Chaos was first analyzed in the classical limit of the YM theory

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<sup>1</sup> Due to the large number of contributions on this topic, we have opted to list only a representative sample of recent publications. As a result, many relevant references are not included here.

and it was shown to exist in both the continuum and lattice formulations of the theory. Particularly, for homogeneous gauge field configurations, it was found that spontaneous symmetry breaking triggers the *transition to chaos* (TC) with the rise of the energy density, whereas the dynamics of YM fields in the absence of spontaneous symmetry breaking remains chaotic at any density of energy. The emergence of chaos in classical dynamics of the  $SU(2) \times U(1)$  theory was numerically explored in [9]. Follow-up research was focused on understanding the TC in the semi-classical regime of quantum mechanics (QM) as well as in quantum field theory (QFT). The investigation of chaos in classical gauge theory has later targeted on a number of specific problems. One of them was confirming the effect induced by the Higgs on the chaotic dynamics of classical YM theory. It was shown that the Higgs scalar regularizes the dynamics of gauge fields at low energy densities [4, 7-8]. It was also discovered that quantum fluctuations leading to symmetry breaking via the Coleman–Weinberg mechanism tend to stabilize chaotic dynamics of spatially homogeneous systems of YM and Higgs fields at low energy densities [7-8, 10-11]. The connection between the chaotic dynamics of a classical field theory and the instability of the one-loop effective action of the associated QFT was analyzed in [10].

Surprisingly, aside from studies like the ones previously cited, the majority of theoretical models in particle physics have overlooked the implications of chaos in QFT. The goal of our report is to contribute to a reversal of this trend. We emphasize here that one-loop corrections to the classical interaction of the Higgs with W, Z bosons or photons are likely to lower the threshold for the TC and destabilize the vacuum in the low or intermediate TeV scale. *By default, a rapidly decaying vacuum in this energy region explains away the long-standing problems of fine-tuning and ultraviolet stability of the SM.*

Our report is organized as follows: section two reviews the Higgs potential and the fine-tuning problem of the SM. Estimates on vacuum stability based on extrapolation of the SM near the Planck scale are briefly addressed in section three. Section four highlights details on the TC that may be relevant for the dynamics of the di-boson and di-photon decay channels. Concluding remarks are presented in the last section.

This work represents a continuation of several studies initiated by the author in [12-15]. It is preliminary in nature and calls out for further clarifications and revisions. Concurrent efforts may refute, amend or consolidate our findings.

## **2. Stability of the Higgs potential and the fine-tuning problem**

Electroweak (EW) symmetry in the SM is broken by a scalar field having the following doublet structure [16]:

$$\Phi = \begin{pmatrix} G^+ \\ \frac{1}{\sqrt{2}}[(H + v) + iG^0] \end{pmatrix} \quad (2.1)$$

Here,  $G^+$  and  $G^0$  represent the charged and neutral Goldstone bosons arisen from spontaneous symmetry breaking,  $H$  is the SM Higgs boson,  $v \approx O(M_{EW}) = 246 \text{ GeV}$  is the Higgs vacuum expectation value (vev) and  $M_{EW}$  stands for the EW scale. Symmetry breaking is caused by the Higgs potential, whose form satisfies the requirements of renormalizability and gauge-invariance:

$$V = \mu^2 \Phi^+ \Phi + \lambda (\Phi^+ \Phi)^2 \quad (2.2)$$

with  $\lambda \approx O(1)$  and  $\mu^2 \approx O(M_{EW}^2)$ . A vanishing quartic coupling ( $\lambda = 0$ ) represents the critical value that separates the ordinary EW phase from an unphysical phase where the Higgs field

assumes unbounded values. Likewise, the coefficient  $\mu^2$  plays the role of an order parameter whose sign describes the transition between a symmetric phase and a broken phase. Minimizing the Higgs potential yields a vev given by:

$$v^2 = -(\mu^2/\lambda) \quad (2.3)$$

where the physical mass of the Higgs is:

$$M_H^2 = -2\lambda v^2 = 2\mu^2 \quad (2.4)$$

The renormalized mass squared of the Higgs scalar contains two contributions:

$$\mu^2 = \mu_0^2 + \Delta\mu^2 \quad (2.5)$$

in which  $\mu_0^2$  represents the ultraviolet (bare) value. This mass parameter picks up quantum corrections  $\Delta\mu^2$  that depend quadratically on the ultraviolet cutoff  $\Lambda$  of the theory. Consider for example the contribution of radiative corrections to  $\mu^2$  from top quarks. The complete one-loop calculation of this contribution reads:

