

The differential coordinate transformation in the general relativity theory

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ABSTRACT

:In the general relativity theory's the gravity field,in the vacuum, save the motive particle's coordinate systems' the differential coordinate transformation.And study the velocity add formula in the general relativity theory

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I.Introduction

In the general relativity theory, Schwarzschild solution is

$$\begin{aligned}
d\tau^2 &= \left(1 - \frac{2GM}{rc^2}\right)dt^2 - \frac{1}{c^2} \left[\frac{dr^2}{1 - \frac{2GM}{rc^2}} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right] \quad (1) \\
d\tau^2 &= \left(1 - \frac{2GM}{rc^2}\right)dt^2 - \frac{1}{c^2} \left[\frac{dr^2}{1 - \frac{2GM}{rc^2}} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right] \\
&= \left(1 - \frac{2GM}{rc^2}\right)dt^2 \left[1 - \frac{1}{c^2} \left\{ \left(\frac{dr}{dt}\right)^2 \frac{1}{\left(1 - \frac{2GM}{rc^2}\right)^2} + r^2 \left(\frac{d\theta}{dt}\right)^2 \frac{1}{\left(1 - \frac{2GM}{rc^2}\right)} \right. \right. \\
&\quad \left. \left. + r^2 \sin^2 \theta \left(\frac{d\phi}{dt}\right)^2 \frac{1}{\left(1 - \frac{2GM}{rc^2}\right)} \right\} \right] \\
&= \left(1 - \frac{2GM}{rc^2}\right)dt^2 \left[1 - \frac{1}{c^2} (\bar{V}_r^2 + \bar{V}_\theta^2 + \bar{V}_\phi^2) / c^2 \right] \\
&= \left(1 - \frac{2GM}{rc^2}\right)dt^2 [1 - \bar{V}^2 / c^2], \\
\bar{V}_r &= \frac{dr}{dt} \frac{1}{\left(1 - \frac{2GM}{rc^2}\right)}, \bar{V}_\theta = r \frac{d\theta}{dt} \frac{1}{\sqrt{1 - \frac{2GM}{rc^2}}}, \bar{V}_\phi = r \sin \theta \frac{d\phi}{dt} \frac{1}{\sqrt{1 - \frac{2GM}{rc^2}}} \\
\bar{V}^2 &= \bar{V}_r^2 + \bar{V}_\theta^2 + \bar{V}_\phi^2 \quad (2)
\end{aligned}$$

But the real velocity V is

$$\begin{aligned}
d\tau^2 &= \left(1 - \frac{2GM}{rc^2} - \frac{V^2}{c^2}\right)dt^2, \\
V^2 &= \left(\frac{ds}{dt}\right)^2 = \frac{1}{1 - \frac{2GM}{rc^2}} \left(\left(\frac{dr}{dt}\right)^2 + r^2 \left(\frac{d\theta}{dt}\right)^2 + r^2 \sin^2 \theta \left(\frac{d\phi}{dt}\right)^2 \right. \\
&\quad \left. = V_r^2 + V_\theta^2 + V_\phi^2 \right) \\
V_r &= \frac{dr}{dt} \frac{1}{\sqrt{1 - \frac{2GM}{rc^2}}}, V_\theta = r \frac{d\theta}{dt}, V_\phi = r \sin \theta \frac{d\phi}{dt}
\end{aligned}$$

$$ds^2 = \frac{1}{1 - \frac{2GM}{rc^2}} dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

Hence

$$\bar{V} = \frac{V}{\sqrt{1 - \frac{2GM}{rc^2}}} \quad (3)$$

II.Additional chapter

In this case, think that a particle move. And think the coordinate systems by the particle's point and velocity. In this time, think the coordinate systems' transformation likely the special relativity theory. In the formula (2), on the coordinate system $S(t, r)$, move the other coordinate system $S'(t', r')$ by the real velocity component u_r, u_θ, u_ϕ , exist a differential coordinate transformation of the coordinate system $S(t, r)$ and the other coordinate system $S'(t', r')$. In this time, the differential coordinate transformation of the coordinate system is likely the Lorentz transformation.

$$\begin{aligned} \sqrt{1 - \frac{2GM}{rc^2}} dt &= \gamma \left(\sqrt{1 - \frac{2GM}{r'c^2}} dt' + \frac{\bar{u}_r}{c^2} \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \frac{\bar{u}_\theta}{c^2} r' d\theta' + \frac{\bar{u}_\phi}{c^2} r' \sin \theta' d\phi' \right), \\ \frac{dr}{\sqrt{1 - \frac{2GM}{rc^2}}} &= \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \gamma \bar{u}_r dt' \sqrt{1 - \frac{2GM}{r'c^2}} \\ &\quad - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta r' d\theta' + \bar{u}_\phi r' \sin \theta' d\phi') \bar{u}_r, \\ rd\theta &= r' d\theta' + \gamma \bar{u}_\theta dt' \sqrt{1 - \frac{2GM}{r'c^2}} \\ &\quad - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta r' d\theta' + \bar{u}_\phi r' \sin \theta' d\phi') \bar{u}_\theta, \\ r \sin \theta d\phi &= r' \sin \theta d\phi' + \gamma \bar{u}_\phi dt' \sqrt{1 - \frac{2GM}{r'c^2}} \\ &\quad - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta r' d\theta' + \bar{u}_\phi r' \sin \theta' d\phi') \bar{u}_\phi, \\ \bar{u}_r &= u_r / \sqrt{1 - \frac{2GM}{rc^2}}, \bar{u}_\theta = u_\theta / \sqrt{1 - \frac{2GM}{rc^2}}, \bar{u}_\phi = u_\phi / \sqrt{1 - \frac{2GM}{rc^2}}, \bar{u} = u / \sqrt{1 - \frac{2GM}{rc^2}} \end{aligned}$$

$$\bar{u}^2 = \bar{u}_r^2 + \bar{u}_\theta^2 + \bar{u}_\phi^2, \quad u^2 = u_r^2 + u_\theta^2 + u_\phi^2, \quad \gamma = \frac{1}{\sqrt{1 - \frac{\bar{u}^2}{c^2}}} \quad (4)$$

The formula (4) insert the formula (1),

$$\begin{aligned}
d\tau^2 &= (1 - \frac{2GM}{rc^2})dt^2 - \frac{1}{c^2} \left[\frac{dr^2}{1 - \frac{2GM}{rc^2}} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \right] \\
&= \gamma^2 \left(\sqrt{1 - \frac{2GM}{r'c^2}} dt' + \frac{\bar{u}_r}{c^2} \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \frac{\bar{u}_\theta}{c^2} r' d\theta' + \frac{\bar{u}_\phi}{c^2} r' \sin \theta' d\phi' \right)^2 \\
&\quad - \frac{1}{c^2} \left[\frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \gamma \bar{u}_r dt' \sqrt{1 - \frac{2GM}{r'c^2}} \right. \\
&\quad \left. - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta r' d\theta' + \bar{u}_\phi r' \sin \theta' d\phi') \bar{u}_r \right]^2 \\
&\quad - \frac{1}{c^2} \left[r' d\theta' + \gamma \bar{u}_\theta dt' \sqrt{1 - \frac{2GM}{r'c^2}} \right. \\
&\quad \left. - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta r' d\theta' + \bar{u}_\phi r' \sin \theta' d\phi') \bar{u}_\theta \right]^2 \\
&\quad - \frac{1}{c^2} \left[r' \sin \theta d\phi' + \gamma \bar{u}_\phi dt' \sqrt{1 - \frac{2GM}{r'c^2}} \right. \\
&\quad \left. - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta r' d\theta' + \bar{u}_\phi r' \sin \theta' d\phi') \bar{u}_\phi \right]^2 \\
&= (1 - \frac{2GM}{r'c^2})dt'^2 - \frac{1}{c^2} \left[\frac{dr'^2}{1 - \frac{2GM}{r'c^2}} + r'^2 d\theta'^2 + r'^2 \sin^2 \theta' d\phi'^2 \right] \\
&\quad \bar{u}_r = u_r / \sqrt{1 - \frac{2GM}{rc^2}}, \quad \bar{u}_\theta = u_\theta / \sqrt{1 - \frac{2GM}{rc^2}}, \quad \bar{u}_\phi = u_\phi / \sqrt{1 - \frac{2GM}{rc^2}}, \quad \bar{u} = u / \sqrt{1 - \frac{2GM}{rc^2}} \\
&\quad \bar{u}^2 = \bar{u}_r^2 + \bar{u}_\theta^2 + \bar{u}_\phi^2, \quad u^2 = u_r^2 + u_\theta^2 + u_\phi^2, \quad \gamma = \frac{1}{\sqrt{1 - \frac{\bar{u}^2}{c^2}}} \quad (5)
\end{aligned}$$