$$\Delta\mu^2 = \frac{N_c \lambda_t^2}{16\pi^2} [-2\Lambda^2 + 6M_t^2 \ln(\frac{\Lambda}{M_t}) + \dots] \quad (2.5)$$

in which  $\lambda_t$  and  $M_t$  are the Yukawa coupling and mass of the top quark. If the bare Higgs mass is set near the cutoff  $\mu_0^2 = O(\Lambda^2) = O(M_{Pl}^2)$ , then  $\Delta\mu^2 \approx -10^{35}$  GeV. This large correction must precisely cancel against  $\mu_0^2$  to protect the EW scale. This is the root cause of the *fine-tuning problem*, which boils down to the implausible requirement that  $\mu_0^2$  and  $\Delta\mu^2$  should offset each other to about 31 decimal places.

Closely related to the fine-tuning problem is the question of whether the SM remains valid all the way up to the Planck scale ( $M_{Pl}$ ). This question is non-trivial because it depends on how the Higgs quartic coupling  $\lambda$  behaves at high energies. Competing trends are at work here, namely [17-19]:

- 1) Radiative corrections from top quarks *drop*  $\lambda$  at higher scales, while those from the self-interacting Higgs *grow*  $\lambda$  at higher scales.
- 2) If  $\lambda(M_{EW})$  is too large, the Higgs loops dominate and  $\lambda$  diverges at some intermediate scale called the *Landau pole*. However, if  $\lambda(M_{EW})$  is too small, the top loops dominate,  $\lambda$  runs negative at some intermediate scale which, in turn, makes the potential unbounded from below and destabilizes the vacuum.

### **3. Radiative corrections and the vacuum stability**

Within the SM, the value of the physical Higgs mass  $M_H \approx 125 \text{ GeV}$  hinted by recent LHC data falls at the border of vacuum stability which, in turn, implies a vanishing quartic coupling near  $M_{Pl}$ . A recent study [17] has undertaken a complete perturbative analysis on the vacuum stability, including the two-loop threshold correction to  $\lambda$  at the EW scale due to QCD and top quark couplings. This analysis has unveiled the following outcomes:

- 1) Vacuum instability develops around a Renormalization Group scale on the order of  $\mu_{cr} = O(10^{11}) \text{ GeV}$ .
- 2) Both parameters of the Higgs potential (2.2) assume near-critical values about  $\mu_{cr}$

$$\mu^2 \ll M_{Pl}, \quad \lambda(\mu_{cr}) \approx 0 \quad (3.1)$$

It was concluded that (3.1) hints at the possibility that the SM behaves as a statistical system approaching criticality at  $\mu_{cr}$  [17]. The next section attempts to refute this conclusion. We find that critical behavior and the approach to chaos in the SM are likely to occur at a scale *appreciably lower* than  $\mu_{cr}$ .

#### **4. Transition to chaos in Higgs decay channels**

The goal of this section is to briefly survey several scenarios describing the TC in systems comprising YM or Maxwell fields (M) in interaction with the Higgs scalar. These scenarios may be relevant for the dynamics of decay channels in which Higgs breaks up in pairs of YM bosons or photons.

4.1) Let us begin by bringing up a numerical study of the TC for the classical and homogeneous SU(2) x U(1) theory [9]. It was found there that the critical energy density for the onset of chaos is given by:

$$\varepsilon_c = \left( \frac{4M_W^2}{\sqrt{2}G_F} \right) \varepsilon \quad (4.1)$$

where  $M_W$  is the mass of the  $W$  boson,  $G_F$  the Fermi constant and  $\varepsilon$  the energy density of the system of Yang-Mills and Higgs bosons. It follows from (4.1) that  $\varepsilon_c \approx 2 \times 10^{10} \text{ GeV/fm}^3$  or  $\varepsilon_c \approx 1.529 \times 10^8 \text{ GeV}^4$ . Multiplying this energy density by the volume corresponding to the EW interaction range  $V_{EW} \approx r^3$  with  $r = O(10^{-3} \text{ fm})$ , results in a critical energy for the TC on the order of  $100 \text{ GeV} \approx O(M_{EW})$ .