If $d\theta = d\phi = 0$ is in the formula (1), treat the light,

$$d\tau^2 = (1 - \frac{2GM}{rc^2})dt^2 - \frac{1}{c^2} \frac{dr^2}{1 - \frac{2GM}{rc^2}} = 0, d\theta = d\phi = 0$$

$$cdt = \frac{dr}{(1 - \frac{2GM}{rc^2})}, \quad ct = r + \frac{2GM}{c^2} \ln |r - \frac{2GM}{c^2}| \quad (6)$$

In the r -axis, in the motive coordinate system $S(t, r)$ and $S'(t', r')$, Energy-momentum transformation is by the formula (4)

$$\begin{aligned} E &= m_0 c^2 \frac{dt}{d\tau}, \quad E' = m_0 c^2 \frac{dt'}{d\tau} \\ p^r &= m_0 \frac{dr}{d\tau}, \quad p^{r'} = m_0 \frac{dr'}{d\tau}, \quad p^\theta = m_0 r \frac{d\theta}{d\tau}, \quad p'^\theta = m_0 r' \frac{d\theta'}{d\tau} \\ p^\phi &= m_0 r \sin \theta \frac{d\phi}{d\tau}, \quad p'^\phi = m_0 r' \sin \theta' \frac{d\phi'}{d\tau} \\ \sqrt{1 - \frac{2GM}{rc^2}} E &= \gamma \left(\sqrt{1 - \frac{2GM}{r'c^2}} E' + \frac{\bar{u}_r p^{r'}}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta p'^\theta + \bar{u}_\phi p'^\phi \right), \\ \frac{p^r}{\sqrt{1 - \frac{2GM}{rc^2}}} &= \frac{p^{r'}}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \gamma \frac{\bar{u}_r}{c^2} E' \sqrt{1 - \frac{2GM}{r'c^2}} \\ &\quad - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r - \frac{p^{r'}}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta p'^\theta + \bar{u}_\phi p'^\phi) \bar{u}_r, \\ p^\theta &= p'^\theta + \gamma \frac{\bar{u}_\theta}{c^2} E' \sqrt{1 - \frac{2GM}{r'c^2}} \\ &\quad - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r - \frac{p^{r'}}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta p'^\theta + \bar{u}_\phi p'^\phi) \bar{u}_\theta, \\ p^\phi &= p'^\phi + \gamma \frac{\bar{u}_\phi}{c^2} E' \sqrt{1 - \frac{2GM}{r'c^2}} \\ &\quad - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r - \frac{p^{r'}}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta p'^\theta + \bar{u}_\phi p'^\phi) \bar{u}_\phi, \\ \bar{u}_r &= u_r / \sqrt{1 - \frac{2GM}{rc^2}}, \quad \bar{u}_\theta = u_\theta / \sqrt{1 - \frac{2GM}{rc^2}}, \quad \bar{u}_\phi = u_\phi / \sqrt{1 - \frac{2GM}{rc^2}}, \quad \bar{u} = u / \sqrt{1 - \frac{2GM}{rc^2}} \end{aligned}$$

$$\bar{u}^2 = \bar{u}_r^2 + \bar{u}_\theta^2 + \bar{u}_\phi^2 , \quad u^2 = u_r^2 + u_\theta^2 + u_\phi^2 , \quad \gamma = \frac{1}{\sqrt{1 - \frac{\bar{u}^2}{c^2}}} \quad (7)$$

$$\begin{aligned}
& (1 - \frac{2GM}{rc^2})E^2 - \frac{1}{1 - \frac{2GM}{rc^2}}(p^r)^2 c^2 - (p^\theta)^2 c^2 - (p^\phi)^2 c^2 \\
& = (1 - \frac{2GM}{rc^2})m_0^2 c^4 (\frac{dt}{d\tau})^2 - \frac{1}{1 - \frac{2GM}{rc^2}}m_0^2 c^2 (\frac{dr}{d\tau})^2 - m_0^2 c^2 r^2 (\frac{d\theta}{d\tau})^2 \\
& \quad - m_0^2 c^2 r^2 \sin^2 \theta (\frac{d\phi}{d\tau})^2 \\
& = m_0^2 c^4 [(1 - \frac{2GM}{rc^2})dt^2 - \frac{1}{c^2} \{ \frac{dr^2}{1 - \frac{2GM}{rc^2}} + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \}] / d\tau^2 \\
& = m_0^2 c^4 , E = m_0 c^2 \frac{dt}{d\tau} , p^r = m_0 \frac{dr}{d\tau} , p^\theta = m_0 r \frac{d\theta}{d\tau} , p^\phi = m_0 r \sin \theta \frac{d\phi}{d\tau} \quad (8)
\end{aligned}$$

$$\begin{aligned}
m_0^2 c^4 & = (1 - \frac{2GM}{rc^2})E^2 - \frac{1}{1 - \frac{2GM}{rc^2}}(p^r)^2 c^2 - (p^\theta)^2 c^2 - (p^\phi)^2 c^2 \\
& = \gamma^2 (\sqrt{1 - \frac{2GM}{r'c^2}} E' + \frac{\bar{u}_r p'^r}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta p'^\theta + \bar{u}_\phi p'^\phi)^2 \\
& \quad - c^2 [\frac{p'^r}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \gamma \frac{\bar{u}_r}{c^2} E' \sqrt{1 - \frac{2GM}{r'c^2}} \\
& \quad - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r \frac{p'^r}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta p'^\theta + \bar{u}_\phi p'^\phi) \bar{u}_r]^2 \\
& \quad - c^2 [p'^\theta + \gamma \frac{\bar{u}_\theta}{c^2} E' \sqrt{1 - \frac{2GM}{r'c^2}} \\
& \quad - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r \frac{p'^r}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta p'^\theta + \bar{u}_\phi p'^\phi) \bar{u}_\theta]^2 , \\
& \quad - c^2 [p'^\phi + \gamma \frac{\bar{u}_\phi}{c^2} E' \sqrt{1 - \frac{2GM}{r'c^2}}
\end{aligned}$$

$$\begin{aligned}
& - (1 - \gamma) \frac{1}{\bar{u}^2} (\bar{u}_r \frac{p'^r}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u}_\theta p'^\theta + \bar{u}_\phi p'^\phi) \bar{u}_\phi]^2 \\
& = (1 - \frac{2GM}{r'c^2}) E'^2 - \frac{1}{1 - \frac{2GM}{r'c^2}} (p'^r)^2 c^2 - (p'^\theta)^2 c^2 - (p'^\phi)^2 c^2 \\
& E = m_0 c^2 \frac{dt}{d\tau}, p^r = m_0 \frac{dr}{d\tau}, p^\theta = m_0 r \frac{d\theta}{d\tau}, p^\phi = m_0 r \sin\theta \frac{d\phi}{d\tau} \\
& \bar{u}_r = u_r / \sqrt{1 - \frac{2GM}{rc^2}}, \bar{u}_\theta = u_\theta / \sqrt{1 - \frac{2GM}{rc^2}}, \bar{u}_\phi = u_\phi / \sqrt{1 - \frac{2GM}{rc^2}}, \bar{u} = u / \sqrt{1 - \frac{2GM}{rc^2}} \\
& \bar{u}^2 = \bar{u}_r^2 + \bar{u}_\theta^2 + \bar{u}_\phi^2, u^2 = u_r^2 + u_\theta^2 + u_\phi^2, \gamma = \frac{1}{\sqrt{1 - \frac{\bar{u}^2}{c^2}}} \quad (9)
\end{aligned}$$

In the r -axis, in the motive coordinate system $S(t, r)$ and $S'(t', r')$, by the formula (4) that if