4.2) We turn next to the case of massless scalar electrodynamics coupled to the Higgs scalar. One-loop corrections to the classical interaction of this system lead to the following critical energy density [10]:

$$\varepsilon_c \approx \frac{11}{108} \lambda v^4 - \frac{4}{81\pi^2} (\lambda^2 + \frac{3e^4}{10}) v^4 \quad (4.2)$$

If the quartic coupling is evaluated at the pole top mass [17] and the system is assumed to be near the EW scale, that is, if  $\lambda = 0.126$  and  $v \approx 246$  GeV, (4.2) leads to  $\varepsilon_c = 4.658 \times 10^7$  GeV<sup>4</sup> and a critical energy for the TC on the order of 10 GeV.

4.3) Finally, let us consider a system of EW bosons coupled to a massless Higgs scalar [11]. The critical energy density defined there includes a contribution from radiative corrections to the Higgs potential ( $\varepsilon_c^0$ ) and a contribution from the vacuum energy density ( $\varepsilon_v$ ). These are estimated to be, respectively:

$$\varepsilon_c^0 = \frac{3\eta^4}{64\pi^2} \exp[2(\alpha_w - \frac{\lambda}{g^4} \beta_w)] (1 + \frac{1}{2 \cos^4 \theta_w}) [1 - 7 \exp(-48\Lambda_w \beta_w)] \quad (4.3)$$

$$\varepsilon_v = O(\lambda \eta^4)$$

Here  $\eta$  stands for the renormalization scale,  $g$  for the weak coupling of the EW bosons,  $\theta_w$  is the Weinberg angle and  $\alpha_w, \beta_w, \Lambda_w$  are numerical coefficients dependent on  $\theta_w$ . Taking  $\lambda = 0.126$ , assuming  $\eta \approx O(v)$  and adding the two contributions gives  $\varepsilon_c = 4.581 \times 10^8$  GeV<sup>4</sup> and again a critical energy on the order of 100 GeV  $\approx O(M_{EW})$ .

The table below summarizes all results discussed in this section.

Relation	Context	Critical energy density ( $\text{GeV}^4$ )	Critical energy (GeV)
(4.1)	classical YM-Higgs	$1.529 \times 10^8$	$O(10^2)$
(4.2)	quantum M-Higgs	$4.658 \times 10^7$	$O(10)$
(4.3)	quantum YM-Higgs	$4.581 \times 10^8$	$O(10^2)$

#### **4. Conclusions**

Our findings suggest that the transition from order to chaos in classical and quantum systems of gauge and Higgs fields is prone to occur at a scale substantially lower than  $\mu_{cr} = O(10^{11})$  GeV. Quantum corrections from the Higgs quartic coupling and from the interaction of the Higgs with heavy particles become irrelevant as the vacuum loses stability and dies out. The inability of the vacuum to survive not too far above the LHC scale explains away the fine-tuning problem and signals the breakdown of the SM in this region.

This conclusion, albeit preliminary, sheds light on the problem of *extrapolating* field theories in the deep ultraviolet sector, specifically near the Planck scale. As it is well known, General Relativity (GR) is exclusively an “effective” low-energy framework and efforts to develop perturbative quantization of classical gravity result in non-renormalizable theories [20]. So far, the ultraviolet completion of GR has been approached from two different directions. Whereas string theories pursues introduction of new fields and symmetries beyond the SM, some quantum gravity theories retain all fields and symmetries of GR and treat gravity as a fundamental non-perturbative interaction. Since gravity cannot be decoupled from the energy-momentum of matter at any level of description, vacuum instability in the TeV sector points out that both approaches to the ultraviolet completion of GR may be missing a critical piece of the puzzle.

As mentioned in the first section, our conclusions need to be further scrutinized and, most importantly, confronted with the experiment. For instance, it is necessary to clarify the relationship between the standard Higgs mechanism of EW symmetry breaking and the vacuum stability problem. We anticipate that this effort may help answer the following open questions:

- a) Is the observed di-photon excess of Higgs decays rooted in the vacuum instability near the LHC scale?
- b) Can one envision that the Higgs scalar is a *short-lived cluster of gauge bosons* (EW or gluons) rather than a fundamental field postulated by the theory, whose function is to preserve the consistency of the SM?
- c) Would this interpretation bring us closer to understanding the physical origin of the Higgs mass?

The results of this investigation are planned to be reported elsewhere [21].

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