$d\theta = d\phi = 0, d\theta' = d\phi' = 0, u_r = u, \bar{u}_r = \bar{u}, u_\theta = u_\phi = 0$ is, the velocity add formula is

$$\begin{aligned}
& \frac{dr}{\sqrt{1 - \frac{2GM}{rc^2}}} = \gamma \left(\frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u} dt' \sqrt{1 - \frac{2GM}{r'c^2}} \right) \\
& \sqrt{1 - \frac{2GM}{rc^2}} dt = \gamma \left(\sqrt{1 - \frac{2GM}{r'c^2}} dt' + \frac{\bar{u}}{c^2} \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} \right) \\
& \bar{u} = u / \sqrt{1 - \frac{2GM}{rc^2}}, \gamma = \frac{1}{\sqrt{1 - \frac{\bar{u}^2}{c^2}}}
\end{aligned}$$

$$d\theta = d\phi = 0, d\theta' = d\phi' = 0, u_r = u, \bar{u}_r = \bar{u}, u_\theta = u_\phi = 0$$

$$\begin{aligned}
& \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \bar{u} dt' \sqrt{1 - \frac{2GM}{r'c^2}} \\
& \frac{dr}{dt} \frac{1}{(1 - \frac{2GM}{rc^2})} = \frac{\sqrt{1 - \frac{2GM}{r'c^2}}}{dt' \sqrt{1 - \frac{2GM}{r'c^2}} + \frac{\bar{u}}{c^2} \frac{dr'}{\sqrt{1 - \frac{2GM}{r'c^2}}}}
\end{aligned}$$

$$\begin{aligned}
&= \frac{\frac{dr'}{\sqrt{1-\frac{2GM}{r'c^2}}} + \frac{u}{\sqrt{1-\frac{2GM}{rc^2}}} dt' \sqrt{1-\frac{2GM}{r'c^2}}}{dt' \sqrt{1-\frac{2GM}{r'c^2}} + \frac{1}{c^2} \frac{u}{\sqrt{1-\frac{2GM}{rc^2}}} \frac{dr'}{\sqrt{1-\frac{2GM}{r'c^2}}}} \\
&= \frac{\frac{dr'}{\sqrt{1-\frac{2GM}{r'c^2}}} + \frac{u/\sqrt{1-\frac{2GM}{rc^2}}}{\sqrt{1-\frac{2GM}{r'c^2}}} dt'(1-\frac{2GM}{r'c^2})}{dt' \frac{(1-\frac{2GM}{r'c^2})}{\sqrt{1-\frac{2GM}{r'c^2}}} + \frac{1}{c^2} \frac{dr'}{\sqrt{1-\frac{2GM}{r'c^2}}} u/\sqrt{1-\frac{2GM}{rc^2}}} \\
&= \frac{dr' + \frac{u}{\sqrt{1-\frac{2GM}{rc^2}}} dt'(1-\frac{2GM}{r'c^2})}{dt'(1-\frac{2GM}{r'c^2}) + \frac{dr'}{c^2} \frac{u}{\sqrt{1-\frac{2GM}{rc^2}}}}
\end{aligned}$$

$$d\theta = d\phi = 0, d\theta' = d\phi' = 0, u_r = u, \bar{u}_r = \bar{u}, u_\theta = u_\phi = 0 \quad (10)$$

III. Conclusion

Hence, in the general relativity theory, the velocity add formula is that if $d\theta = d\phi = 0$

, $d\theta' = d\phi' = 0, u_r = u, u_\theta = u_\phi = 0$ is

$$\begin{aligned}
\bar{V}_r &= \bar{V} = \frac{V_r}{\sqrt{1-\frac{2GM}{rc^2}}} = \frac{V}{\sqrt{1-\frac{2GM}{rc^2}}} = \frac{dr}{dt} \frac{1}{(1-\frac{2GM}{rc^2})} \\
&= \frac{\frac{dr'}{dt'} \frac{1}{(1-\frac{2GM}{r'c^2})} + \frac{u}{\sqrt{1-\frac{2GM}{rc^2}}}}{1 + \frac{1}{c^2} \cdot \frac{dr'}{dt'} \frac{1}{(1-\frac{2GM}{r'c^2})} \frac{u}{\sqrt{1-\frac{2GM}{rc^2}}}}
\end{aligned}$$

$$= \frac{\frac{V'}{\sqrt{1 - \frac{2GM}{r'c^2}}} + \frac{u}{\sqrt{1 - \frac{2GM}{rc^2}}}}{1 + \frac{1}{c^2} \cdot \frac{\frac{V'}{\sqrt{1 - \frac{2GM}{r'c^2}}}}{\frac{u}{\sqrt{1 - \frac{2GM}{rc^2}}}}} \quad (11)$$

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